



NATIONAL CENTRE FOR
SCIENTIFIC RESEARCH "DEMOKRITOS"
INSTITUTE OF NUCLEAR AND PARTICLE PHYSICS



H.F.R.I.
Hellenic Foundation for
Research & Innovation

Top quark pair production and decay with jet activity at the LHC

Giuseppe Bevilacqua
NCSR "Demokritos"

Theory Seminar

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Work in collaboration with H.Y.Bi, B.Hartanto, M.Kraus, M. Lupattelli, D. Stremmer and M.Worek

Based on: [JHEP 08 \(2021\) 008](#) [Phys.Rev.D 107 \(2023\) 11, 014028](#) [Phys.Rev.D 107 \(2023\) 11, 114027](#)

Fundamentals of top quark physics

Why is top quark physics interesting?

- The top quark is the heaviest particle of the Standard Model

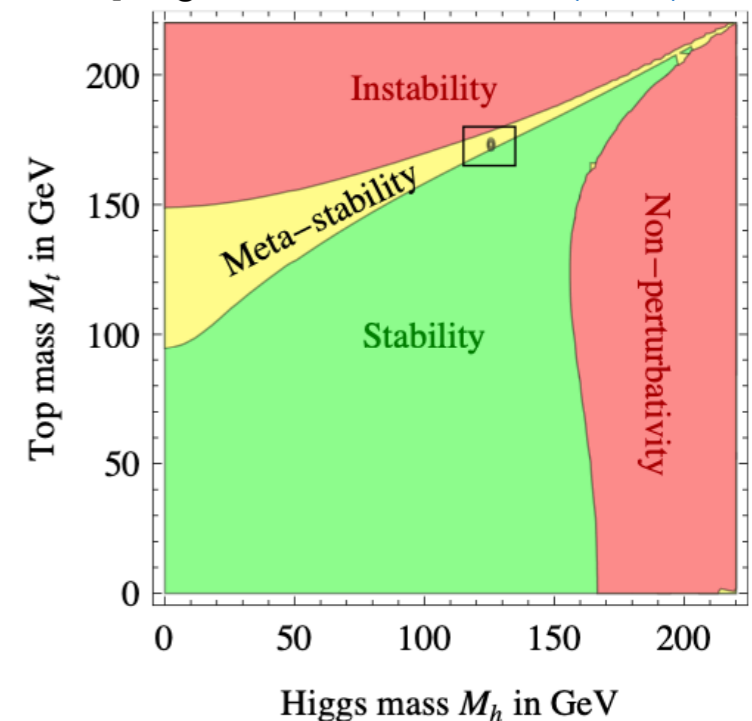
- ↳ • Substantial Yukawa coupling: sensitive to EWSB mechanism
- Top mass is a free parameter: need precise measurements

- Extremely short-lived $\sim 5 \cdot 10^{-25} \text{ s} \ll \Lambda_{QCD}^{-1}$

- ↳ • Decays before bound states can be formed
- $\text{BR}(t \rightarrow Wb) \approx 100\%$: unique experimental signature
- Direct handle of top quark properties from decay products

$$y_t = \sqrt{2} \frac{m_t}{v} \approx 1$$

[Degrassi et al., [JHEP 08 \(2012\) 098](#)]



Production

- Inclusive rates
- Differential distributions
- New production mechanisms
- ...

Decays

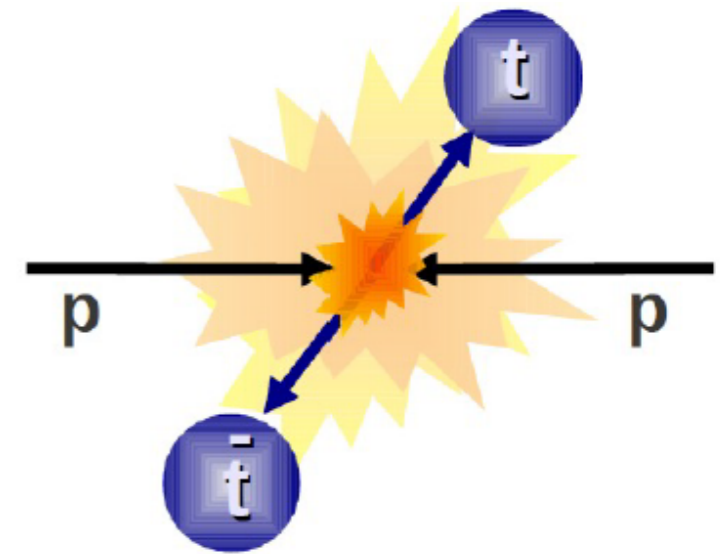
- Decay channels (SM & BSM)
- Couplings to SM particles
- Spin correlations
- ...

Intrinsic properties

- Mass
- Charge
- Lifetime
- Width
- ...

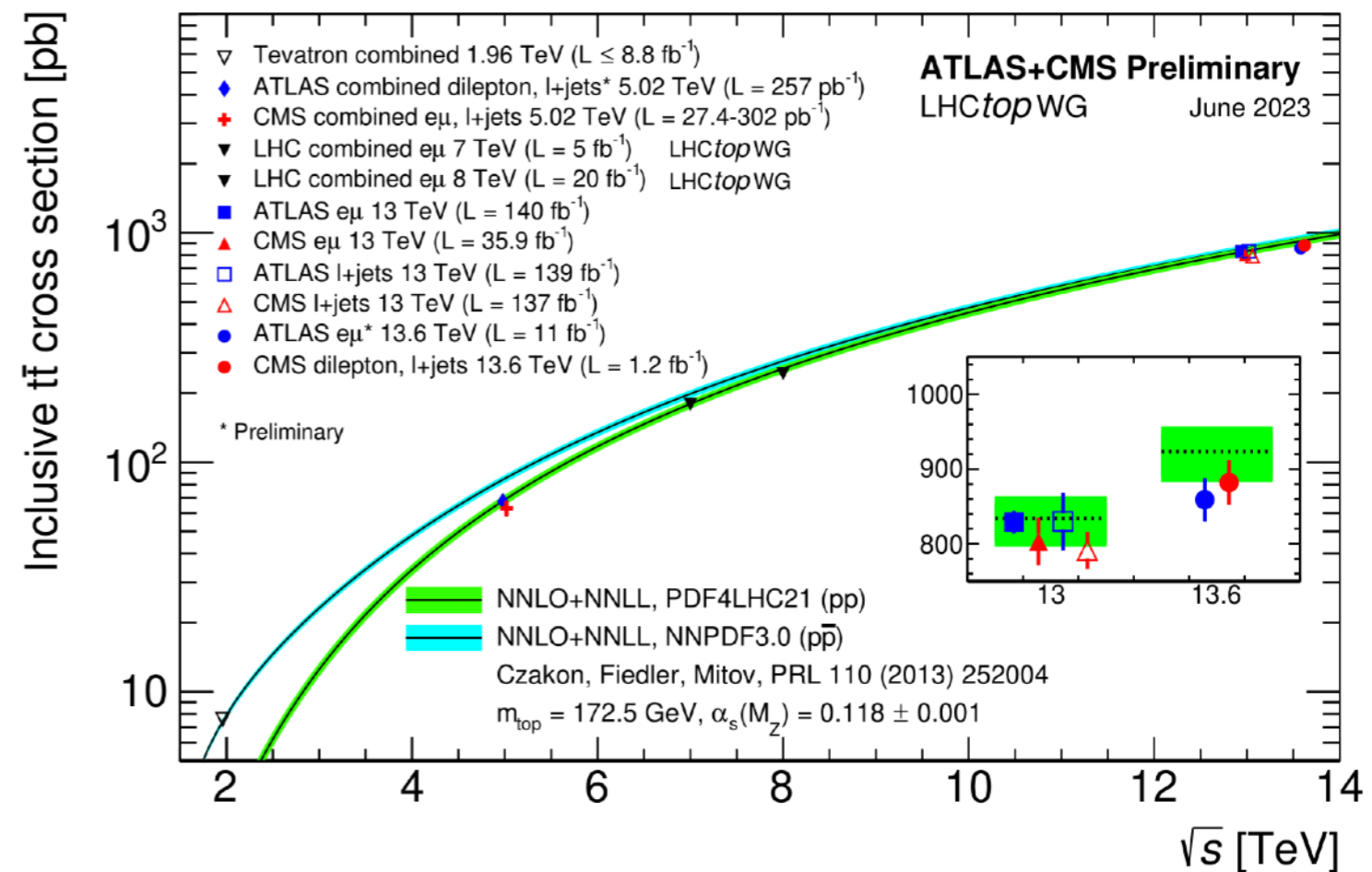
$t\bar{t}$ production at the LHC

- Dominant mode of top quark production at the LHC
- Can be selected with high purity (dilepton, but also ℓ +jets)
- $\mathcal{O}(100 \text{ M})$ $t\bar{t}$ events in Run 2
- “Standard candle” for precision tests of SM



[[LHCPhysics/TtbarNNLO](https://lhcb.physics.stonybrook.edu/TtbarNNLO)]

7 TeV			
Central value	Scale uncert.	PDF+ α_s uncert.	Mass uncert. up / down
179.6	+4.8 -6.2	± 6.1	-5.26 +5.44
8 TeV			
Central value	Scale uncert.	PDF+ α_s uncert.	Mass uncert. up / down
256.0	+6.7 -8.9	± 8.0	-7.33 +7.58
13 TeV			
Central value	Scale uncert.	PDF+ α_s uncert.	Mass uncert. up / down
833.9	+20.5 -30.0	± 21.0	-22.5 +23.2
13.6 TeV			
Central value	Scale uncert.	PDF+ α_s uncert.	Mass uncert. up / down
923.6	+22.6 -33.4	± 22.8	-24.6 +25.4

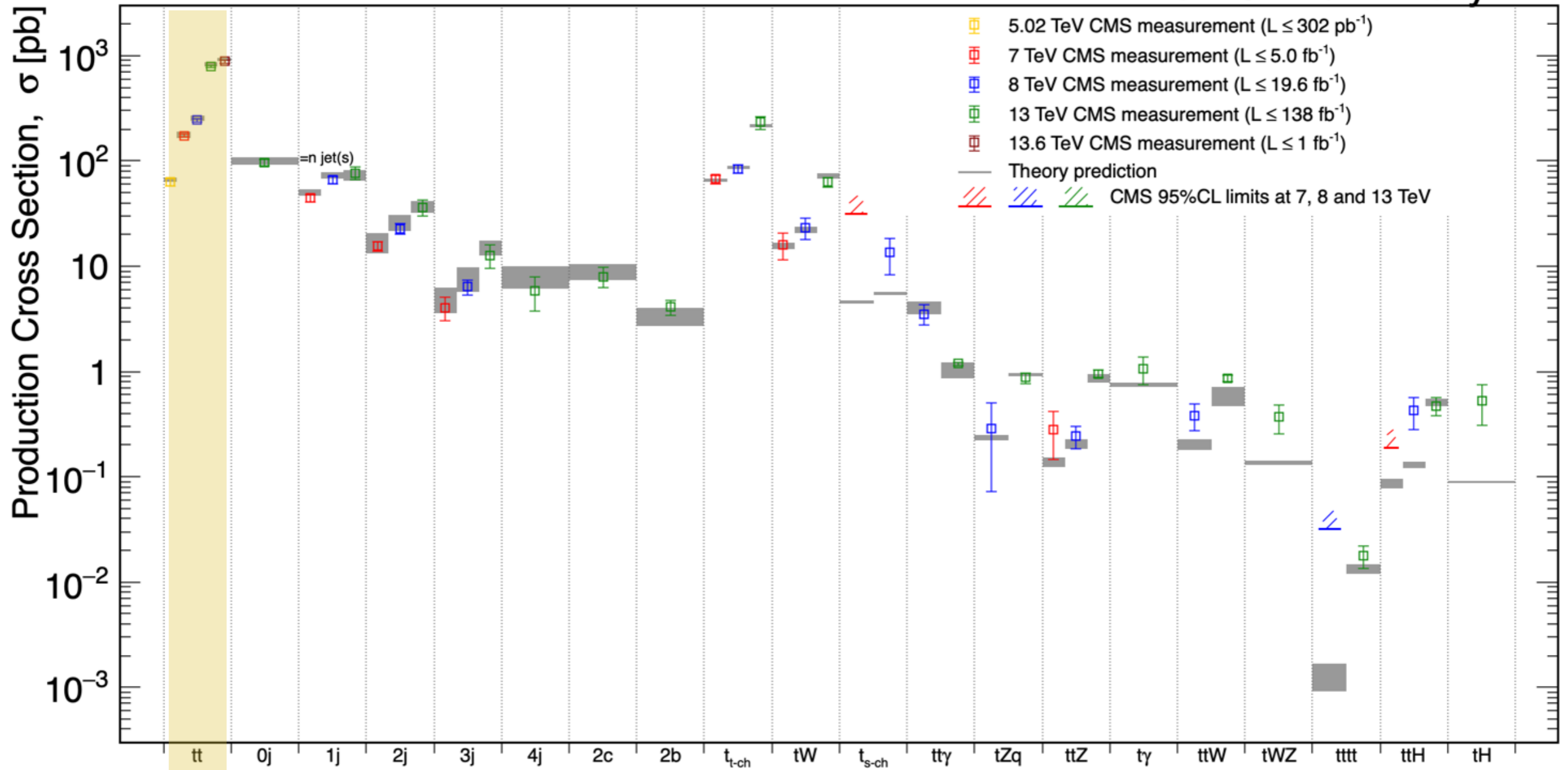


Top quark production cross sections

[CMSPublic/PhysicsResultsCombined]

Aug 2023

CMS Preliminary



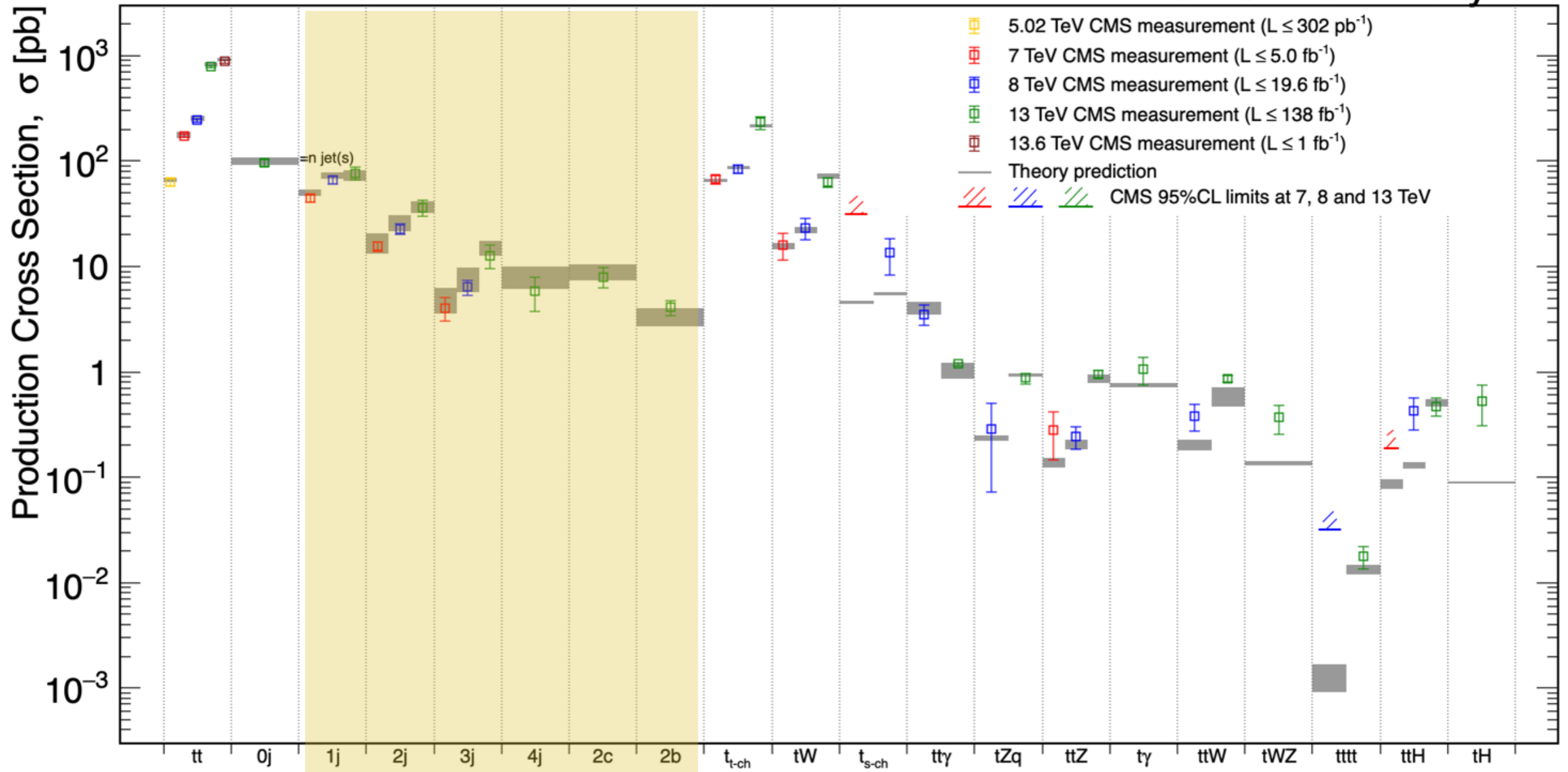
Inclusive $t\bar{t}$ sample

Top quark production cross sections

[CMSPublic/PhysicsResultsCombined]

