

Corfu 1999

Particle & Condensed Matter

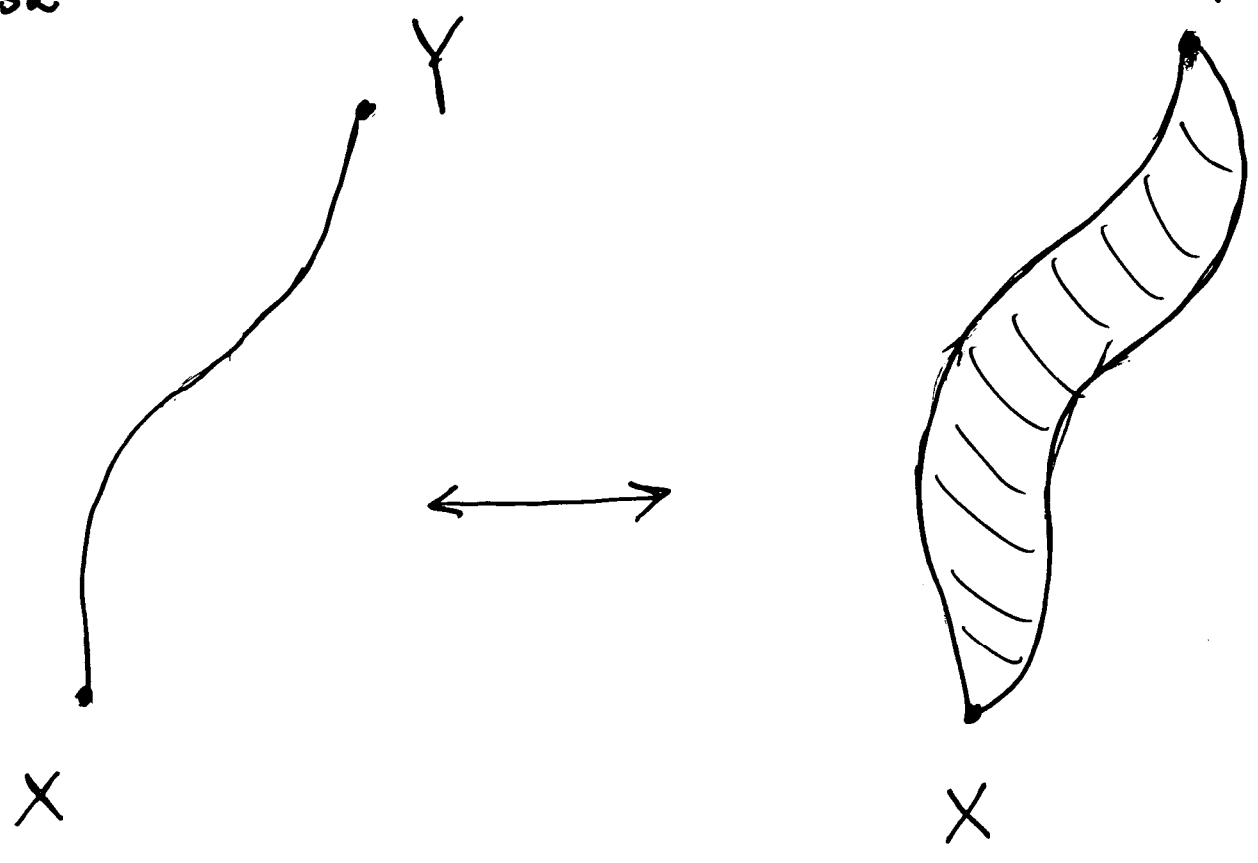
Vacuum structure and
phase transitions of gauge
theories with linear
action.

