

1. Publications

2. Monte Carlo Method

3. Metropolis algorithm

4. Many particle systems $\# 10^{30}$

5. Kolmogorov discrepancy $D_N \sim c\sqrt{N}$

6. Dynamical Systems $T^t x_0 \Rightarrow x_t$

7. Classification of Dynamical Systems:

K-systems, Entropy.

8. $D_N(T)$; $\tau_0 = 1/h(T)$

9. High Dimensional K-systems

10. Period of Generator on Galois field

$$\tau = \frac{p^D - 1}{p - 1}$$

11. Examples

$$p = 2^{61} - 1 \quad D = 256$$

$$\tau \approx 2^{15000}$$

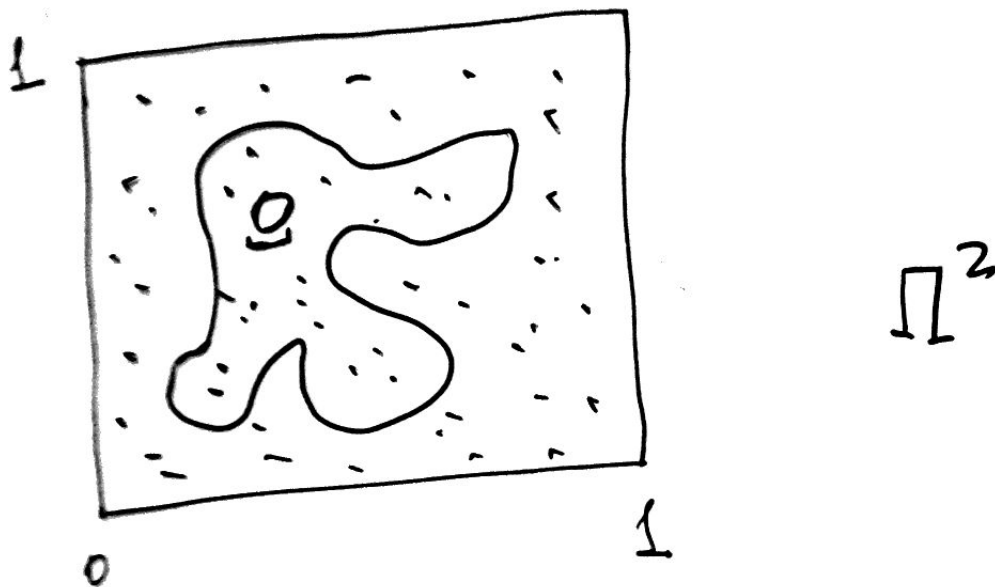
12.

Publications.

1. J. Comp. Phys. 97 (1991) 566; Preprint EPI-1986
2. J. Comp. Phys. 97 (1991) 573; Preprint EPI-1986.
3. Preprint EPI-1986 "Sinai Billiards as Pseudorandom number generators"
4. Int. J. Mod. Phys. C 7 (1996) 73 "K-system generators on Galois field"
5. F. James, Chaos, Solitons & Fractals 6 (1995) 221 "Chaos and Randomness"
6. M. Lüscher "A portable high-quality random number generator ..."
Comp. Phys. Comm. 79 (1994) 100
7. F. James "RANLUX: A Fortran implementation of RNG" Comp. Phys. Comm 1994

In Yerevan: collaboration was with
N. Akopyan - at DESY.mos.