Aug 2023

CMS Preliminary



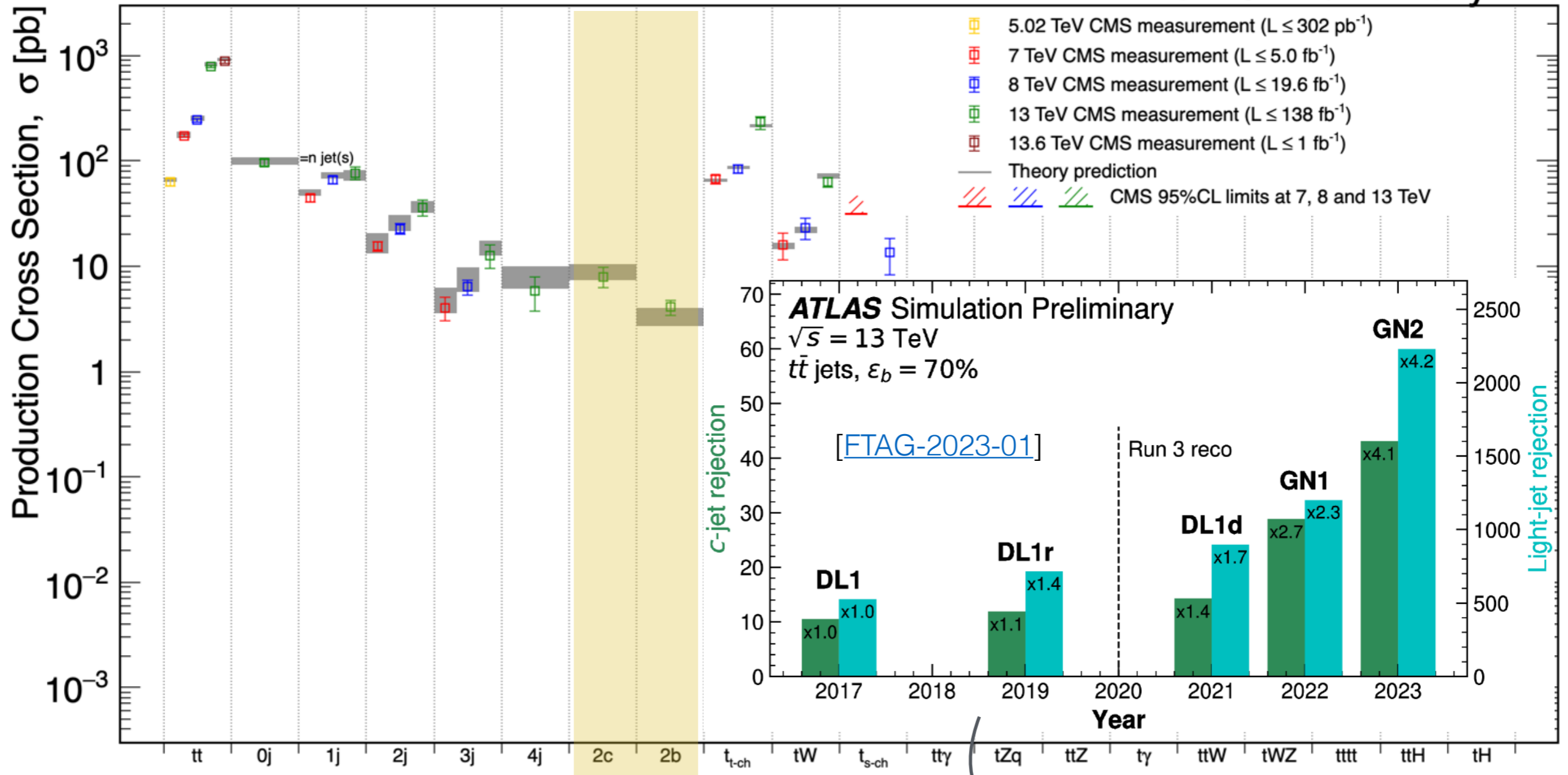
Hadronic jet activity in $t\bar{t}$ events

Top quark production cross sections

[CMSPublic/PhysicsResultsCombined]

Aug 2023

CMS Preliminary

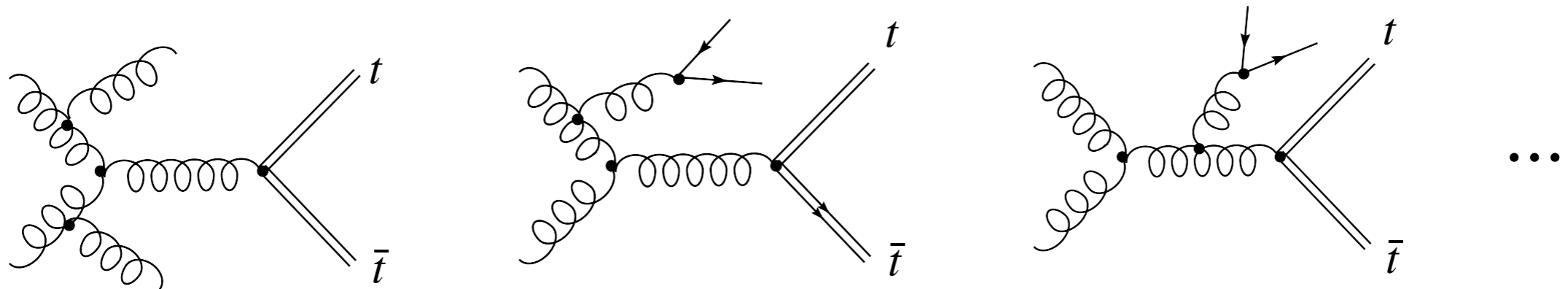


$t\bar{t} + \text{HF}$

b- and *c*-tagging algorithms are constantly improving

$t\bar{t}$ + jets as a signal

- A significant fraction of the inclusive $t\bar{t}$ sample is accompanied by **additional hard jet(s)** arising from QCD radiation



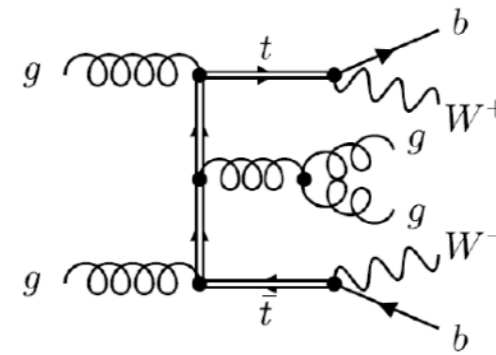
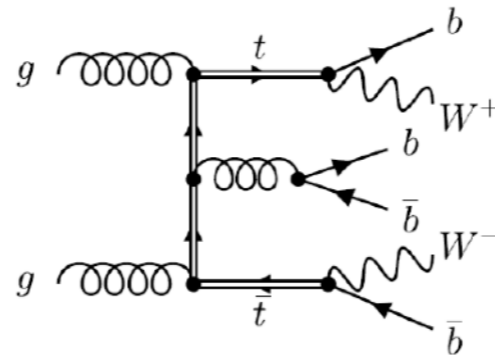
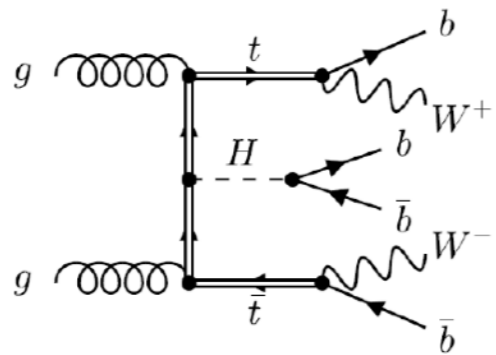
- **Genuine multiscale process** with characteristic scales typically separated by one order of magnitude \rightarrow test of perturbative QCD
- Several $t\bar{t}$ observables are **sensitive to extra-jet radiation**
- $t\bar{t} + 1$ jet provides one method to extract top quark mass at the LHC

Alioli, Fuster, Irlles, Moch and Uwer '13

Alioli, Fuster, Garzelli, Gavardi, Irlles, Melini, Moch, Uwer and Voss '22

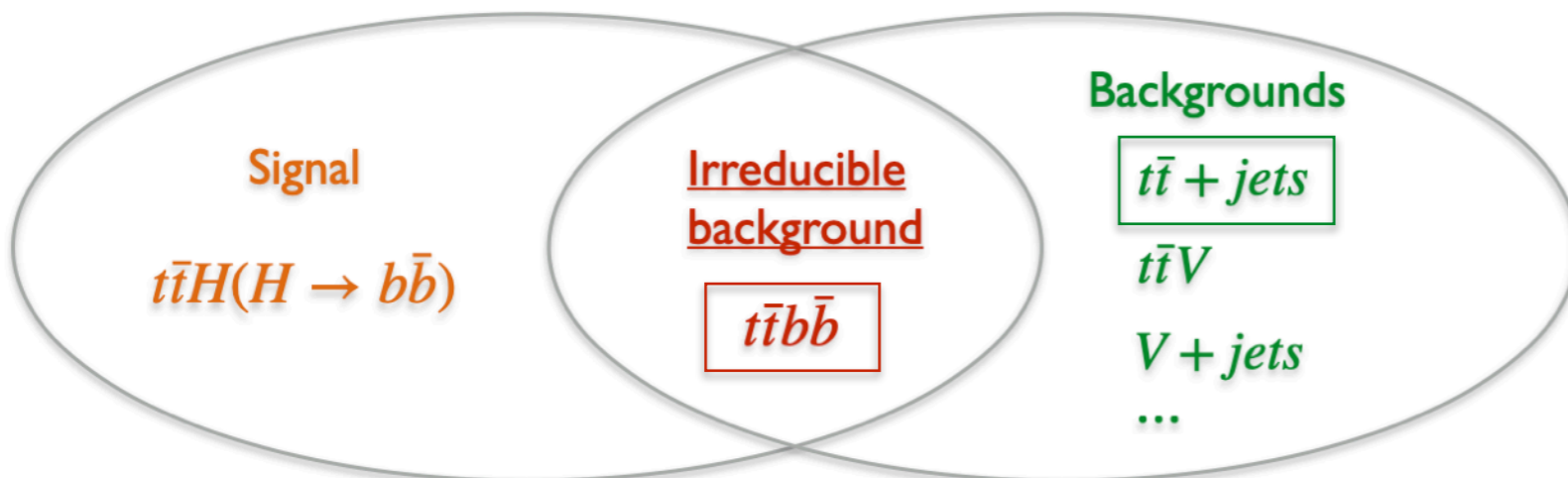
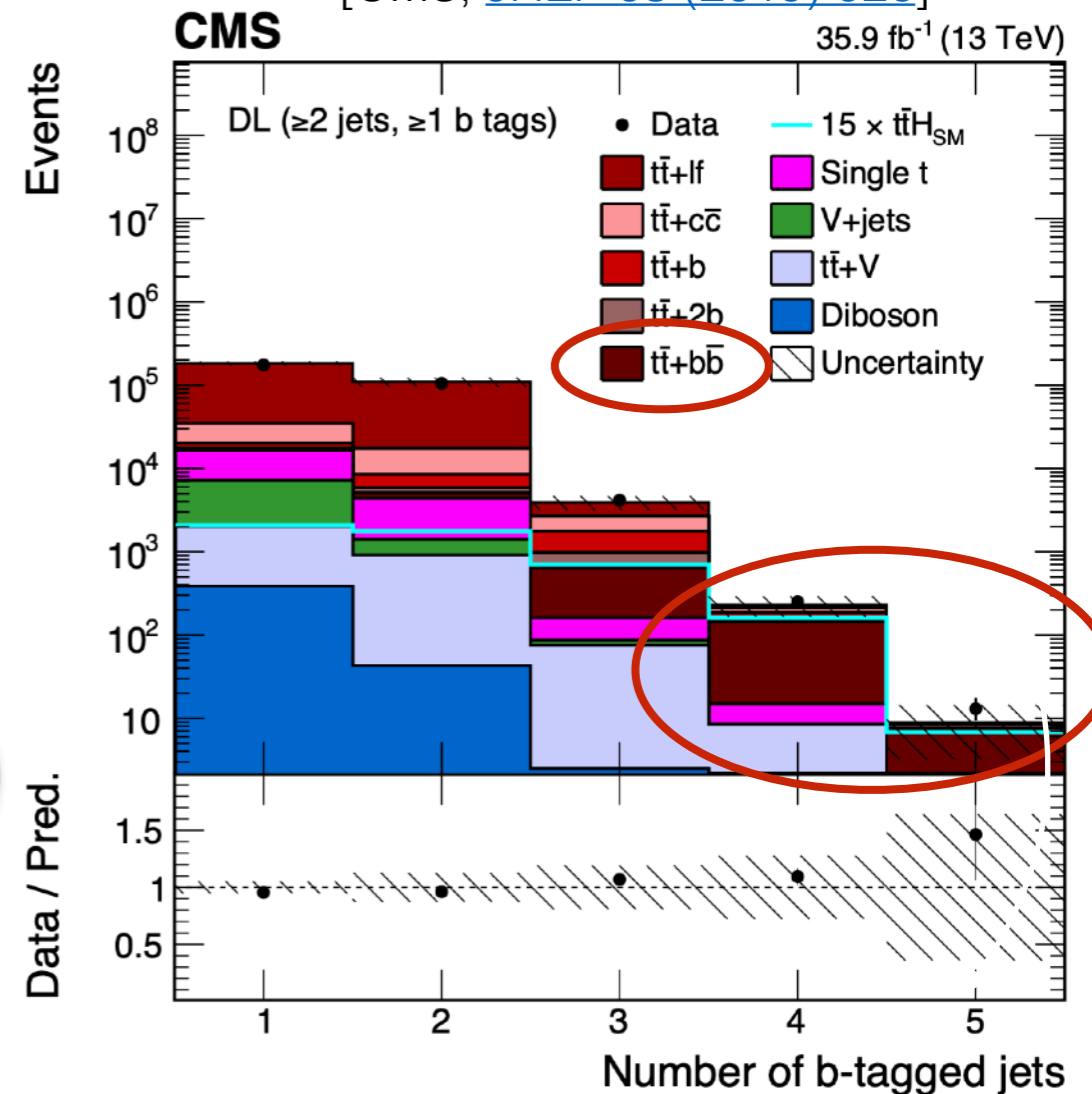
$t\bar{t}$ +jets as a background