G. Savvidy

Nat. Res. Cent. Demokritos

Athens

G.S. & K.S.
Phys. Lett
1982



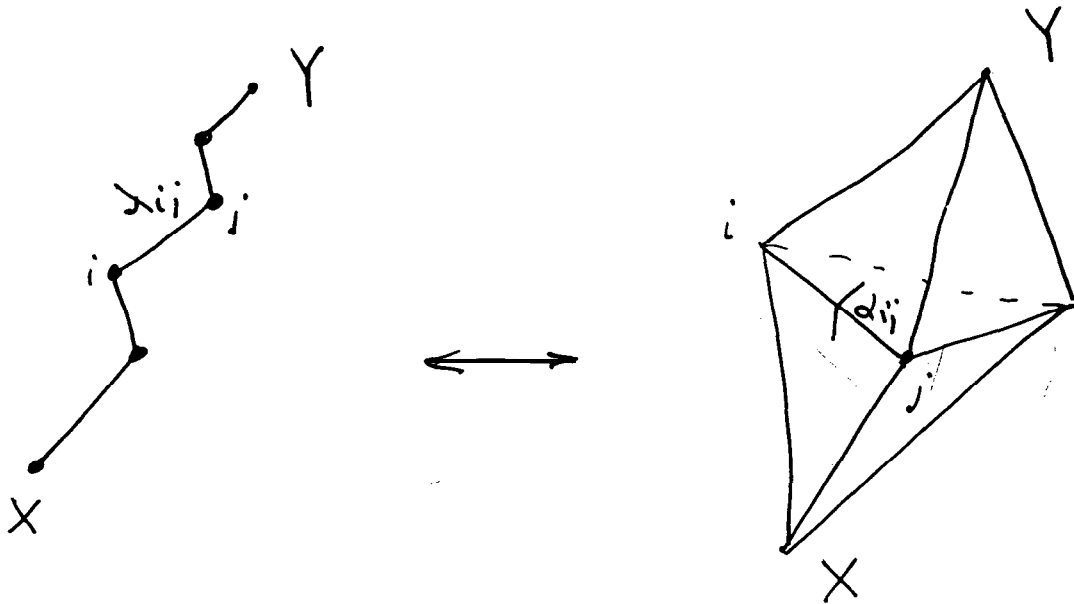
$$K(X, Y) = \sum_{\{\text{paths}\}} e^{-A_{xy}} \longleftrightarrow$$

$$\longleftrightarrow \sum_{\{\text{surfaces}\}} e^{-A_{xy}}$$

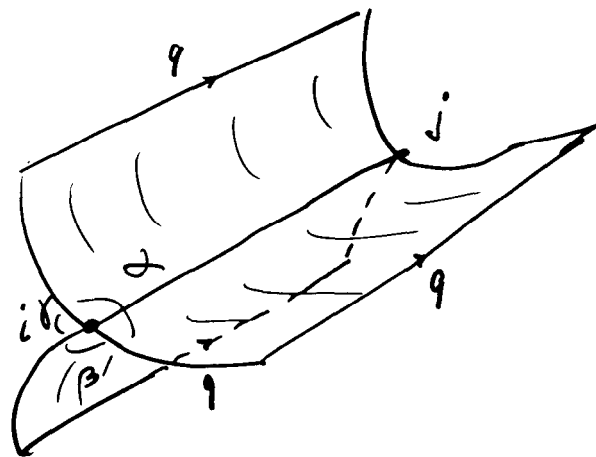
$$[A_{xy}] = \text{cm}$$

Gonihedric Action

$$A_{xy} = m \cdot \sum_{\langle ij \rangle} \lambda_{ij} \iff m \cdot \sum_{\langle ij \rangle} \lambda_{ij} \cdot (\pi - \alpha_{ij})^3 = A_{xy}$$



λ_{ij} - is the length of the edge $\langle ij \rangle$
 α_{ij} - is the dihedral angle

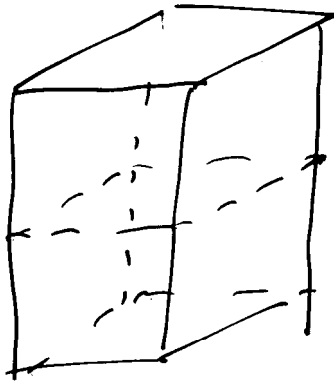
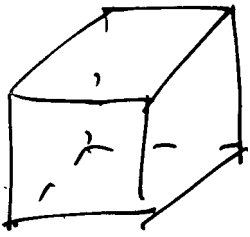


$$k \cdot m \sum_{\langle ij \rangle} \lambda_{ij} \cdot \left[|\pi - \alpha_{ij}|^3 + |\pi - \beta_{ij}|^3 + |\pi - \gamma_{ij}|^3 \right]$$

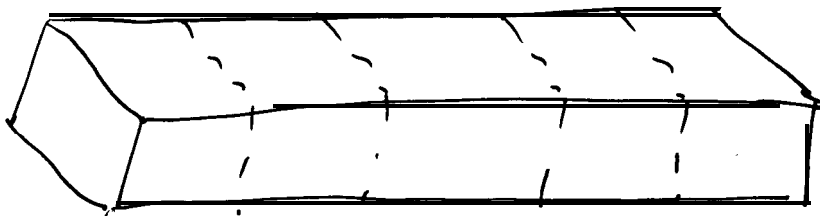
k : interaction constant

F. Weyner
G. Sarrily
1994

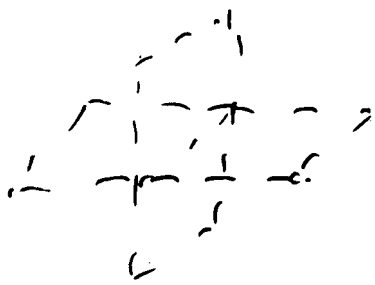
3 -



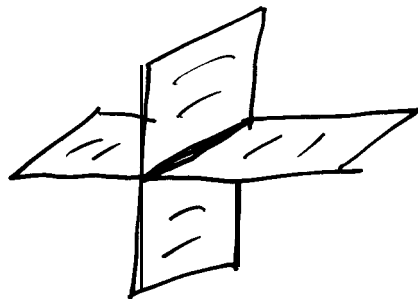
$$A = 12a \cdot \theta\left(\frac{\sqrt{2}}{2}\right) ; A = 16a \theta\left(\frac{\sqrt{2}}{2}\right).$$