Monte Carlo method



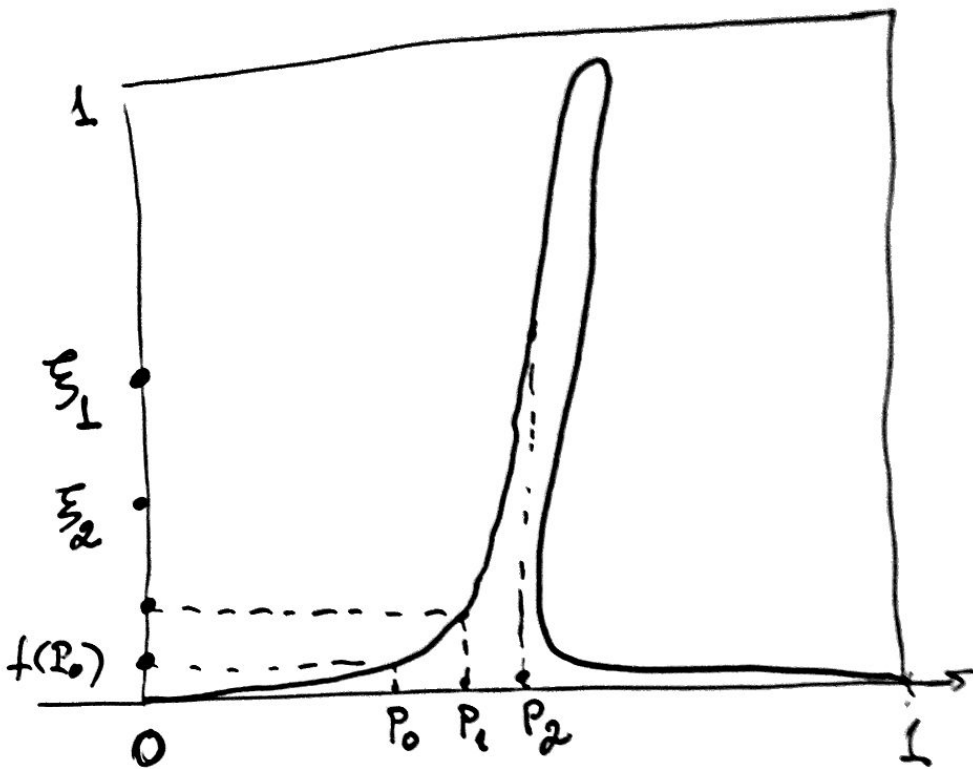
The area of O ?

In probability theory :

$$\xi \in \Pi^2 \quad p(\xi) = 1$$

$$\text{Area } O = \frac{\# \text{ inside } O}{N}$$

Metropolis algorithm.