- Background to $t\bar{t}H(H \rightarrow b\bar{b})$ production (and to many BSM searches)



[CMS, [JHEP 03 \(2019\) 026](#)]

- $pp \rightarrow t\bar{t}H$: direct sensitivity to top Yukawa coupling
- $H \rightarrow b\bar{b}$: largest BR ($\sim 58\%$)
- $t\bar{t}H(H \rightarrow b\bar{b})$ is a tiny signal in a huge background



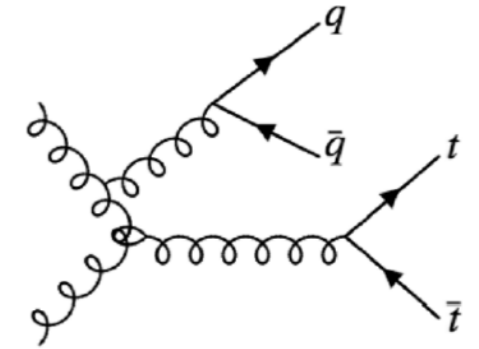
$t\bar{t}$ differential measurements (I)

[ATLAS, [JHEP 10 \(2018\) 159](#)]

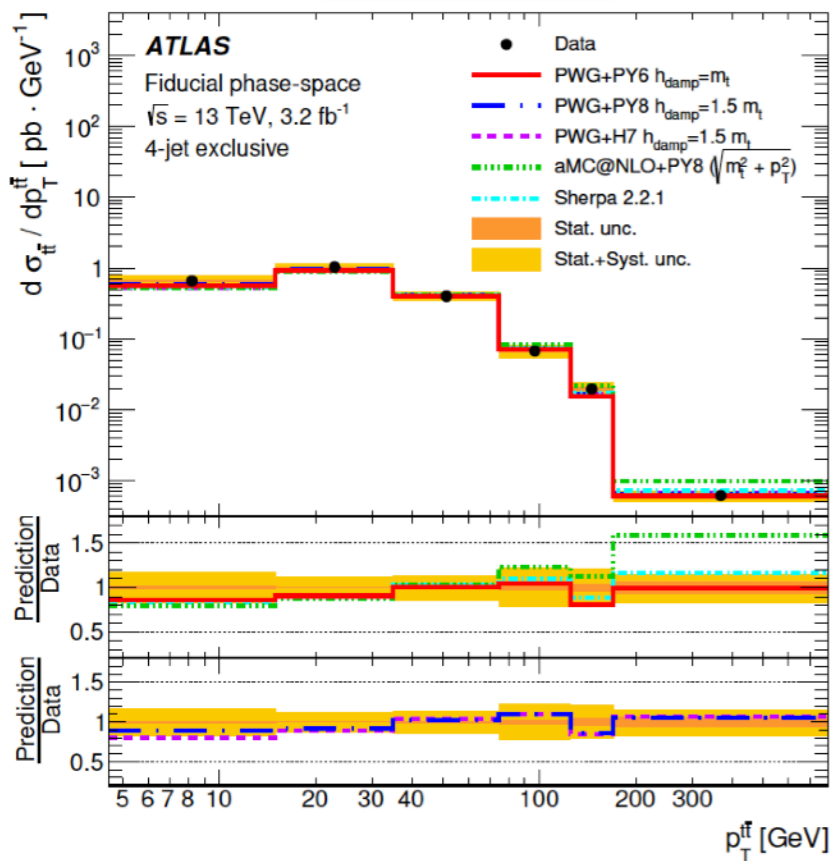
- Reconstruction of $t\bar{t}$ kinematics in $\ell + \text{jets}$ channel

- $t\bar{t}$ kinematics measured for different additional jet multiplicities: $N_{j,\text{add}}$

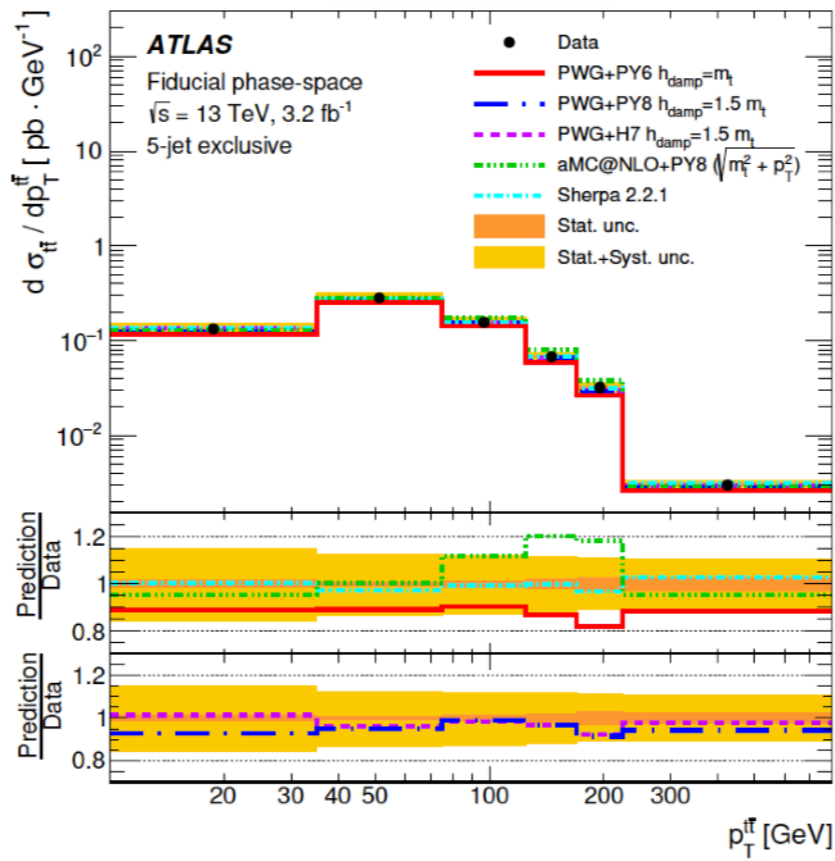
- Some observables sensitive to extra-jet radiation, e.g. $p_T(t\bar{t})$:



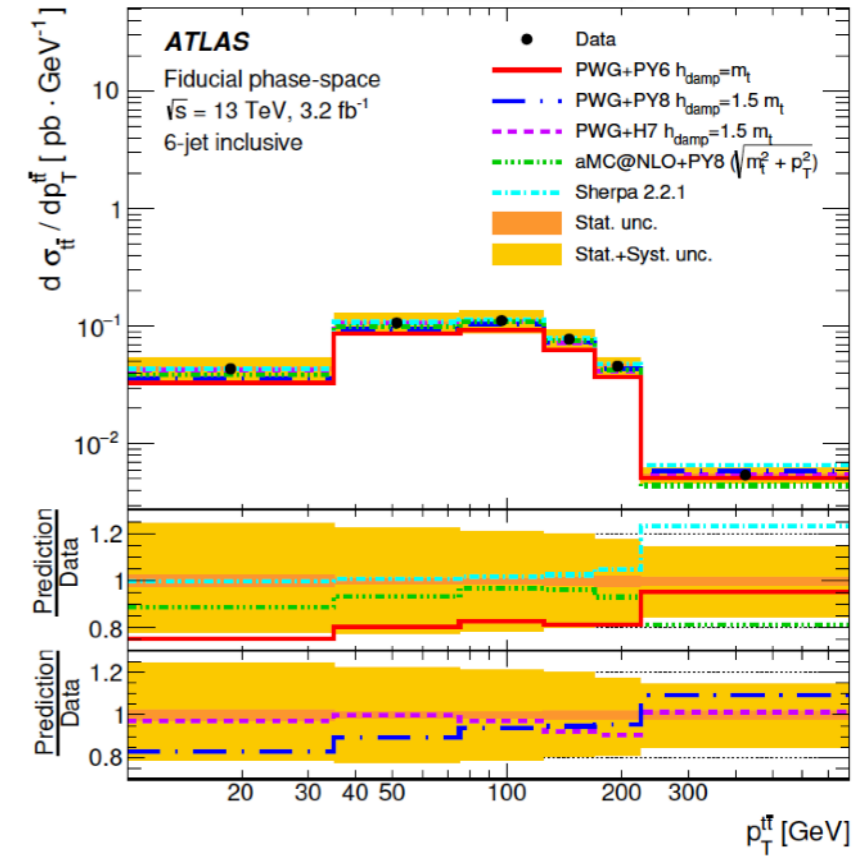
$N_{j,\text{add}} = 0$



$N_{j,\text{add}} = 1$



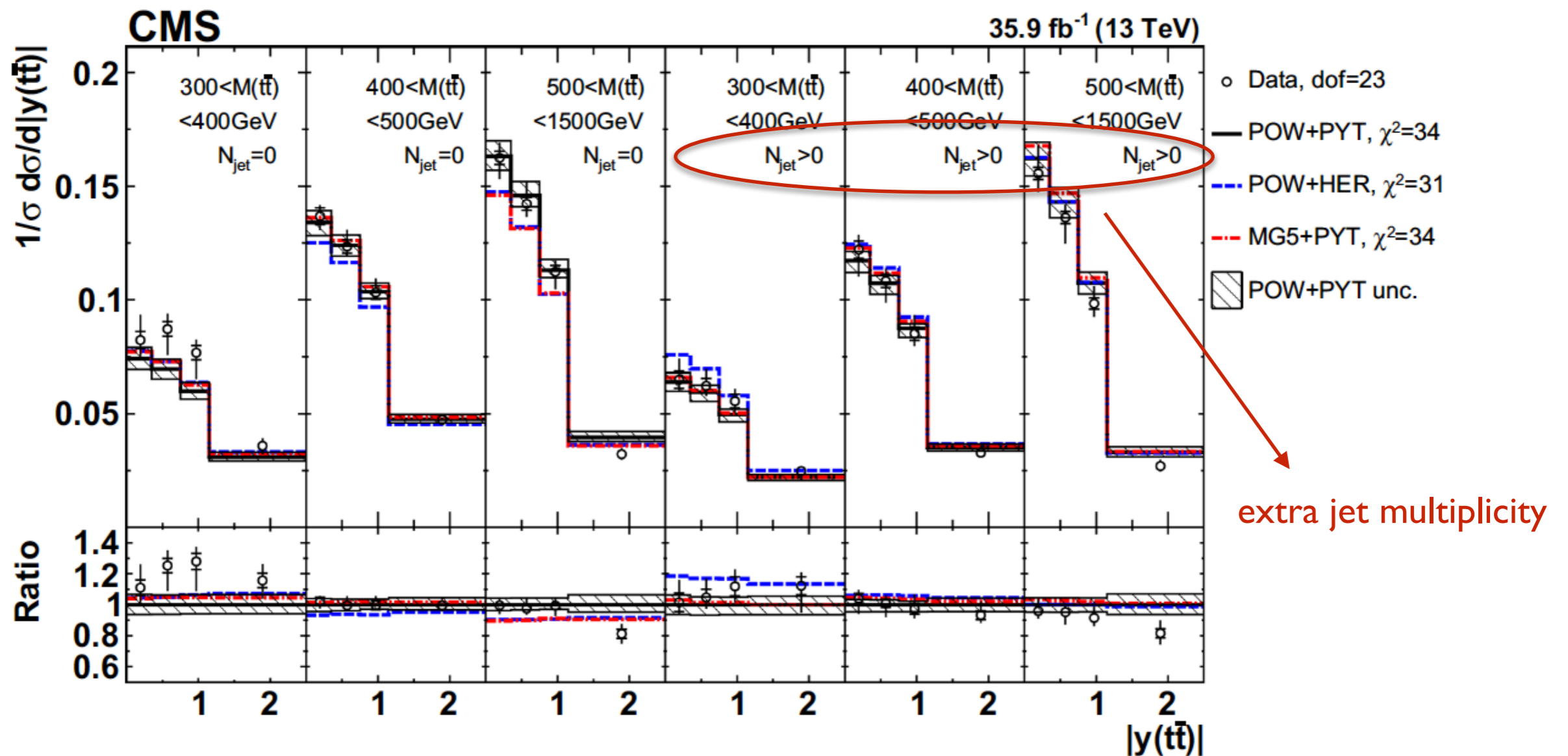
$N_{j,\text{add}} \geq 2$



$t\bar{t}$ differential measurements (II)

[CMS, [EPJ C80 \(2020\) 658](#)]