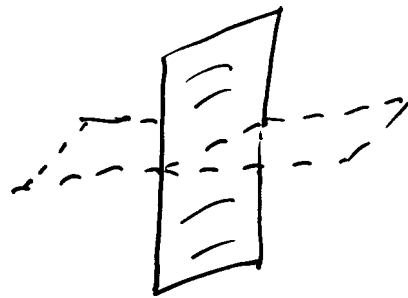
$$A \approx L \cdot a \cdot \theta\left(\frac{\sqrt{2}}{2}\right) \approx L$$



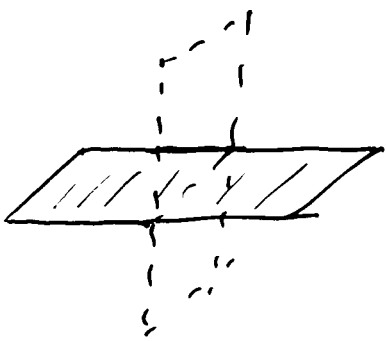
Ω_1



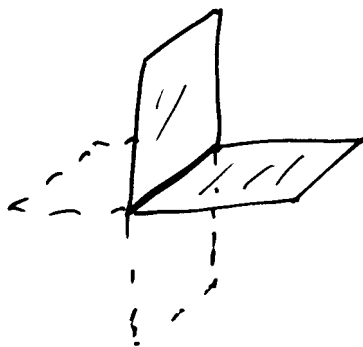
Ω_2



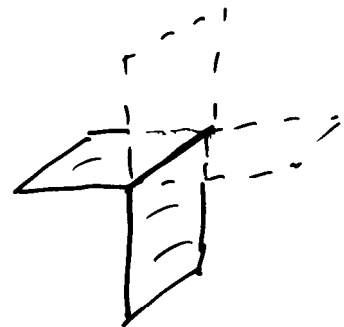
Ω_3



Ω_4



Ω_5

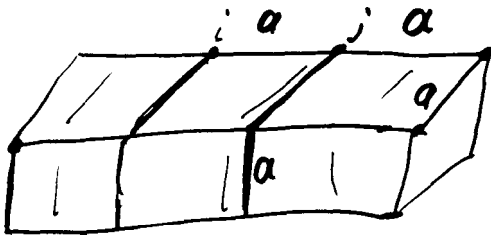


Ω_6

3. Strig on Jeu lattice

1994 F.J.W.

G.K.S.
K.G.S.



$$\theta(0) = \infty$$

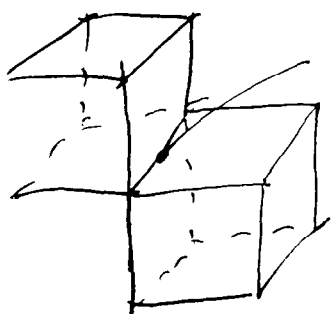
$$\theta\left(\frac{J}{2}\right) = J$$

$$\theta(J) = 0$$

$$A = \sum_{\langle ij \rangle} \lambda_{ij} \cdot \theta(\alpha_{ij})$$

$$\lambda_{ij} = a$$

$$= \sum_{\langle ij \rangle} a \cdot \theta(\alpha_{ij})$$



line of intersection

$$4 \cdot a \cdot \theta\left(\frac{J}{2}\right) \cdot k$$

Nambu

-5-

Linear

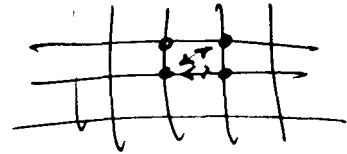
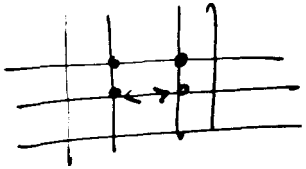
$$D = 3$$

$$\underline{k = 1}$$

$$D = 3$$

$$H = - \sum_{\langle ij \rangle} \sigma_i \sigma_j$$

$$H = - \sum_{\langle ij \rangle} \sigma_i \sigma_j + \frac{1}{4} \sum_{\langle ij \mu \rangle} \sigma_i \sigma_{j+\mu}$$