The area is small

$$p \in [0, 1] \quad \xi \in [0, 1]$$

take p_1 and ξ_1 and compare

$$f(p_1) \quad \xi_1$$

if $f(p_1) < \xi_1$ then stay at p_0 and generate p_2 ξ_2

if $f(p_2) > \xi_2$ then jump to p_2 and so on



$$\rho(P) = f(P)$$

$$\int_{\Pi} f(P) dP = \frac{1}{N} \sum_{i=0}^{N-1} P_i$$

$$\int_{\Pi} f(P) \cdot g(P) dP = \frac{1}{N} \sum_{i=0}^{N-1} g(P_i)$$

$$Z(\beta) = \int e^{-\beta V(x_1, \dots, x_N)} dx_1 \dots dx_N$$

$$D=3 \quad N = 10^{30}$$

$$\rho \equiv e^{-\beta V(x_1, \dots, x_N)}$$

3D Ising

1979

1985

Parisi

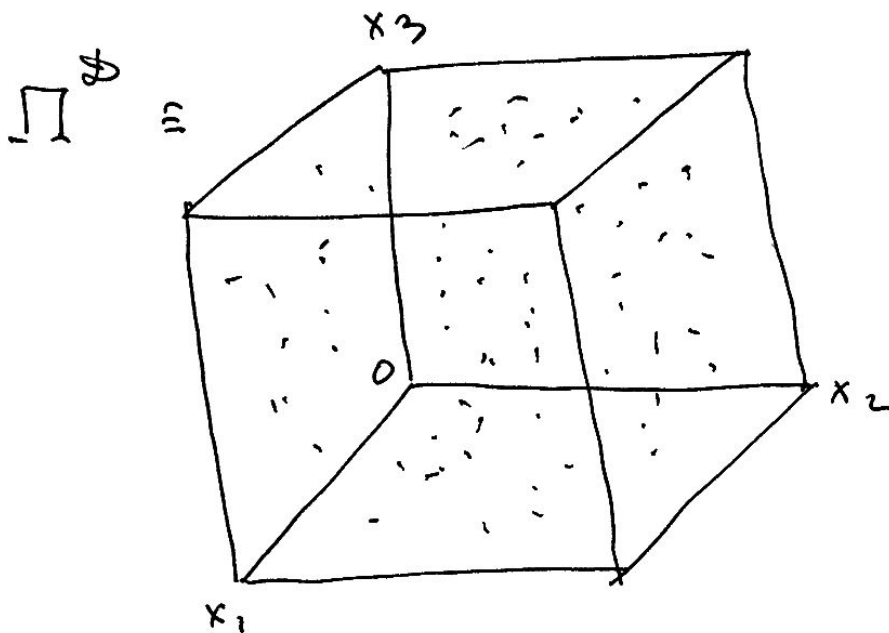
Miyake

K-system generator of Pseudorandom Numbers.

$$P_0, P_1, P_2, \dots, P_N$$

$$P = (x_1, x_2, \dots, x_g)$$

$$0 \leq x_i \leq 1.$$



Quality of pseudorandom numbers?

Kolmogorov discrepancy D_N

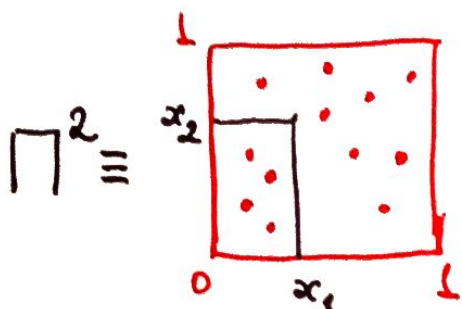
$$D_N(P_0, \dots, P_{N-1}) = \sup_{\{P\}} |N \cdot x_1 \dots x_D - A|$$

A - is number of points P_i with coordinates

$$0 \leq x_1^{(i)} \leq x_1; \dots; 0 \leq x_D^{(i)} \leq x_D$$

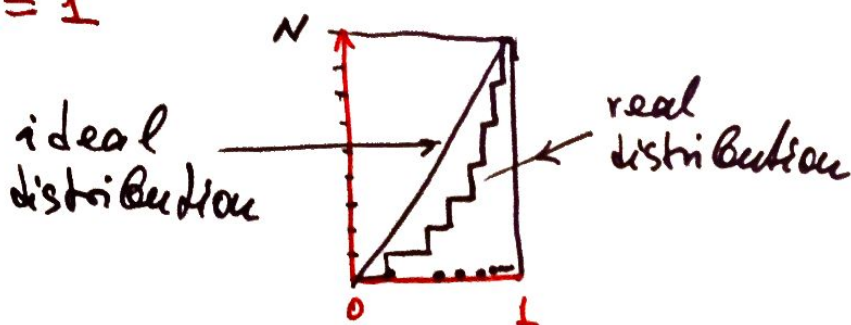
$N x_1 \dots x_D$ - is the number of points for ideal uniform distribution.

$D=2$



$N=12$
 $A=4$

$D=1$



Theorem.

$$\left| \frac{1}{N} \sum_{k=0}^{N-1} f(P_k) - \int_{\Pi^D} f(P) dP \right| \leq \text{const.} \cdot \frac{D_N}{N}$$

D_N - estimates a maximal deviation of real distribution of points from ideal one. $D_N \leq N$.

One should generate sequence P_k , so that D_N would grow as slowly as possible.

Dynamical origin of P_k

a) For random quantity ξ , $\rho(\xi) = 1$, then by central limiting theorem

$$D_N(\xi) \approx \sqrt{N}$$

so

$$\left| \frac{1}{N} \sum_{k=0}^{N-1} f(P_k) - \int f(P) dP \right| \leq \text{Const} \frac{1}{\sqrt{N}}$$

b) trajectory of dynamical system T

$$P_1 = TP_0, \quad P_2 = TP_1 = T^2 P_0, \dots, P_{N-1} = T^{N-1} P_0$$

and Π^D as the phase space of the dynamical system T , Liouville theorem should hold!

The rate of convergence is provided by the dynamical properties of a system T

$$\left| \frac{1}{N} \sum_{k=0}^{N-1} f(T^k P_0) - \int f(P) dP \right| \leq \text{Const} \cdot \frac{D_N(T)}{N}$$

How to get the best rate of convergence? $D_N(T)$

1. Area preserving map of the phase space M

$$M \xrightarrow{T^t} M$$

$$T^t x_0 = x_t$$

$$T^t A = A_t$$

in classical mechanics $x = \begin{pmatrix} q \\ p \end{pmatrix}$.