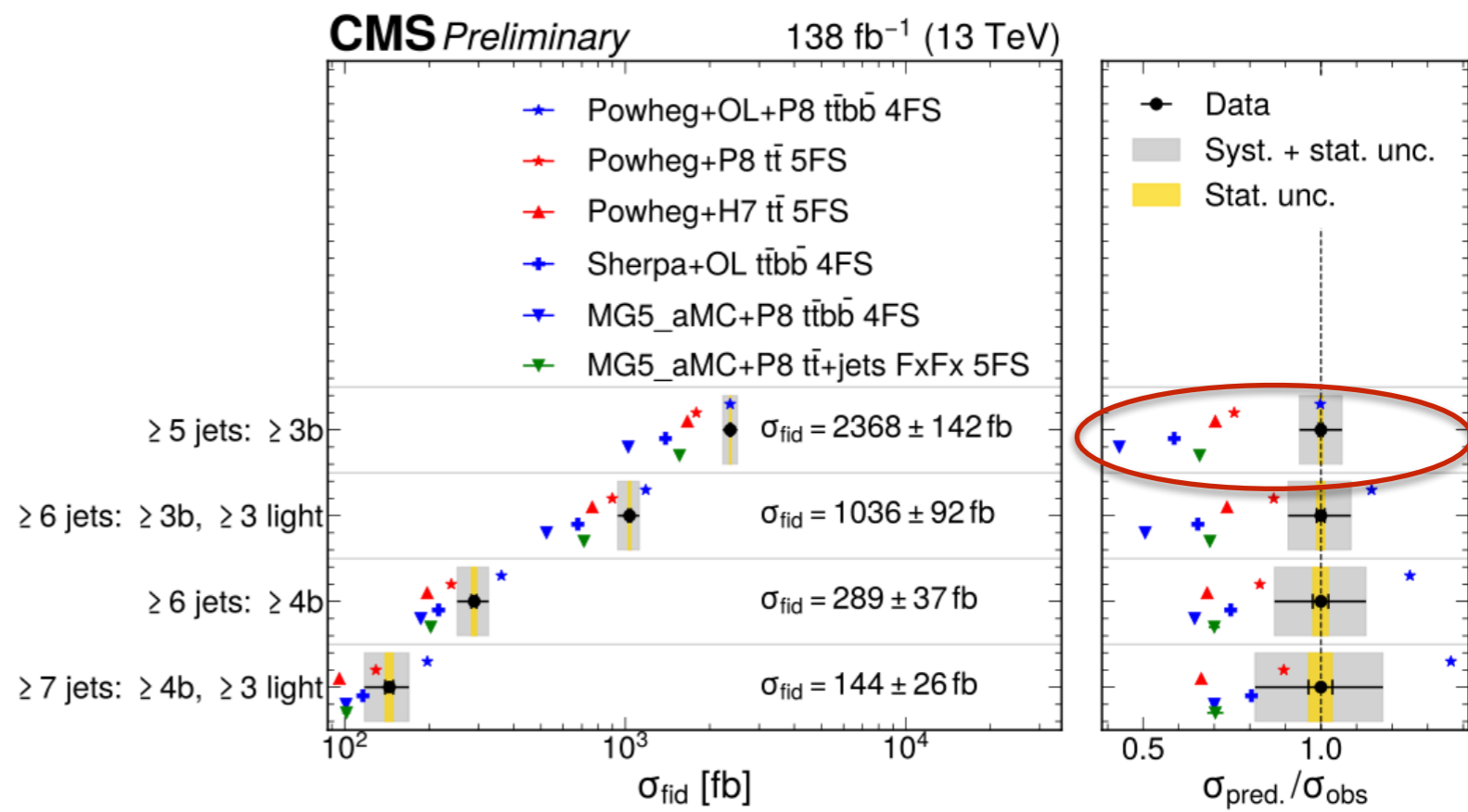
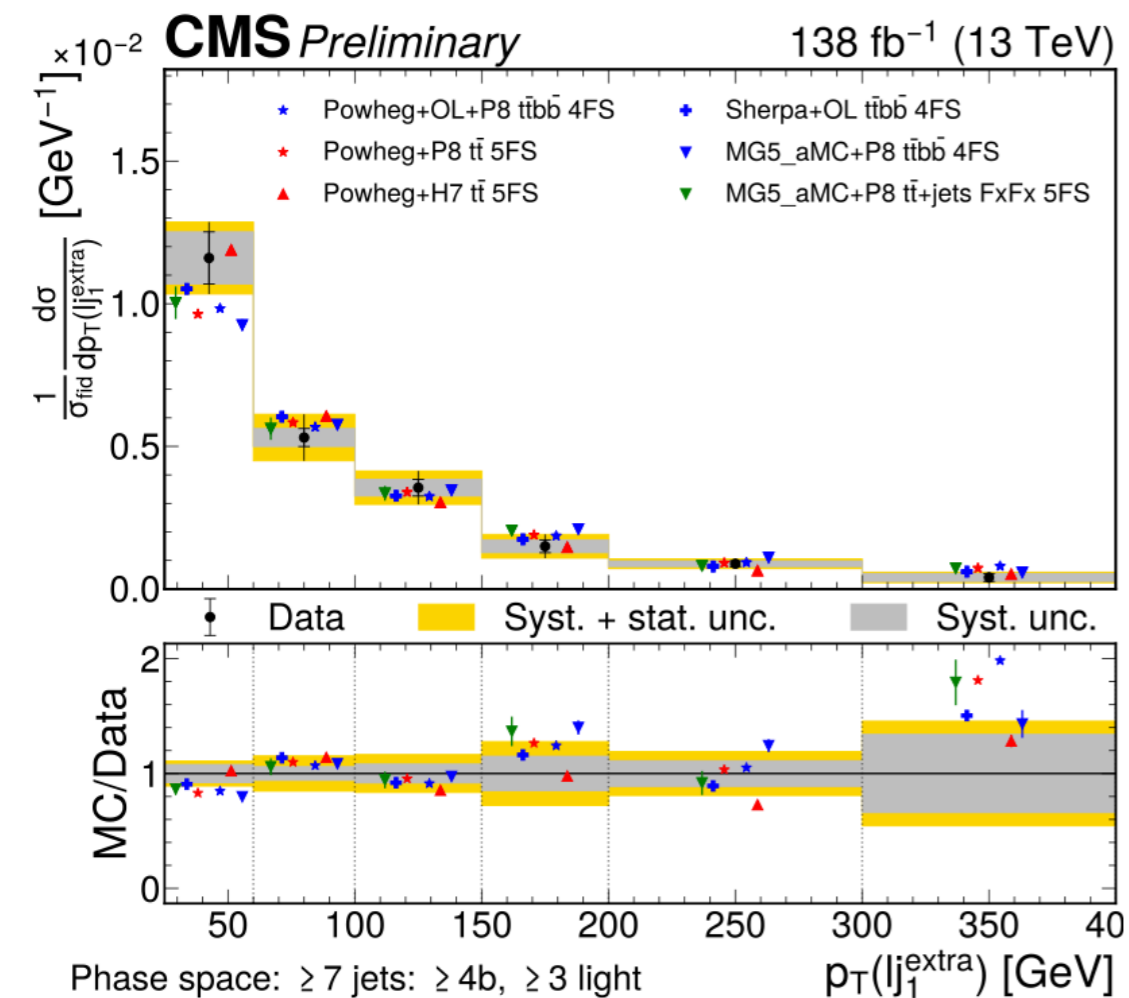
- Dilepton channel (e^+e^- , $\mu^+\mu^-$, $e^\pm\mu^\pm$)
- At least two jets in the final state; at least one of them b -tagged
- Triple-differential measurements \rightarrow simultaneous extraction of α_s , m_t , PDFs



Measurements of $t\bar{t}b\bar{b}$ production

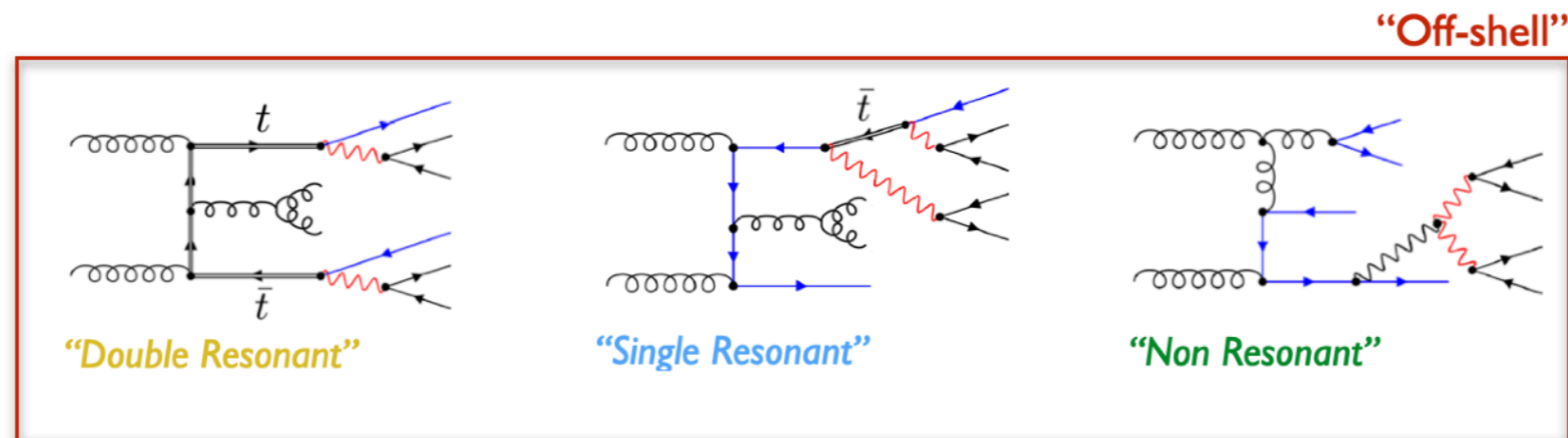
[[CMS-PAS-TOP-22-009](#)]

- ℓ +jets channel
- Inclusive and differential measurements \rightarrow 4 fiducial regions: $5j3b$, $6j4b$, $6j3b3\ell$, $7j4b3\ell$
- Inclusive σ higher than theory predictions (consistent with previous measurements)



$t\bar{t}$ + multijets at LHC: theory challenges

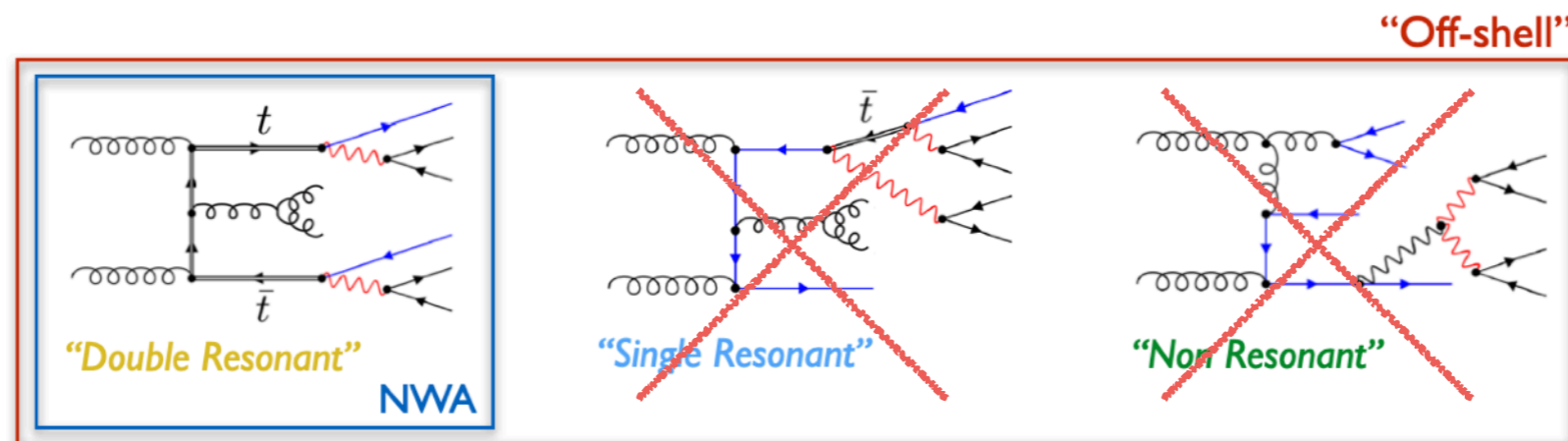
- ME calculations for fully decayed final states are often challenging



$$\text{“Off-shell”} = \text{DR} + \text{SR} + \text{NR} + \text{interferences}$$

$t\bar{t}$ + multijets at LHC: theory challenges

- ME calculations for fully decayed final states are often challenging \rightarrow use NWA

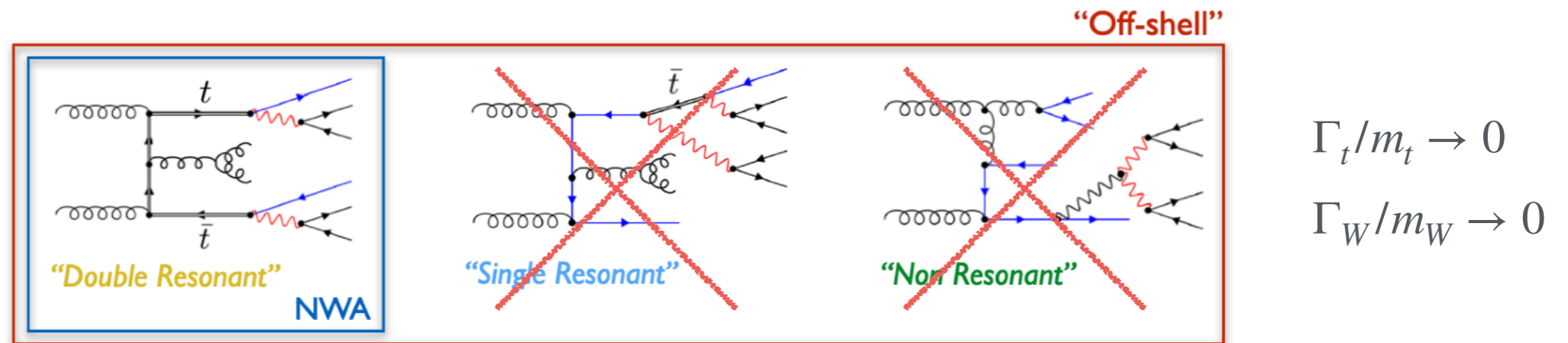


$$\Gamma_t/m_t \rightarrow 0$$

$$\Gamma_W/m_W \rightarrow 0$$

$t\bar{t}$ + multijets at LHC: theory challenges

- ME calculations for fully decayed final states are often challenging \rightarrow use NWA



$$\Gamma_t/m_t \rightarrow 0 \quad \Rightarrow \quad \frac{1}{(p_t^2 - m_t^2)^2 + m_t^2 \Gamma_t^2} \stackrel{\Gamma_t \rightarrow 0}{\sim} \frac{\pi}{m_t \Gamma_t} \delta(p_t^2 - m_t^2) + \mathcal{O}\left(\frac{\Gamma_t}{m_t}\right)$$

- Off-shell effects are suppressed by powers of $\Gamma_t/m_t = \mathcal{O}(1\%)$ for sufficiently inclusive observables

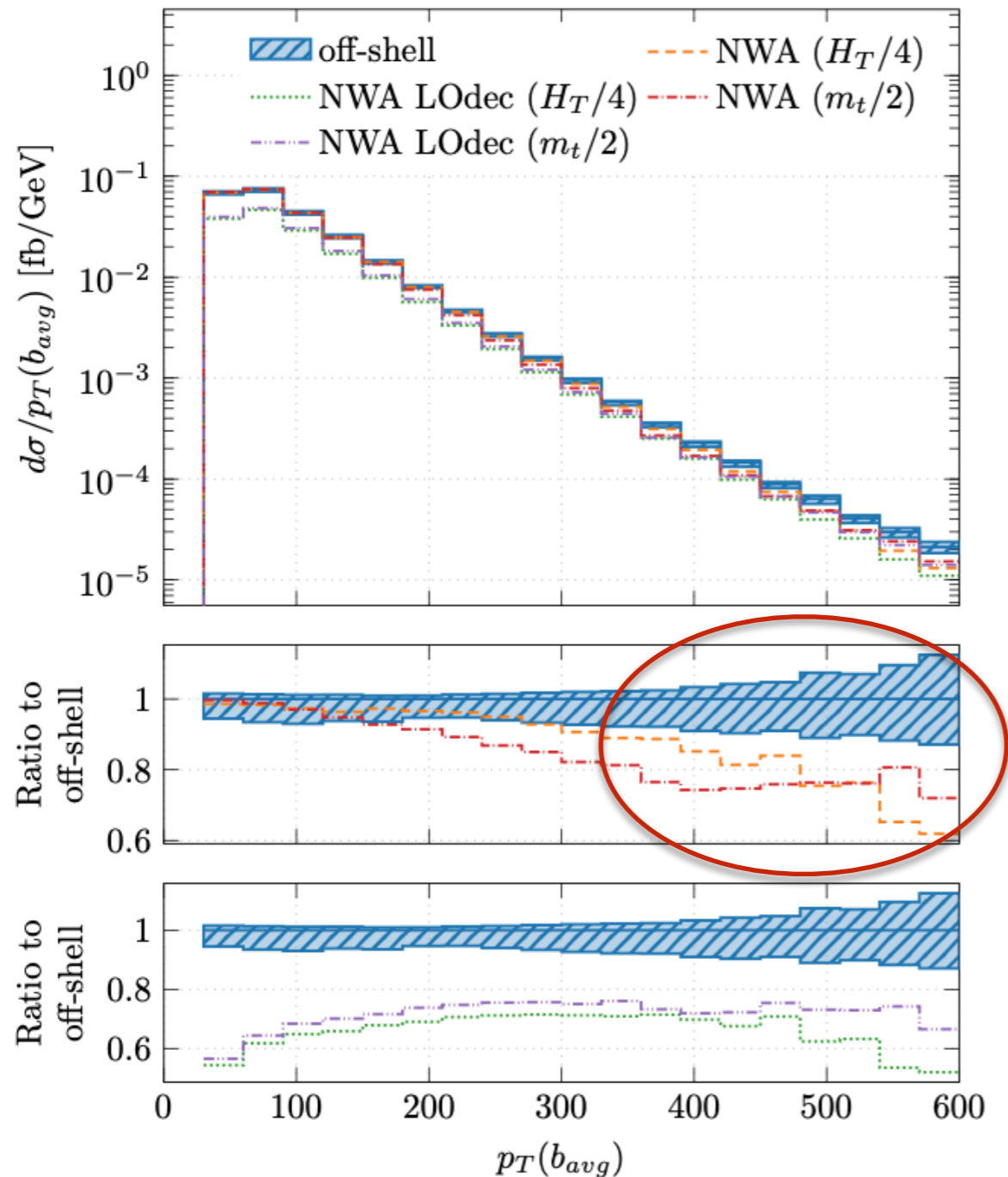
[Fadin, Khoze and Martin, *Phys.Lett.B* 320 (1994) 141-144]