$$\underline{k \neq 1}$$

$$H = -k \sum_{\langle ij \rangle} \sigma_i \sigma_j + \frac{k}{4} \sum_{\langle ij \mu \rangle} \sigma_i \sigma_{j+\mu} + \frac{k-1}{4} \sum_{\langle ij \mu \nu \rangle} \sigma_i \sigma_j \sigma_{i+\mu} \sigma_{j+\nu}$$

$$Z(\beta) = \sum_{\{\sigma\}} e^{-\beta H} = \sum_{\{u, d\}} e^{-\beta A}$$

$$\underline{k = 0}$$

$$H = - \frac{1}{4} \sum_{\langle ij \mu \nu \rangle} \sigma_i \sigma_j \sigma_{i+\mu} \sigma_{j+\nu} \xleftrightarrow{\text{dual}} H = - \sum_{\langle ij \rangle} u_i u_j + d_i d_j + s_i s_j$$



$\mathcal{D}=4$ Observables total energy

$$E_{\text{total}} = \frac{1-5k}{4} \sum_{\{\square\}} + \frac{k}{16} \sum_{\{\square\}} + \frac{k-1}{32} \sum_{\{\square\}}$$

 Simple edges

$$n_2 = \sum_{\{\square\}} - 2 \sum_{\{\square\}} - \sum_{\{\square\}}$$

Intersection energy

$$4k \cdot \bar{n}_4 + 6k \cdot \bar{n}_4 + 12k \cdot n_6 = E_{\text{tot}} - n_2$$

Area

$$s = \sum_{\{\square\}} (1 - \text{diagram})$$

There is no transfer matrix representation in $\mathcal{D}=4$.