2. Classification of the map.

i) ergodic if

$$\lim_{t \rightarrow \infty} \frac{1}{t} \int dt \mu[T^t A \cap B] = \mu[A] \cdot \mu[B].$$

ii) mixing if

$$\lim_{t \rightarrow \infty} \mu[T^t A \cap B] = \mu[A] \cdot \mu[B]$$

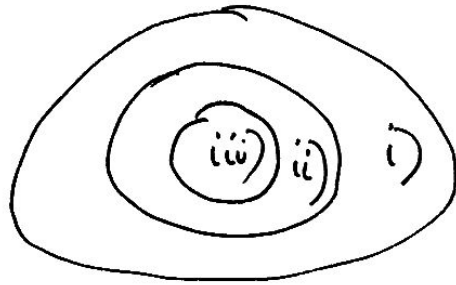
iii) n-fold mixing

$$\lim_{t_1, \dots, t_n \rightarrow \infty} \mu[A_0 \cap T^{t_1} A_1 \cap T^{t_2+t_1} A_2 \cap \dots \cap T^{t_n+\dots+t_1} A_n]$$

$$= \mu[A_1] \dots \mu[A_n].$$

There is hierarchy of systems

$$\text{iii) } \supset \text{ii) } \supset \text{i)}$$



iv) $n \rightarrow \infty$ mixing of any multiplicity

v) K-systems

Split-up $\Xi = \{C\} : \bigcup_{C \in \Xi} C = M, C \cap C' = \emptyset$

A) $\bigvee_{-\infty}^{+\infty} T^t \Xi = \epsilon$ split-up into separate points of M

B) $\bigwedge_{-\infty}^{+\infty} T^t \Xi = \gamma$ split-up coarsening of M

$$\text{v) } \supset \text{iv) } \supset \text{iii) } \supset \text{ii) } \supset \text{i)}$$

K-systems have mixing of any multiplicity n -

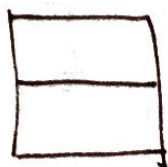
The best statistical property among all dynamical systems!

$h(T)$ - called the Kolmogorov entropy

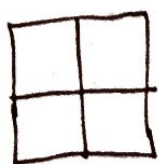
$$0 \leq h(T) \leq \infty$$

The entropy of $\Sigma = \{C\} \rightarrow H(\Sigma)$

$$H = - \sum_{C \in \Sigma} \mu[C] \ln \mu[C]$$



$$H = - \frac{1}{2} \ln \frac{1}{2} - \frac{1}{2} \ln \frac{1}{2} = - \ln \frac{1}{2} = \ln 2$$



$$H = - \frac{1}{4} \ln \frac{1}{4} - \frac{1}{4} \ln \frac{1}{4} \dots = - \ln \frac{1}{4} = 2 \ln 2$$



$$H = - \frac{1}{3} \ln \frac{1}{3} - \frac{2}{3} \ln \frac{2}{3} = - \ln \frac{1}{3} - \frac{2}{3} \ln 2 = \ln 3 - \frac{2}{3} \ln 2$$

if many splits $\Sigma_1, \dots, \Sigma_2, \dots$

$$H(\bigvee_{\alpha} \Sigma_{\alpha}) = - \sum \mu[\bigcap_{\alpha} C_{\alpha}] \ln \mu[\bigcap_{\alpha} C_{\alpha}]$$

$$T^{t=1} = T$$

$$\xi, T\xi, T^2\xi, \dots$$

$$H_n(T, \xi) = H(\xi \vee T\xi \vee T^2\xi \dots \vee T^{n-1}\xi)$$

$$h(T) = \sup_{\xi} \lim_{n \rightarrow \infty} \frac{1}{n} H(\xi \vee T\xi \vee \dots \vee T^{n-1}\xi)$$

$H_n(T, \xi)$ - quantity of information along the time $t = n$.

$h(T, \xi)$ - information per unit time.