Note: in general, this is not true at differential level

$t\bar{t}$ + multijets at LHC: theory challenges

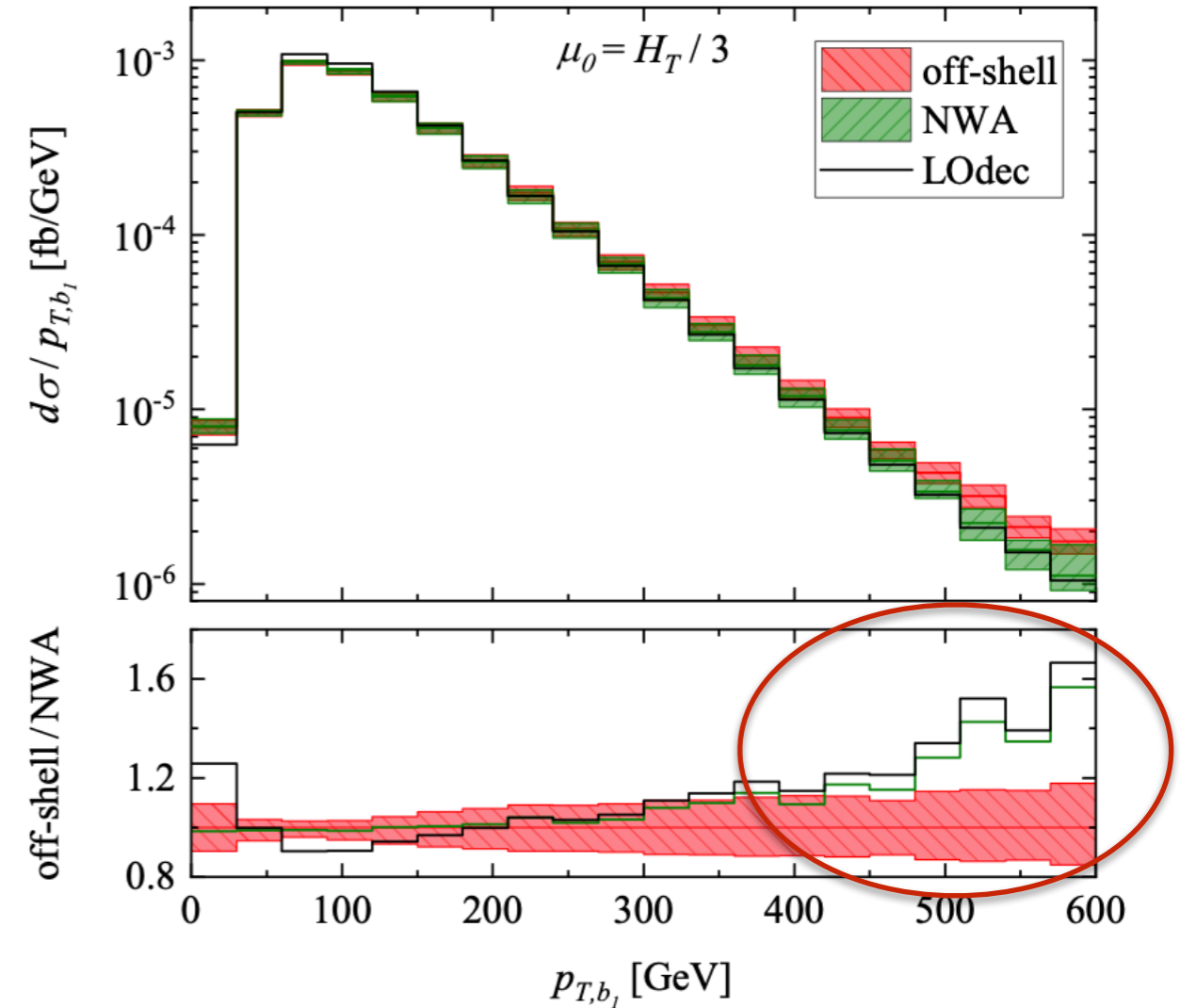
$$pp \rightarrow t\bar{t}\gamma \text{ (dilepton)}$$

GB, Hartanto, Kraus, Weber and Worek, *JHEP* 03 (2020) 154



$$pp \rightarrow t\bar{t}W^\pm \text{ (dilepton)}$$

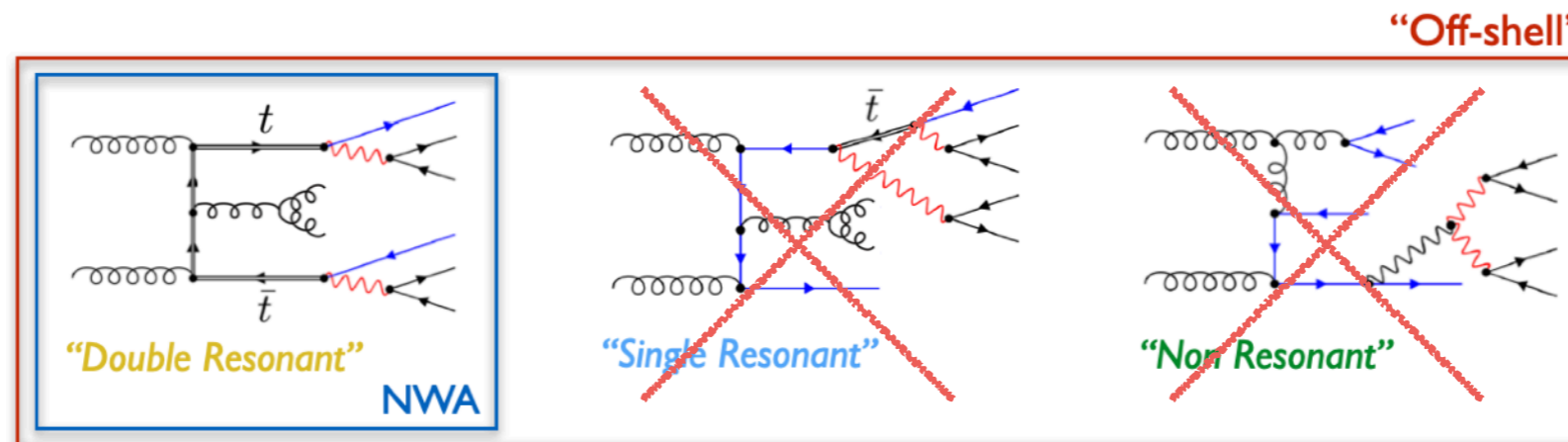
GB, Bi, Hartanto, Kraus and Worek, *JHEP* 08 (2020) 043



- Off-shell effects can reach $\mathcal{O}(50\%)$ and more differentially!

$t\bar{t}$ + multijets at LHC: theory challenges

- ME calculations for fully decayed final states are often challenging \rightarrow use NWA

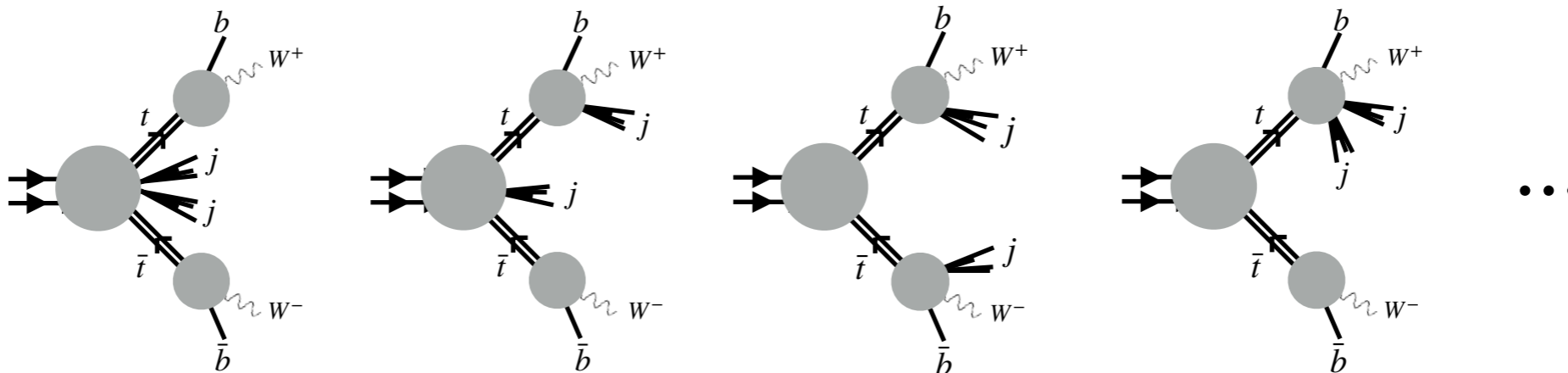


$$\Gamma_t/m_t \rightarrow 0$$

$$\Gamma_W/m_W \rightarrow 0$$

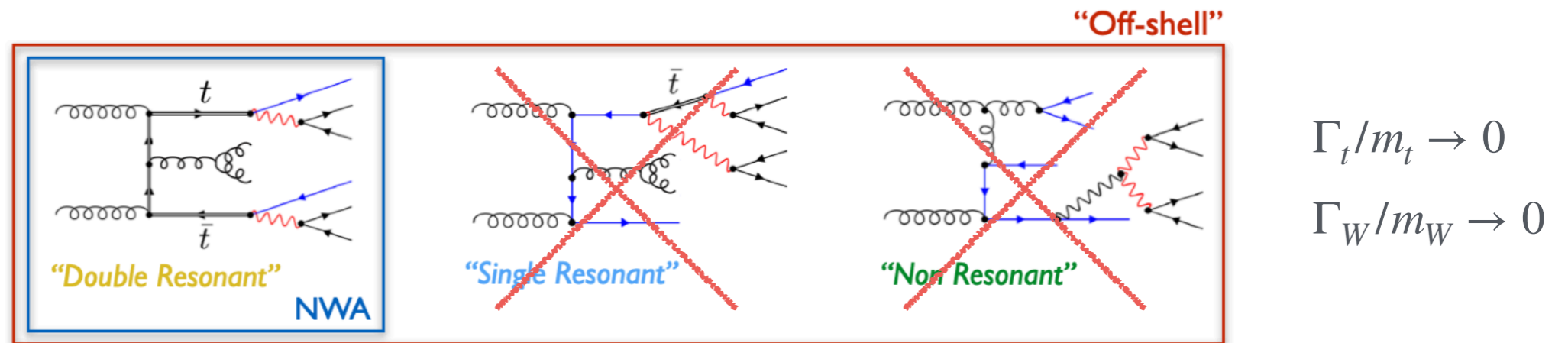
- The number of resonant structures entering the $t\bar{t} + n$ jets cross section in NWA increases rapidly with n

e.g.: $n = 2$



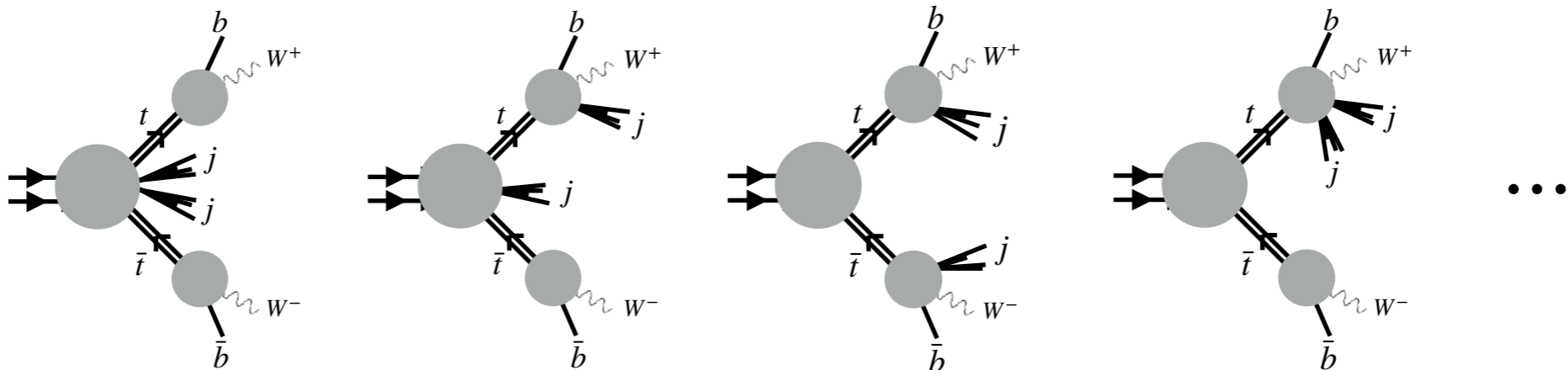
$t\bar{t}$ + multijets at LHC: theory challenges

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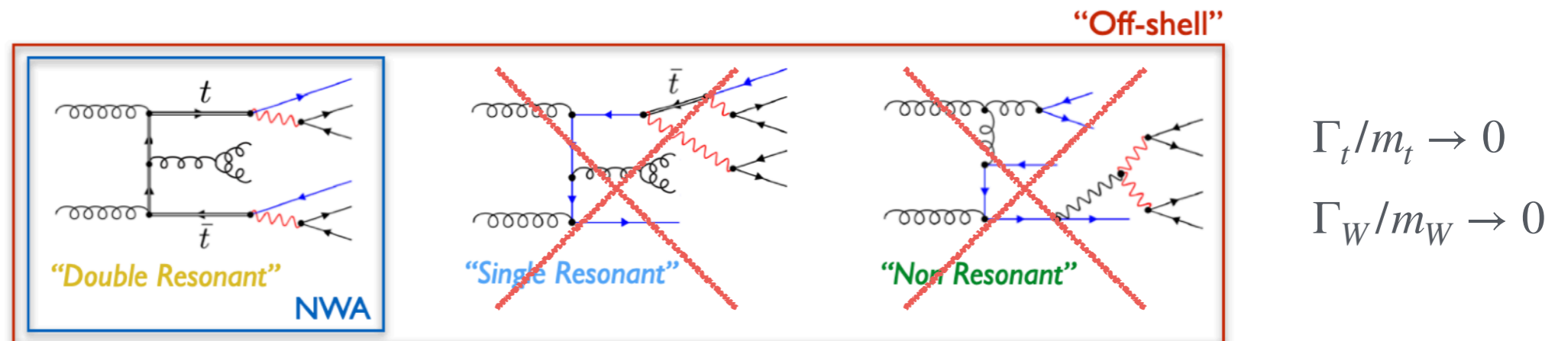
e.g.: $n = 2$



- QCD corrections to both *Production* and *Decay* ME's should be considered for accurate estimates of full NLO cross section

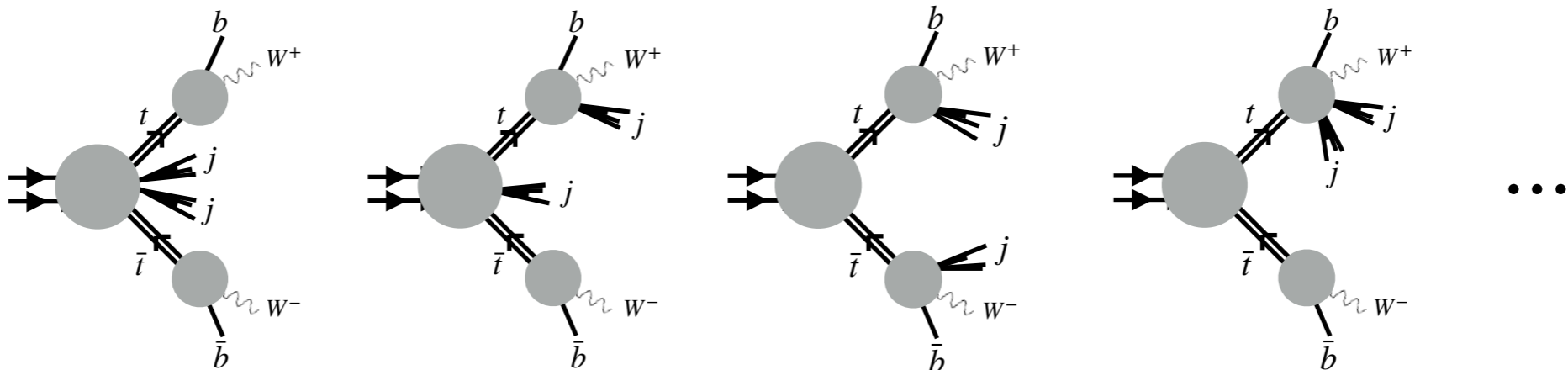
$t\bar{t}$ + multijets at LHC: theory challenges