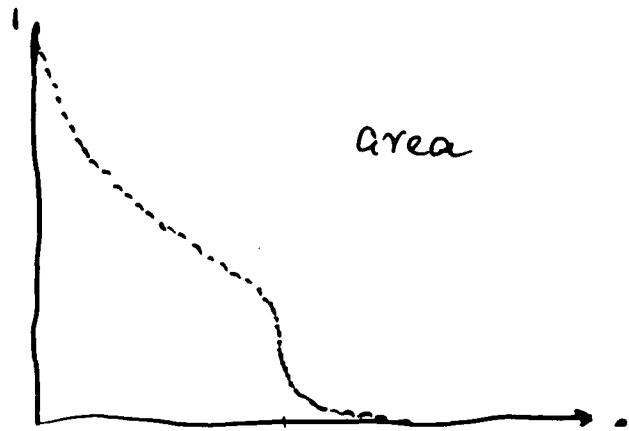
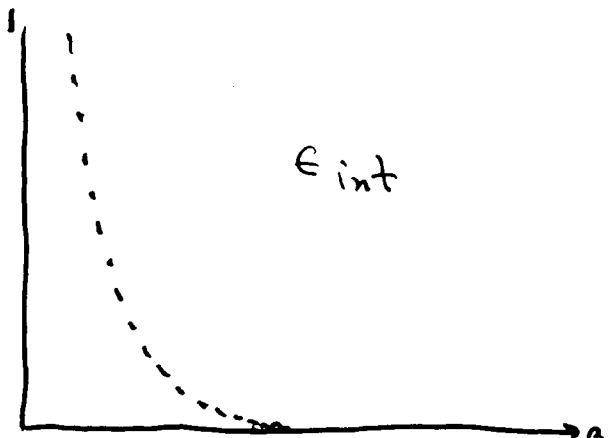
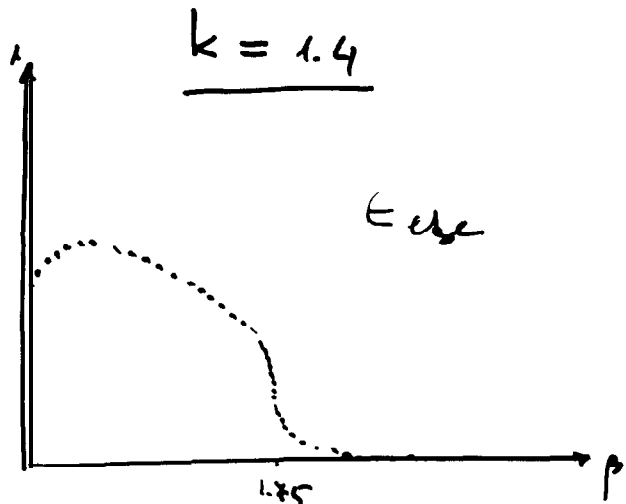
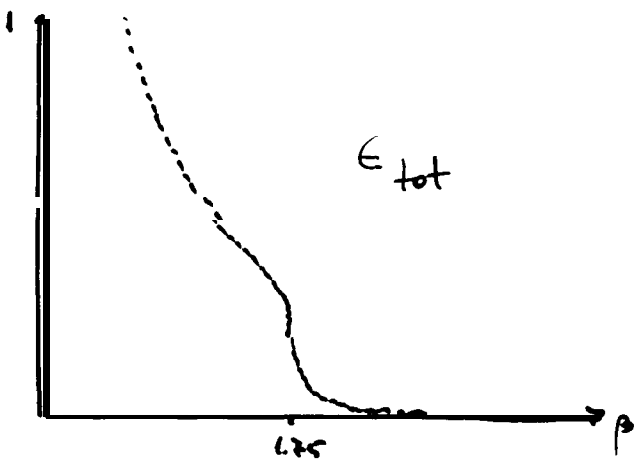
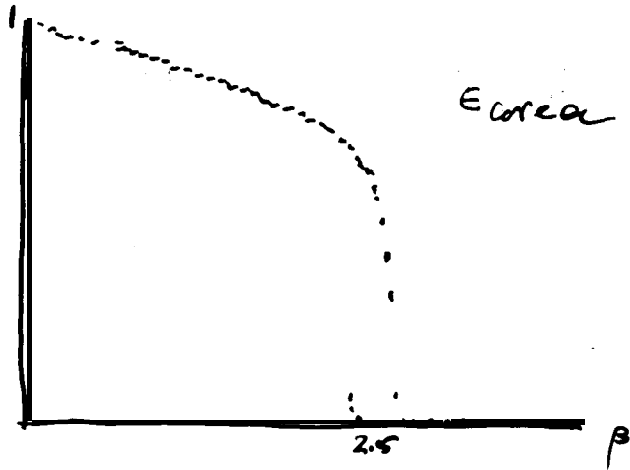
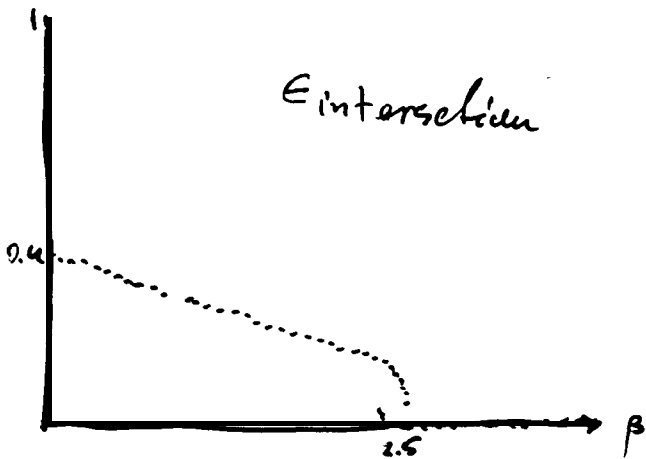
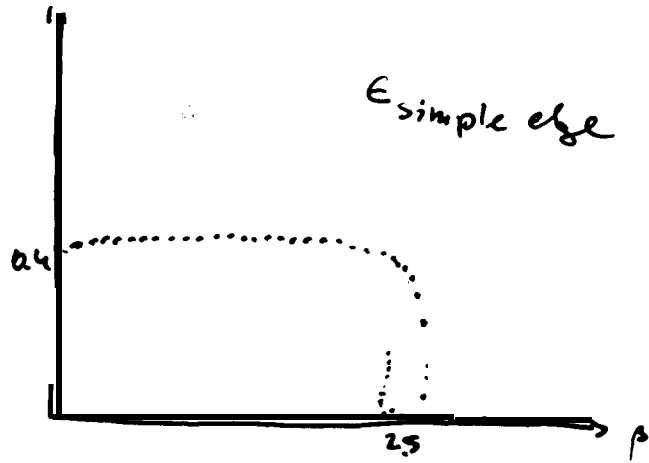
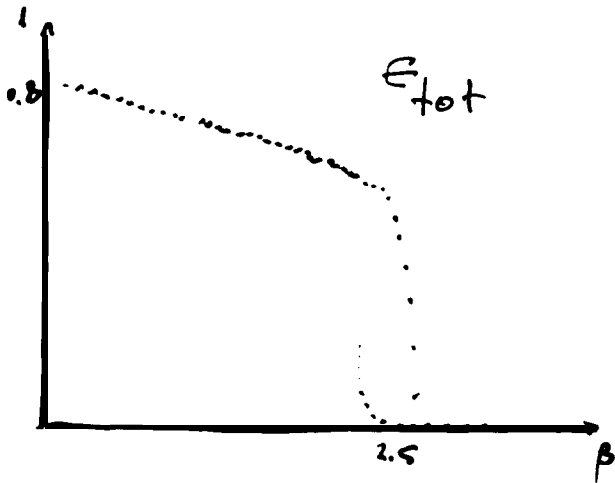
1979
Greutz
Jacobs
Rebbi