Another definition $f \in L_2(M)$

$$U^t f(x) = f(T^t x)$$

U^t as one parameter unitary operator, all

$$|U_t| = 1.$$

i) ii) ... v) - are spectral properties.
countable-multiple Lebesgue

Mixing, n -fold mixing and K -systems

(property of relaxation) \rightarrow (to uniform distribution)

Relaxation of K -systems is most rapid because of their exponential instability.

(Slow growth of due discrepancy $D_N(T)$) \equiv (quick relaxation of dynamical system T)

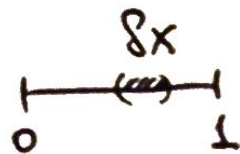
Exp. $P_k = r P_{k-1} \pmod{1}$

$$R_N = \frac{\langle (P^{L+N} - \langle P^{L+N} \rangle)(P^L - \langle P^L \rangle) \rangle}{\langle (P^L - \langle P^L \rangle)^2 \rangle} \sim e^{-N \ln r}$$

Scale of correlation splitting

$$\tau_0 = \frac{1}{\ln r}$$

(time of uniform fill)
of $[0,1]$ from δx) = $\tau_0 \ln(1/\delta x)$



Exp. Anosov K-systems.

$$P_K = T P_{K-1} \quad T = \|a_{K\ell}\|$$

a) $\det \|a_{K\ell}\| = 1$

b) $|\lambda_j| \neq 1$

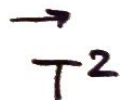
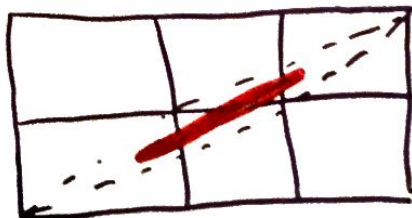
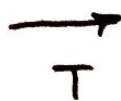
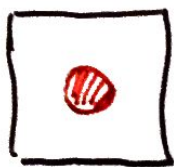
entropy $h(T) = \sum_{|\lambda_j| > 1} \ln |\lambda_j|$

correlation splitting time

$$\tau_0 = \frac{1}{h(T)} = \frac{1}{\sum_{|\lambda_j| > 1} \ln |\lambda_j|}$$

$d=2$

$$\begin{pmatrix} x_1 \\ x_2 \end{pmatrix}_K = \begin{pmatrix} 2 & 1 \\ 1 & 1 \end{pmatrix} \begin{pmatrix} x_1 \\ x_2 \end{pmatrix}_{K-1}$$



$$\lambda_{1,2} = \frac{3 \pm \sqrt{5}}{2}$$

$$h(T) = \ln \frac{3 + \sqrt{5}}{2}$$

High dimensional K-systems.

$$T_2 = \begin{vmatrix} 2 & 1 \\ 1 & 1 \end{vmatrix}$$

$$T_3 = \begin{vmatrix} 2 & 2 & 1 \\ 1 & 2 & 1 \\ 1 & 1 & 1 \end{vmatrix}$$

$$T_4 = \begin{vmatrix} 2 & 3 & 4 & 1 \\ 1 & 2 & 2 & 1 \\ 1 & 1 & 2 & 1 \\ 1 & 1 & 1 & 1 \end{vmatrix}$$

$$\dots T_d = \begin{vmatrix} 2 & 3 & 4 & \dots & d & 1 \\ 1 & 2 & 3 & \dots & d-1 & 1 \\ 1 & 1 & 2 & \dots & d-2 & 1 \\ \dots & \dots & \dots & \dots & \dots & \dots \\ 1 & 1 & 1 & \dots & 2 & 1 \\ 1 & 1 & 1 & \dots & 1 & 1 \end{vmatrix}$$

$$T_{170}$$

$$\lambda_{\max} = 1539.9$$

$$h(T) = 108.9$$

N	1	2	...	10×10^5
D_N/\sqrt{N}	1.714	-	-	3,302
D_N/\sqrt{N}	1.233	-	-	1,102


The entropy $h(\tau)$ defines the number $\mathcal{N}(\tau)$ of periodic trajectories with period less or equal τ