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e.g.: $n = 2$



- QCD corrections to both *Production* and *Decay* ME's should be considered for accurate estimates of full NLO cross section

We'll focus on the case of $t\bar{t} + 2$ jets

Stable top quarks

- $pp \rightarrow t\bar{t} + 2 \text{ jets}$

[GB, Czakon, Papadopoulos and Worek '10,'11]

↪ NLO QCD: fixed-order

- $pp \rightarrow t\bar{t} + 0,1,2,3 \text{ jets}$

[Höche, Maierhöfer, Moretti, Pozzorini and Siegert '17]

↪ NLO QCD: NLO vs MiNLO

Exclusive final states

- $pp \rightarrow t\bar{t} + 0,1,2 \text{ jets}$

[Höche, Krauss, Maierhöfer, Pozzorini, Schönherr and Siegert '15]

↪ NLO QCD: MEPS@NLO multi-jet merging

- $pp \rightarrow t\bar{t} + 0,1,2,3,4 \text{ jets}$

[Gütschow, Lindert and Schönherr '18]

↪ NLO QCD+EW ($n \leq 1$ jet) & LO QCD+EW ($n > 1$): MEPS@NLO multi-jet merging

Fixed order

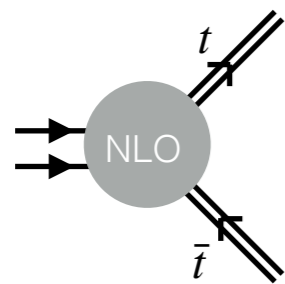
- $pp \rightarrow t\bar{t}b\bar{b}$ → stable tops [Bredenstein, Denner, Dittmaier and Pozzorini '09]
[GB, Czakon, Papadopoulos, Pittau and Worek '09]
 - $pp \rightarrow t\bar{t}b\bar{b}j$ → stable tops [Buccioni, Kallweit, Pozzorini and Zoller '19]
 - $pp \rightarrow e^+\nu_e\mu^-\bar{\nu}_\mu b\bar{b}b\bar{b}$ [Denner, Lang and Pellen '21]
[GB, Bi, Hartanto, Kraus, Lupattelli and Worek '21]
- ↪ Full off-shell

Matched to PS

- $pp \rightarrow t\bar{t}b\bar{b}$ [Garzelli, Kardos and Trocsanyi '14] 5FS
[GB, Garzelli and Kardos '17] 4FS
[Jezo, Lindert, Moretti and Pozzorini '18] 4FS
↪ POWHEG matching
- $pp \rightarrow t\bar{t}b\bar{b}$ [Cascioli, Maierhöfer, Moretti, Pozzorini and Siebert '14] 4FS
↪ MC@NLO matching

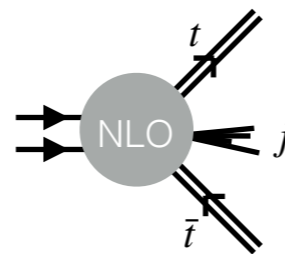
$t\bar{t}jj$: state-of-the-art in a nutshell

- State-of-the-art of $t\bar{t}jj$ MC simulations: **NLO+PS** (merging multijet samples)



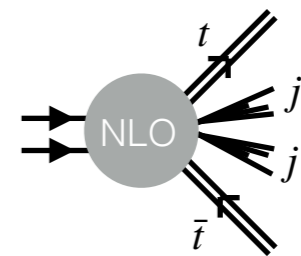
$pp \rightarrow t\bar{t} + 0 \text{ jet}$

\oplus



$pp \rightarrow t\bar{t} + 1 \text{ jet}$

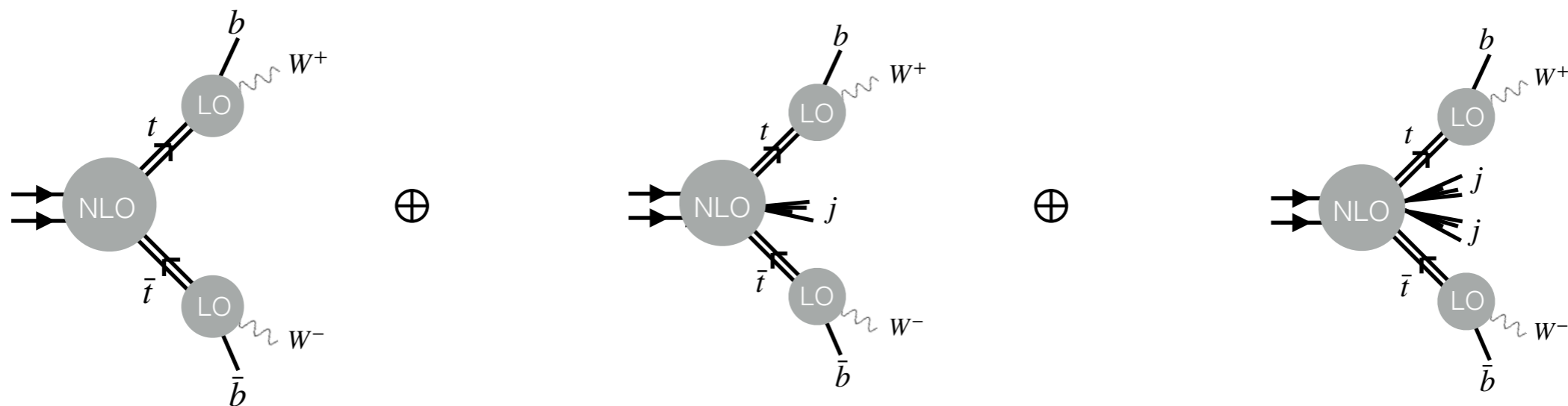
\oplus



$pp \rightarrow t\bar{t} + 2 \text{ jets}$

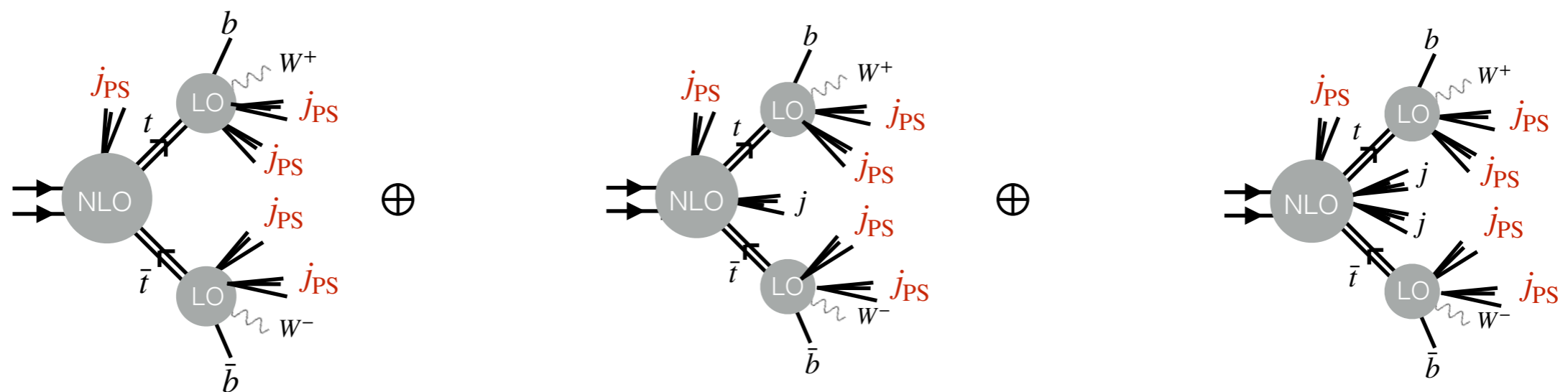
$t\bar{t}jj$: state-of-the-art in a nutshell

- State-of-the-art of $t\bar{t}jj$ MC simulations: NLO+PS (merging multijet samples)
- Top quarks produced on-shell and decayed at **LO** with spin correlations



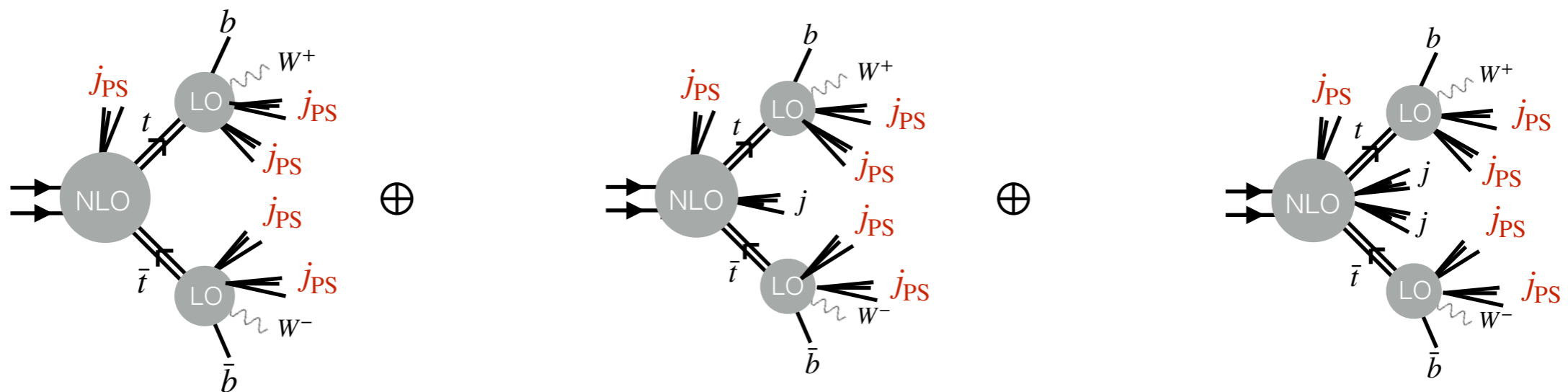
$t\bar{t}jj$: state-of-the-art in a nutshell

- State-of-the-art of $t\bar{t}jj$ MC simulations: NLO+PS (merging multijet samples)
- Top quarks produced on-shell and decayed at LO with spin correlations
- **Parton Shower** evolution (ISR/FSR) accounts for additional jet activity



Some interesting questions

- I. To what extent do **QCD corrections to decays** impact fiducial cross sections?
→ normalisation
- II. Which phase space regions are more sensitive to **radiation off decays**?
→ shapes
- III. What's the impact of **full off-shell effects**?
→ shapes, normalisation



- NLO QCD analysis of $pp \rightarrow t\bar{t}jj$ (dilepton) — NWA

- Anatomy of resonant contributions at NLO QCD
- Effects of hard radiation off top quark decays

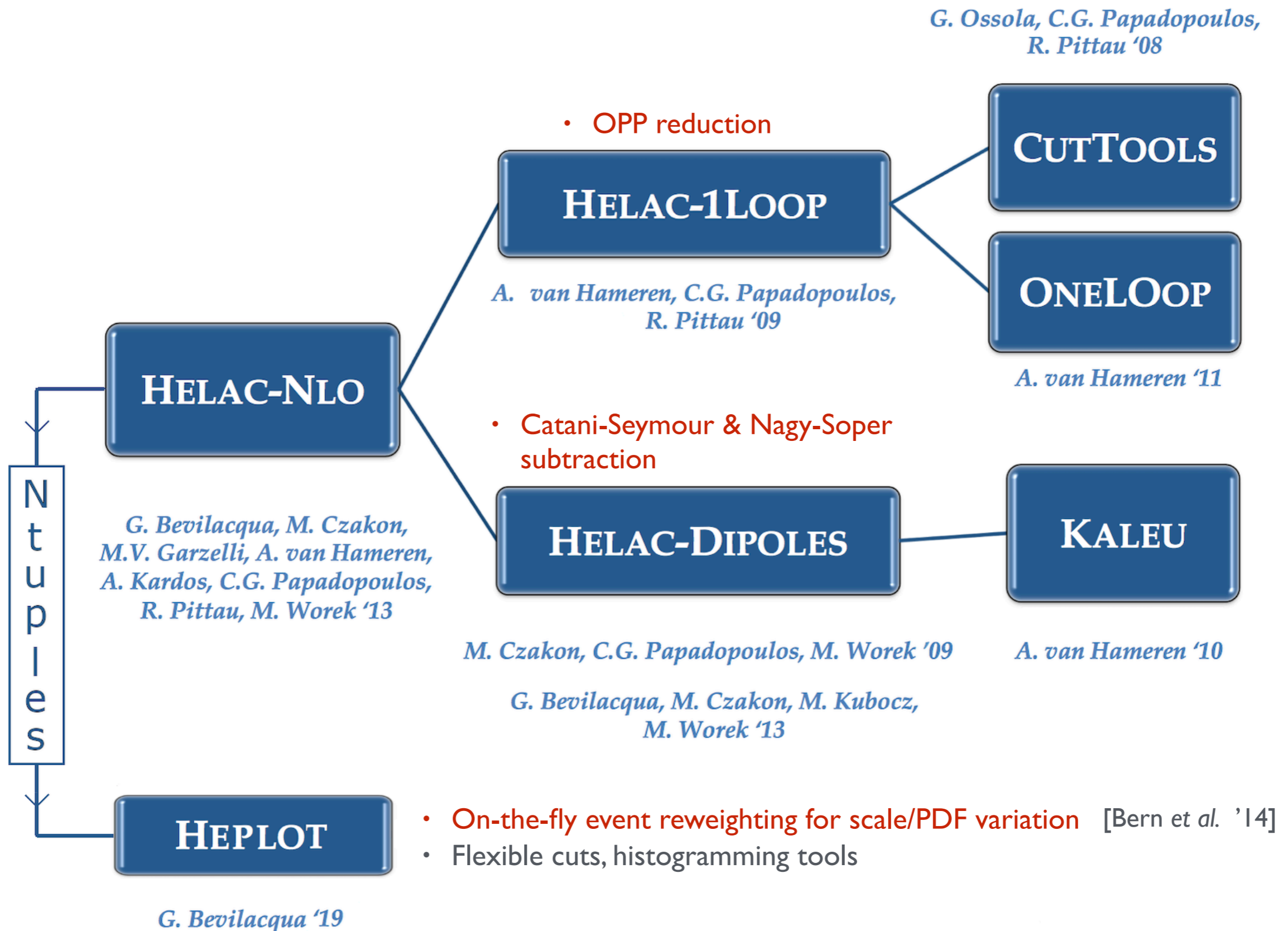
- NLO QCD analysis of $pp \rightarrow t\bar{t}b\bar{b}$ (dilepton) — full off-shell

- Impact of genuine off-shell effects
- Categorization of prompt b -jets

Results obtained with the HELAC-NLO framework

[GB, Czakon, Garzelli, van Hameren, Kardos, Papadopoulos, Pittau and Worek '13]