4 ⊕

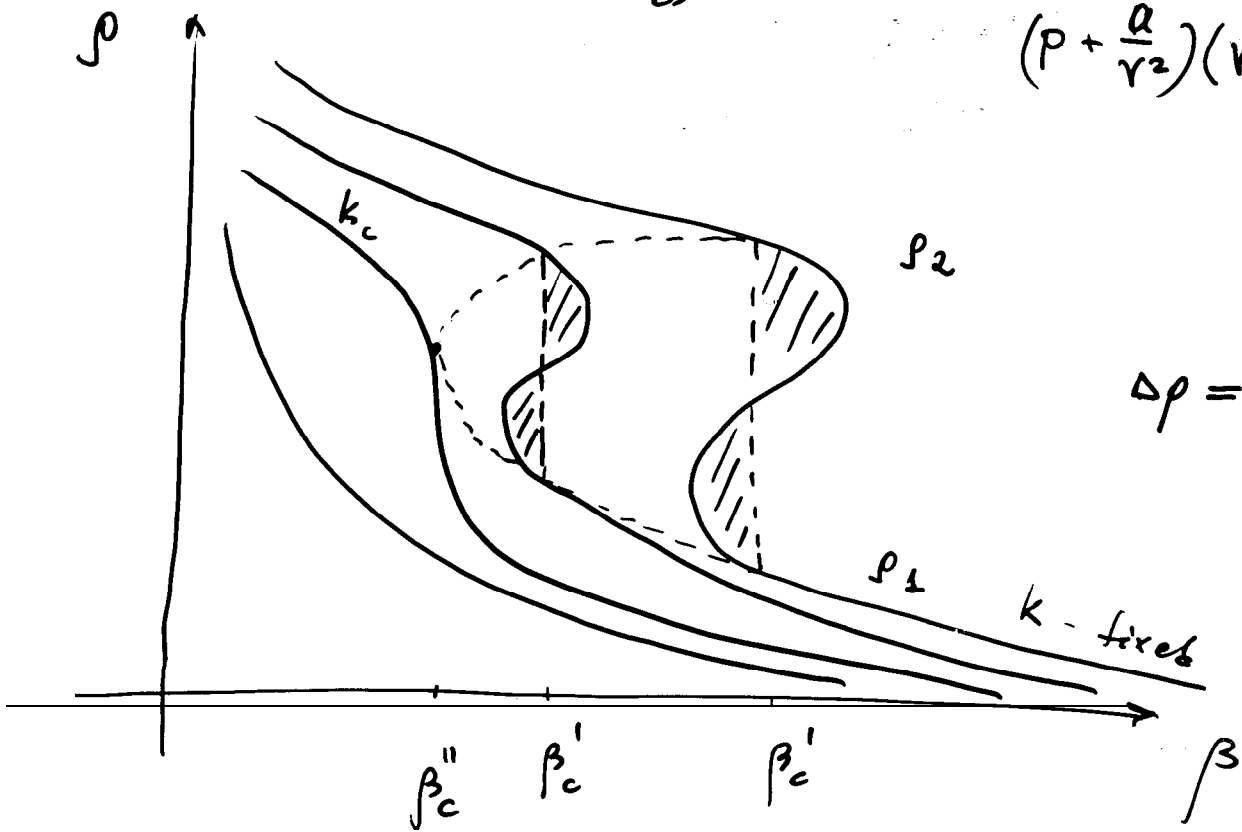
Simulation

$k=0.3$

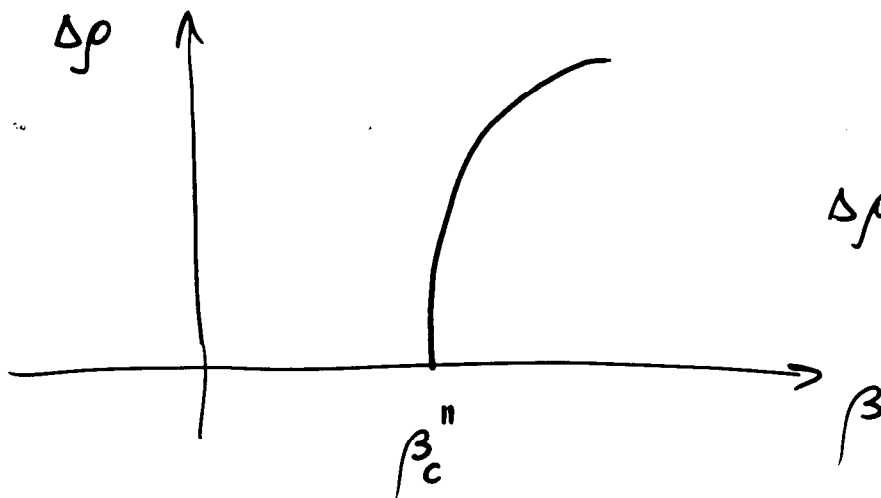
G. Katsoumbas
G.K. & K.G.