$$\mathcal{N}(\tau) \rightarrow \frac{e^{h \cdot \tau}}{h \cdot \tau}$$

in T_{170}

$$\mathcal{N}(\tau) \rightarrow \frac{e^{109 \cdot \tau}}{109 \cdot \tau} !$$

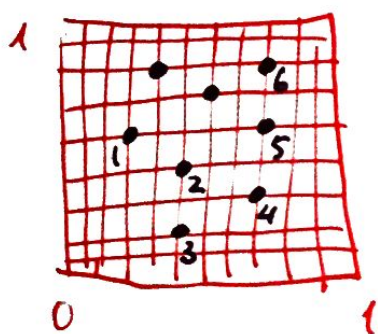
Only periodic trajectories can be simulated on a computer.



The periodic trajectories of K-system.

$$\text{If } P_0 = \left(\frac{q_1}{p_1}; \frac{q_2}{p_2}; \dots; \frac{q_D}{p_D} \right) \quad T = \|a_{ik}\|$$

$\|a_{ik}\|$ are integer and q_i, p_i also
then $P_0 = T^k P_0$ is periodic trajectory
with period k .



$$0 \leq q_i \leq p$$

rational sublattice

$$\text{If } P_0 = \left(\frac{q_1}{p}; \frac{q_2}{p}; \dots; \frac{q_D}{p} \right)$$

then trajectory will stay
on the same sublattice.

The period τ_p - on sublattice $\leq p^D$

$$\frac{q_i^{(k+1)}}{p} = \sum_j a_{ij} \frac{q_j^{(k)}}{p} \quad \text{mod } 1$$

is equivalent to

$$q_i^{(k+1)} = \sum_j a_{ij} q_j^{(k)} \quad \text{mod } p$$

If p - is prime number then

g_i - belongs to Galois field $GF(p)$
 $\{0, 1, \dots, p-1\}$

$GF[3] \quad \{0, 1, 2\} \quad g=2 \quad p=3$
 $g=2 \quad g^2=4=1$

$GF[5] \quad \{0, 1, 2, 3, 4\} \quad g=3 \quad p=5$
 $g=3 \quad g^2=9=4 \quad g^3=2 \quad g^4=1$

g - is a primitive element of $GF(p)$
 $g^{p-1} = 1 \pmod{p}$.

i) If eigenvalue λ of $\|a_{ik}\|$ coincides with primitive element g then maximal period $\tau_p = p-1$.

ii) If eigenvalue λ coincides with primitive element of quadratic extension $GF[\sqrt{p}]$, then maximal period $\tau_p = p^2-1$

$$GF[\sqrt{3}]$$

$$g=2$$

$$h=\sqrt{2}$$

$$a+b \cdot h$$

$$a, b \in GF[3]$$

$$w = 1 + \sqrt{2}$$

$$w^2 = 2\sqrt{2}$$

$$w^3 = 1 + 2\sqrt{2}$$

$$w^4 = 2$$

$$w^5 = 2 + 2\sqrt{2}$$

$$w^6 = \sqrt{2}$$

$$w^7 = 2 + 2\sqrt{2}$$

$$w^8 = 1$$

$$\tau_3 = 3^2 - 1 = 9$$

$$x^2 - 2 = 0$$

$$x = \sqrt{2}$$

iii) If \mathcal{L} coincide with \mathcal{D} -dim.
extension of Galois field

$$GF[\sqrt[p]{p}] \text{ then } \tau_p = p^{\mathcal{D}} - 1$$

the elements of $GF[\sqrt[p]{p}]$ have the
form

$$a + bh + \dots + eh^{\mathcal{D}-1}$$

$$a, b, \dots, e, \in GF[p]$$

h - is primitive element of
 $GF[\sqrt[p]{p}]$.

The period is:

$$\tau = \frac{p^N - 1}{p - 1}$$

If we use the largest Mersenne number

$$p = 2^{61} - 1$$

and dimension of the generator $N = 256$

we can get:

$$\tau \approx 2^{61 \cdot 256} = 2^{15552}$$

In 2013 the largest known prime number is:

$$p = 2^{57885161} - 1$$

by "Great Internet Mersenne Prime Search"