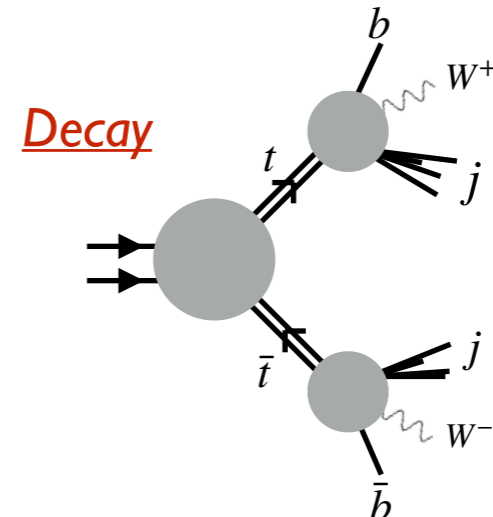
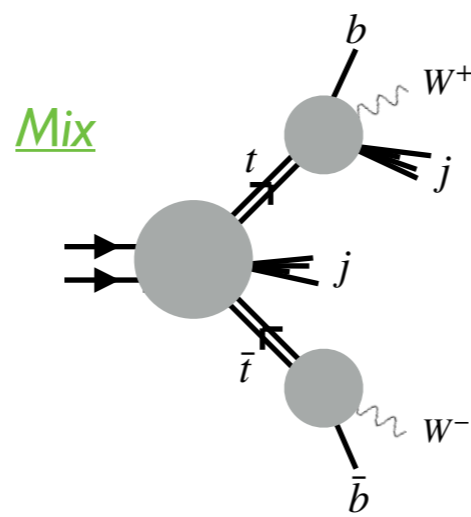
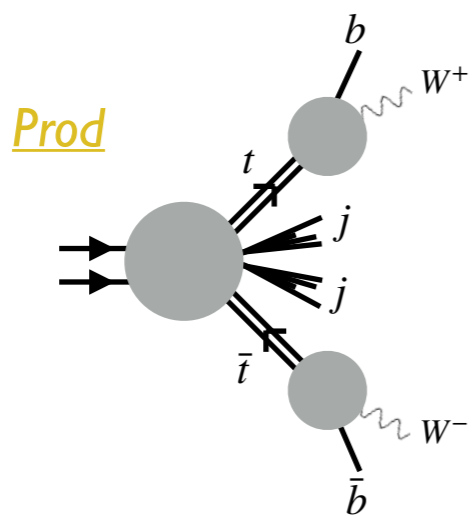
Computational framework



I. A study of jet activity in $t\bar{t}jj$ at the LHC

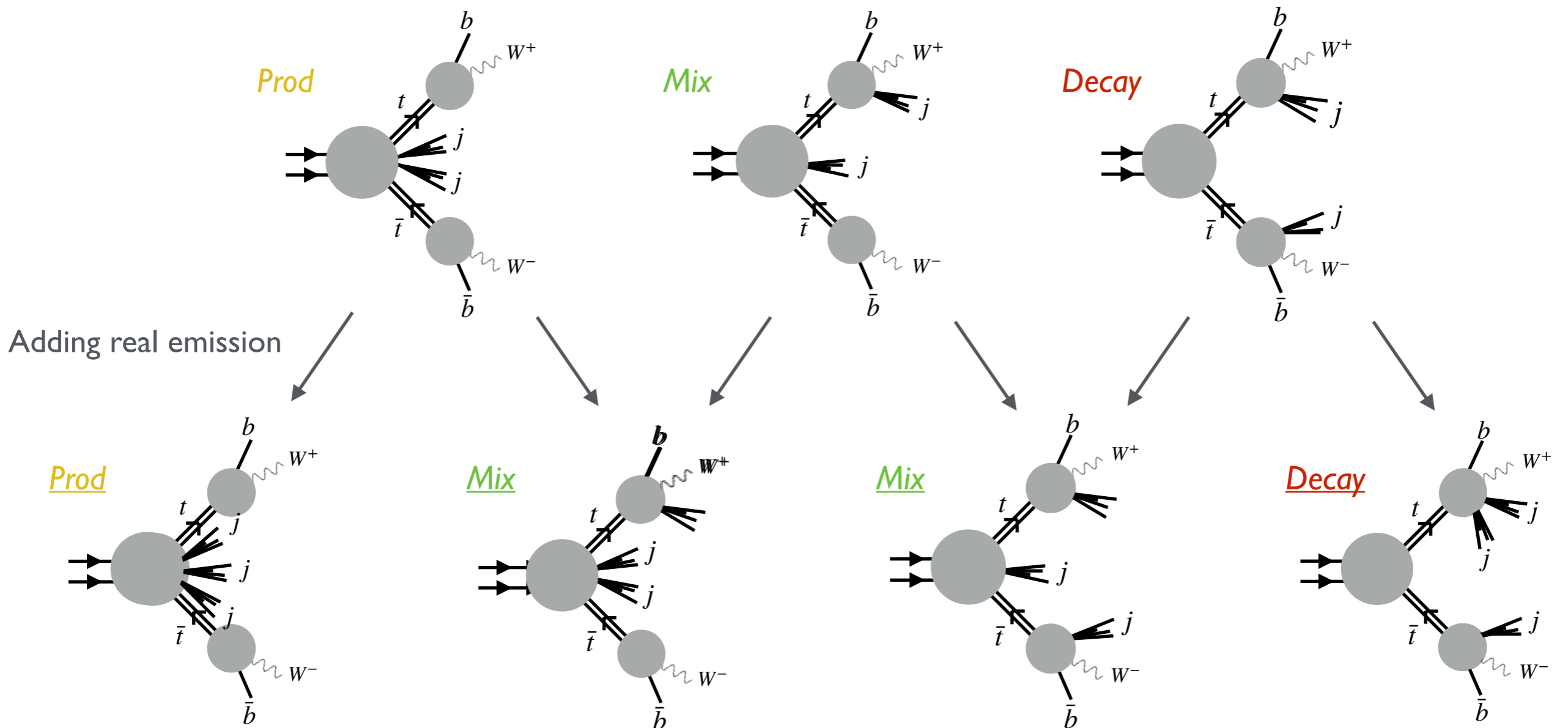
Resonant contributions to $t\bar{t}jj$ in NWA: LO

$$d\sigma_{t\bar{t}jj}^{\text{LO}} = \Gamma_t^{-2} \left(\underbrace{d\sigma_{t\bar{t}jj}^{\text{LO}} d\Gamma_{t\bar{t}}^{\text{LO}}}_{\text{Prod.}} + \underbrace{d\sigma_{t\bar{t}j}^{\text{LO}} d\Gamma_{t\bar{t}j}^{\text{LO}}}_{\text{Mix}} + \underbrace{d\sigma_{t\bar{t}}^{\text{LO}} d\Gamma_{t\bar{t}jj}^{\text{LO}}}_{\text{Decay}} \right)$$



Resonant contributions to $t\bar{t}jj$ in NWA: NLO

$$d\sigma_{t\bar{t}jj}^{\text{NLO}} = \Gamma_t^{-2} \left(\underbrace{\left(d\sigma_{t\bar{t}jj}^{\text{LO}} + d\sigma_{t\bar{t}jj}^{\text{virt}} + d\sigma_{t\bar{t}jj}^{\text{real}} \right)}_{\text{Prod.}} d\Gamma_{t\bar{t}}^{\text{LO}} + \underbrace{d\sigma_{t\bar{t}}^{\text{LO}} \left(d\Gamma_{t\bar{t}jj}^{\text{LO}} + d\Gamma_{t\bar{t}jj}^{\text{virt}} + d\Gamma_{t\bar{t}jj}^{\text{real}} \right)}_{\text{Decay}} \right. \\ \left. + \underbrace{d\sigma_{t\bar{t}j}^{\text{LO}} d\Gamma_{t\bar{t}j}^{\text{LO}} + d\sigma_{t\bar{t}jj}^{\text{LO}} d\Gamma_{t\bar{t}}^{\text{virt}} + d\sigma_{t\bar{t}}^{\text{virt}} d\Gamma_{t\bar{t}jj}^{\text{LO}} + d\sigma_{t\bar{t}j}^{\text{virt}} d\Gamma_{t\bar{t}j}^{\text{LO}} + d\sigma_{t\bar{t}j}^{\text{LO}} d\Gamma_{t\bar{t}j}^{\text{virt}} + d\sigma_{t\bar{t}jj}^{\text{real}} d\Gamma_{t\bar{t}j}^{\text{real}} + d\sigma_{t\bar{t}j}^{\text{real}} d\Gamma_{t\bar{t}jj}^{\text{real}}}_{\text{Mix}} \right)$$



$t\bar{t}jj$: fiducial cross sections

Integrated fiducial cross sections

[GB, Lupattelli, Stremmer and Worek, [Phys. Rev. D \(107\) 2023, 114027](#)]

Modelling	σ^{LO} [fb]	σ^{NLO} [fb]	$\frac{\sigma_i^{\text{LO}}}{\sigma_{\text{NWA}}^{\text{LO}}}$	$\frac{\sigma_i^{\text{NLO}}}{\sigma_{\text{NWA}}^{\text{NLO}}}$
NWA _{full}	868.8(2) ^{+60%} _{-35%}	1225(1) ^{+1%} _{-14%}	1.00	1.00
Prod	843.2(2) ^{+60%} _{-35%}	1462(1) ^{+12%} _{-19%}	0.97	1.19
Mix	25.465(5)	-236(1)	0.029	-0.19
Decay	0.2099(1)	0.1840(8)	0.0002	0.0002
NWA _{full,exp}	—	1173(1) ^{+7%} _{-16%}	—	0.96
NWA _{LOdec}	—	1222(1) ^{+12%} _{-19%}	—	0.998
$\mu_0 = H_T/2$ NNPDF3.1 PDF		$\Delta R(jb) > 0.8$		

Fiducial cuts:

$p_{T,\ell} > 20$ GeV ,
 $p_{T,b} > 30$ GeV ,
 $p_{T,j} > 40$ GeV ,
 $\Delta R_{bl} > 0.4$,
 $\Delta R_{\ell\ell} > 0.4$,
 $\Delta R_{jj} > 0.4$,
 $|y_\ell| < 2.4$,
 $|y_b| < 2.4$,
 $|y_j| < 2.4$,
 $\Delta R_{jl} > 0.4$,
 $\Delta R_{bb} > 0.4$,
 $\Delta R_{jb} > 0.8$ (0.4)

- Moderate QCD corrections: **+41 %**
- NLO uncertainties — Scale: **$\mathcal{O}(15\%)$** PDF: **$\mathcal{O}(2\% - 3\%)$**

$t\bar{t}jj$: fiducial cross sections

Integrated fiducial cross sections

[GB, Lupattelli, Stremmer and Worek, [Phys. Rev. D \(107\) 2023, 114027](#)]

Modelling	σ^{LO} [fb]	σ^{NLO} [fb]	$\frac{\sigma_i^{\text{LO}}}{\sigma_{\text{NWA}}^{\text{LO}}}$	$\frac{\sigma_i^{\text{NLO}}}{\sigma_{\text{NWA}}^{\text{NLO}}}$
NWA _{full}	868.8(2) ^{+60%} _{-35%}	1225(1) ^{+1%} _{-14%}	1.00	1.00
Prod	843.2(2) ^{+60%} _{-35%}	1462(1) ^{+12%} _{-19%}	0.97	1.19
Mix	25.465(5)	-236(1)	0.029	-0.19
Decay	0.2099(1)	0.1840(8)	0.0002	0.0002
NWA _{full,exp}	—	1173(1) ^{+7%} _{-16%}	—	0.96
NWA _{LOdec}	—	1222(1) ^{+12%} _{-19%}	—	0.998
$\mu_0 = H_T/2$ NNPDF3.1 PDF		$\Delta R(jb) > 0.8$		

Fiducial cuts:

$p_{T,\ell} > 20$ GeV ,
 $p_{T,b} > 30$ GeV ,
 $p_{T,j} > 40$ GeV ,
 $\Delta R_{bl} > 0.4$,
 $\Delta R_{\ell\ell} > 0.4$,
 $\Delta R_{jj} > 0.4$,
 $|y_\ell| < 2.4$,
 $|y_b| < 2.4$,
 $|y_j| < 2.4$,
 $\Delta R_{jl} > 0.4$,
 $\Delta R_{bb} > 0.4$,
 $\Delta R_{jb} > 0.8$ (0.4)

- At **LO**: *Prod* is dominant, *Mix* and *Decay* are negligible (and all positive)
- At **NLO**: non-negligible and *negative* contribution from *Mix*: **-19%**

$t\bar{t}jj$: fiducial cross sections

Integrated fiducial cross sections

[GB, Lupattelli, Stremmer and Worek, [Phys. Rev. D \(107\) 2023, 114027](#)]

Modelling	σ^{LO} [fb]	σ^{NLO} [fb]	$\frac{\sigma_i^{\text{LO}}}{\sigma_{\text{NWA}}^{\text{LO}}}$	$\frac{\sigma_i^{\text{NLO}}}{\sigma_{\text{NWA}}^{\text{NLO}}}$
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$\mu_0 = H_T/2$ NNP3.1 PDF		$\Delta R(jb) > 0.8$		

Fiducial cuts:

$p_{T,\ell} > 20$ GeV ,
 $p_{T,b} > 30$ GeV ,
 $p_{T,j} > 40$ GeV ,
 $\Delta R_{bl} > 0.4$,
 $\Delta R_{\ell\ell} > 0.4$,
 $\Delta R_{jj} > 0.4$,
 $|y_\ell| < 2.4$,
 $|y_b| < 2.4$,
 $|y_j| < 2.4$,
 $\Delta R_{jl} > 0.4$,
 $\Delta R_{bb} > 0.4$,
 $\Delta R_{jb} > 0.8$ (0.4)

- NWA_{full} vs NWA_{LOdec} : permille level difference

How stable are these conclusions under different kinematical cuts ?