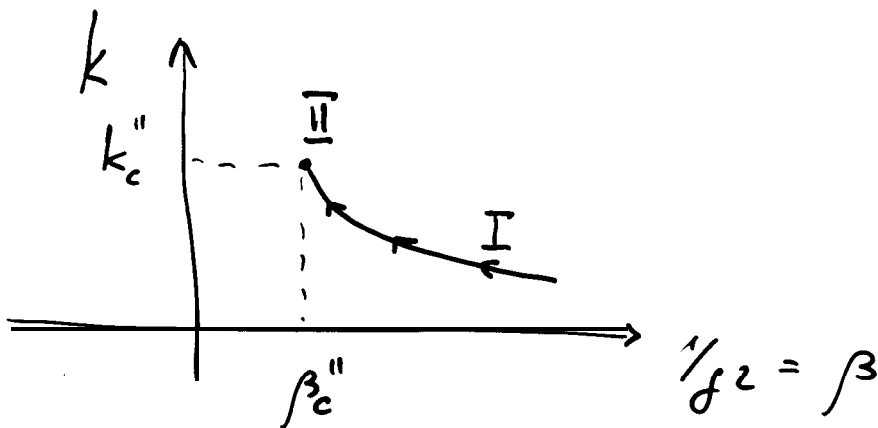
$$\left(P + \frac{a}{V^2}\right)(V - b) = RT$$



$$\Delta p = p_2 - p_1$$



$$\Delta p = (p_c' - p_c'')^\beta$$



$$\beta_{1\text{sig}}(2\sigma) = 0.12$$

$$\beta_{1\text{sig}}(3\sigma) = 0.33$$

$$\beta_{\text{HF}}(4\sigma) = 0.5$$

$$\beta(4\sigma) = 1.95$$

$$H = -\frac{1}{g^2} \sum \left[\text{Diagram of a plaquette with arrows} \right] \rightarrow -9-$$

$$H = \frac{n}{g^2} \sum \left\{ 1 - \frac{1}{n^2} \text{Re Tr } U_{\text{plaq}} \cdot \text{Re Tr } U_{\text{plaq}} \right\}$$

for smooth classical fields

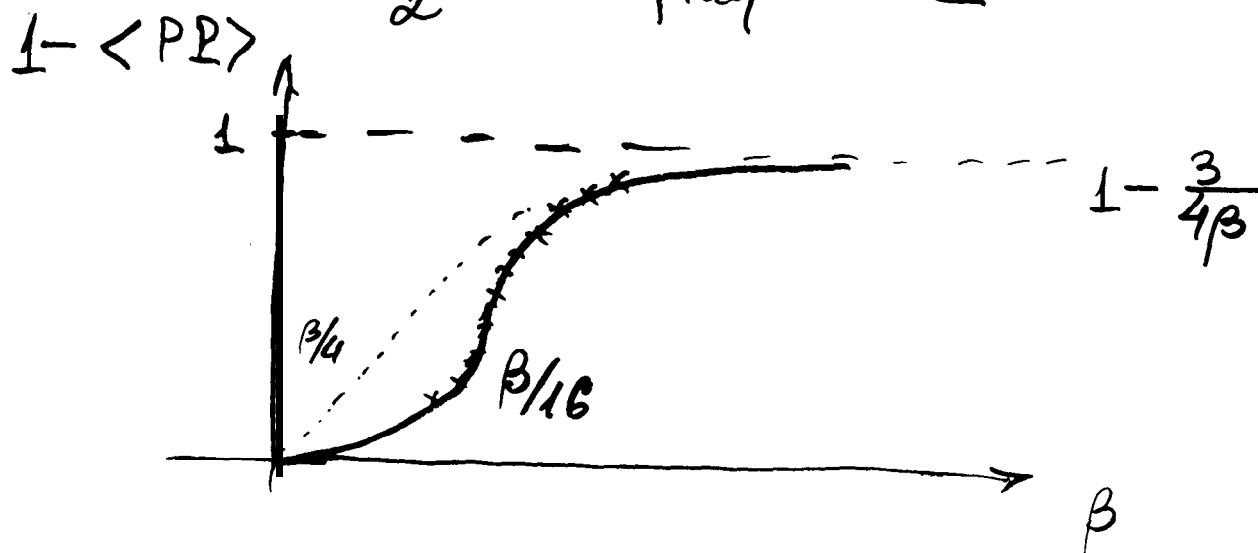
$$\rightarrow \frac{g^2 a^4}{n} \text{Tr } F^2$$

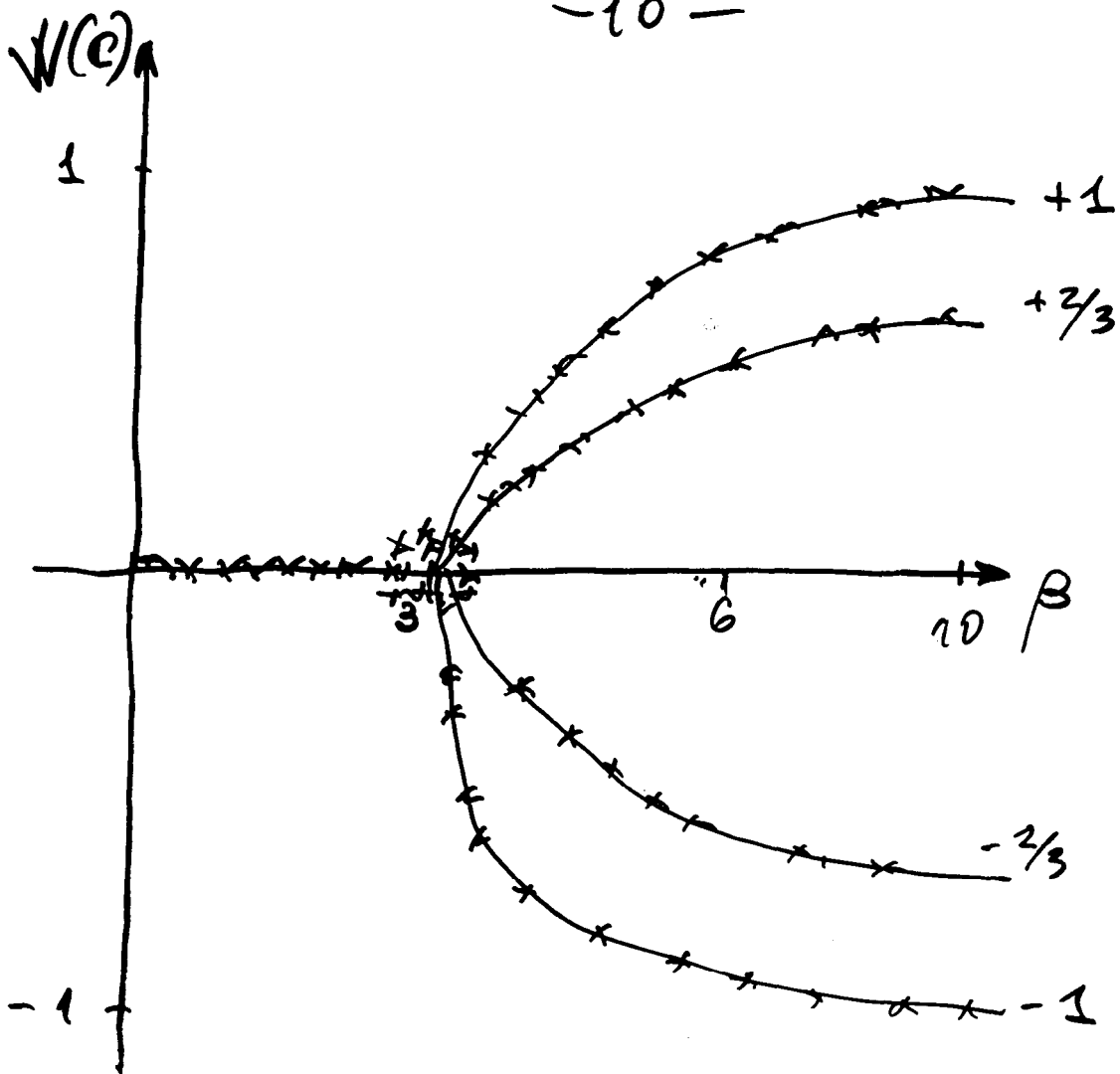
In addition to vacuum configurations

$$\frac{1}{2} \text{Tr } U_{\text{plaq}} = +1$$

there are 2^{3N} different vacuum configurations

$$\frac{1}{2} \text{Tr } U_{\text{plaq}} = -1$$





$$-\sigma L^2 - mL - \text{Const}$$

$$W(L, L) \approx \text{Const} \cdot L$$

$$\sigma_1^{\text{phys}} \approx \sigma_{1/3}^{\text{phys}} \approx \sigma_{-1/3}^{\text{phys}} \approx \sigma_{-1}^{\text{phys}}$$