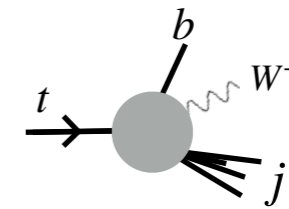
$t\bar{t}jj$: fiducial cross sections

[GB, Lupattelli, Stremmer and Worek, [Phys. Rev. D \(107\) 2023, 114027](#)]

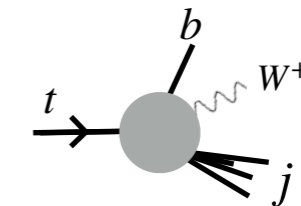
$\Delta R(jb) > 0.8$

Modelling	σ^{LO} [fb]	σ^{NLO} [fb]	$\frac{\sigma_i^{\text{LO}}}{\sigma_{\text{NWA}}^{\text{LO}}}$	$\frac{\sigma_i^{\text{NLO}}}{\sigma_{\text{NWA}}^{\text{NLO}}}$
NWA _{full}	868.8(2) ^{+60%} _{-35%}	1225(1) ^{+1%} _{-14%}	1.00	1.00
Prod	843.2(2) ^{+60%} _{-35%}	1462(1) ^{+12%} _{-19%}	0.97	1.19
Mix	25.465(5)	-236(1)	0.029	-0.19
Decay	0.2099(1)	0.1840(8)	0.0002	0.0002
NWA _{LOdec}	-	1222(1) ^{+12%} _{-19%}	-	0.998

“more suppressing”



“less suppressing”



$\Delta R(jb) > 0.4$

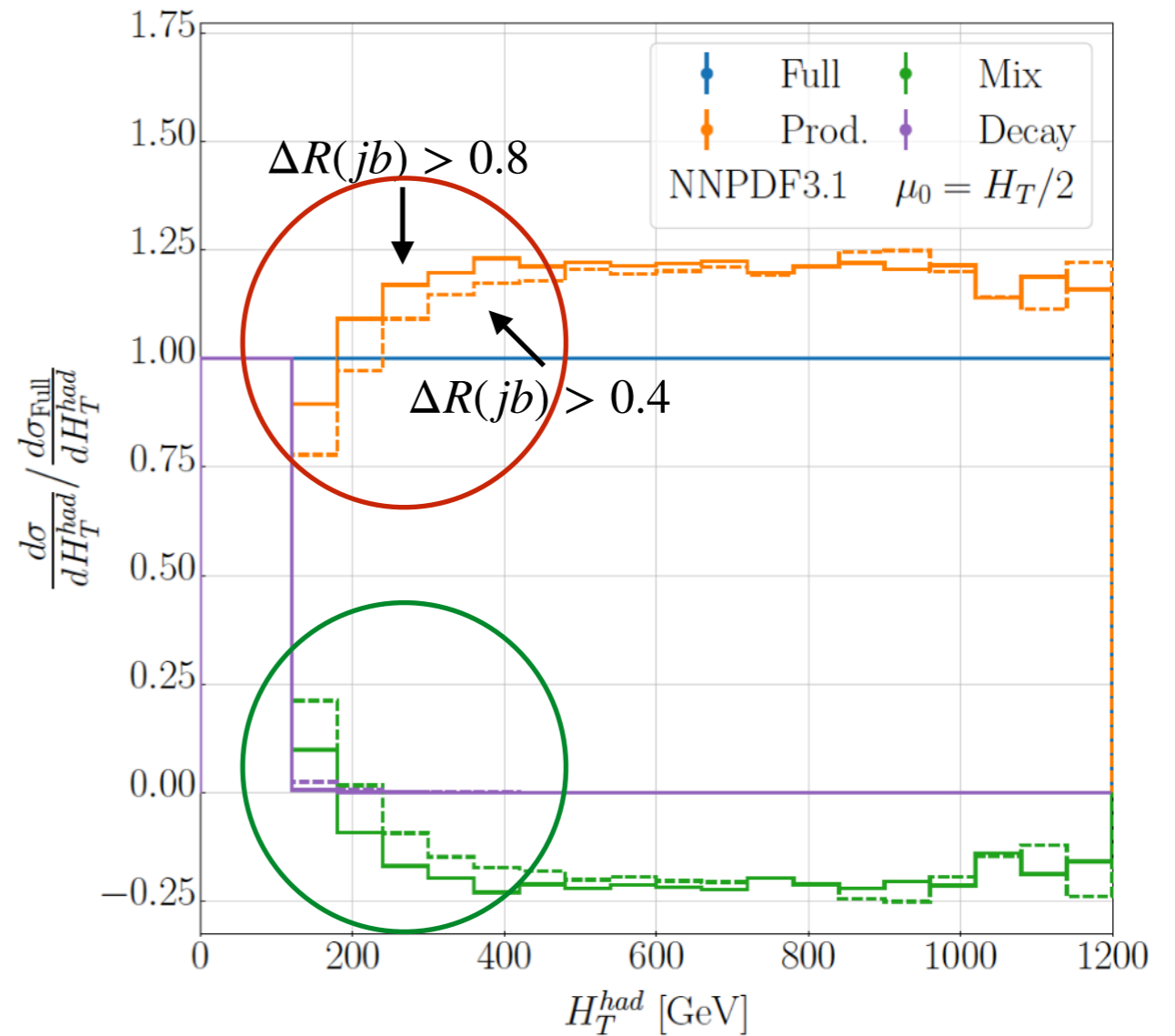
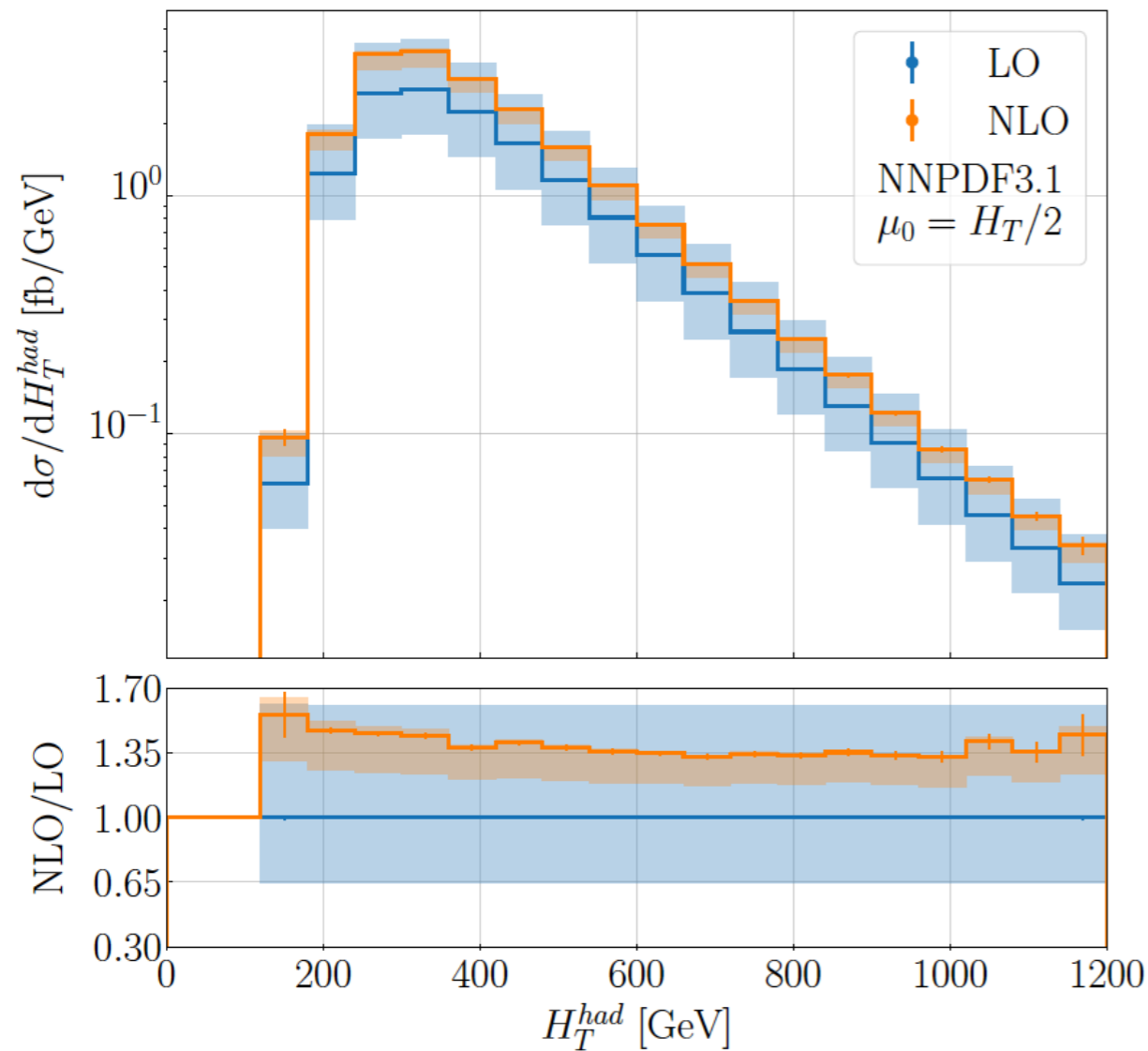
Modelling	σ^{LO} [fb]	σ^{NLO} [fb]	$\frac{\sigma_i^{\text{LO}}}{\sigma_{\text{NWA}}^{\text{LO}}}$	$\frac{\sigma_i^{\text{NLO}}}{\sigma_{\text{NWA}}^{\text{NLO}}}$
NWA _{full}	1074.5(3) ^{+60%} _{-35%}	1460(1) ^{+1%} _{-13%}	1.00	1.00
Prod	983.1(3) ^{+60%} _{-35%}	1662(1) ^{+11%} _{-18%}	0.91	1.14
Mix	89.42(3)	-205(1)	0.083	-0.14
Decay	1.909(1)	2.436(6)	0.002	0.002
NWA _{LOdec}	-	1390(2) ^{+11%} _{-18%}	-	0.95

- *Prod-Mix* interplay varies when jet radiation off top quarks is less suppressed
- $|\text{NWA}_{\text{full}} - \text{NWA}_{\text{LOdec}}| \sim 5\%$ for $\Delta R(jb) > 0.4$

$t\bar{t}jj$: differential cross sections

$$H_T^{had} = \sum_{i=1}^2 p_{Tj_i} + \sum_{i=1}^2 p_{Tb_i}$$

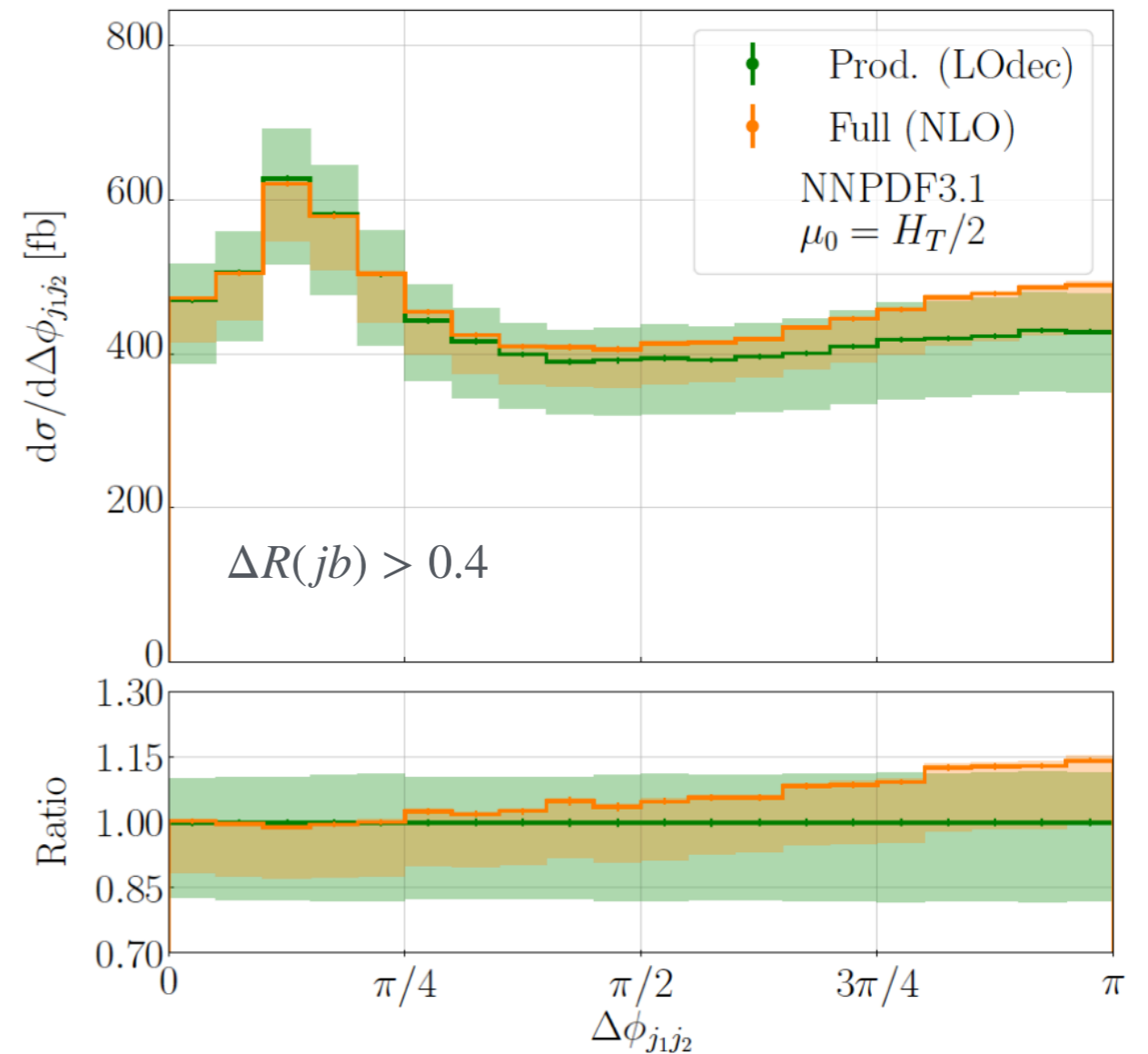
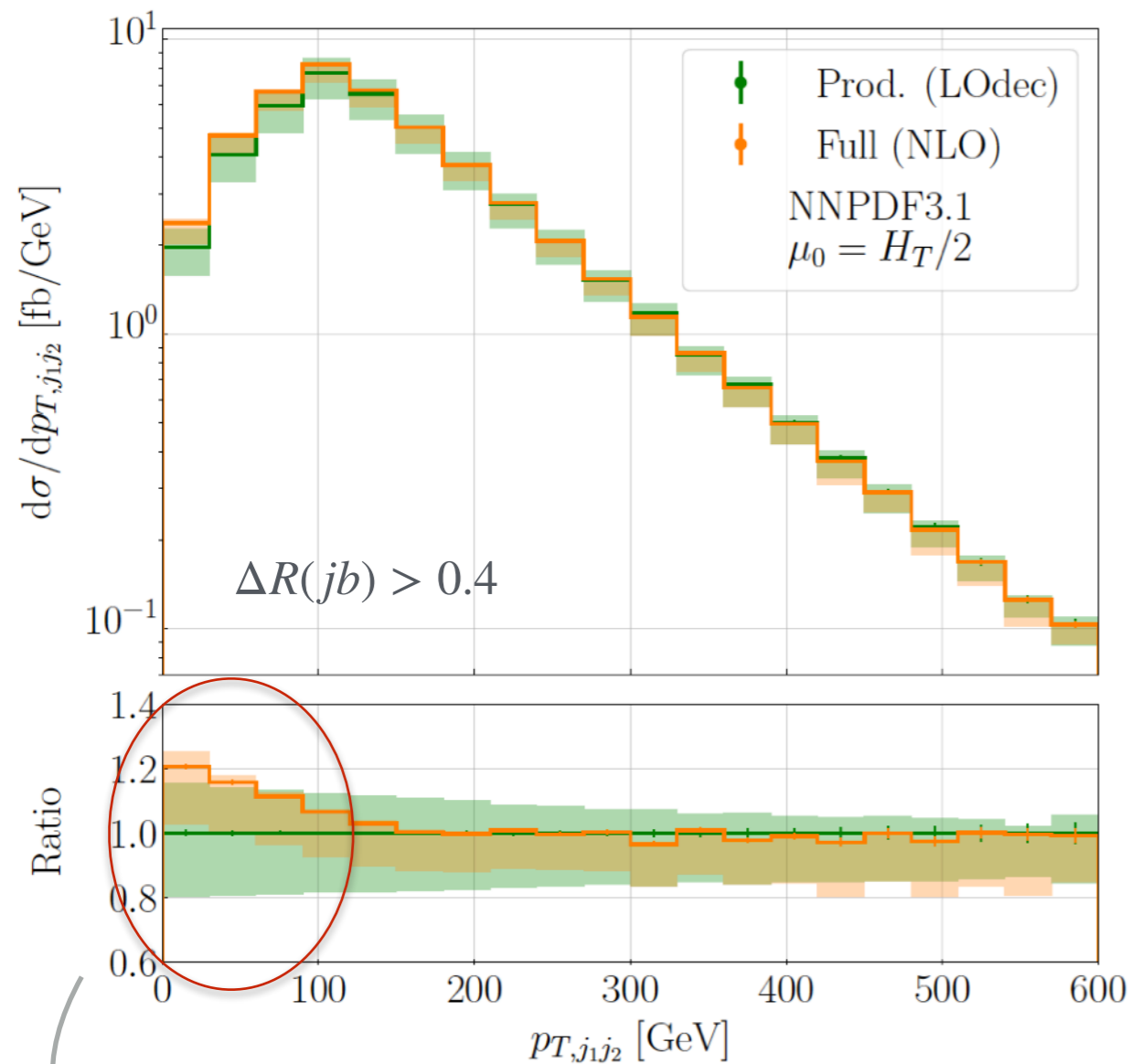
[GB, Lupattelli, Stremmer and Worek, [Phys. Rev. D \(107\) 2023, 114027](#)]



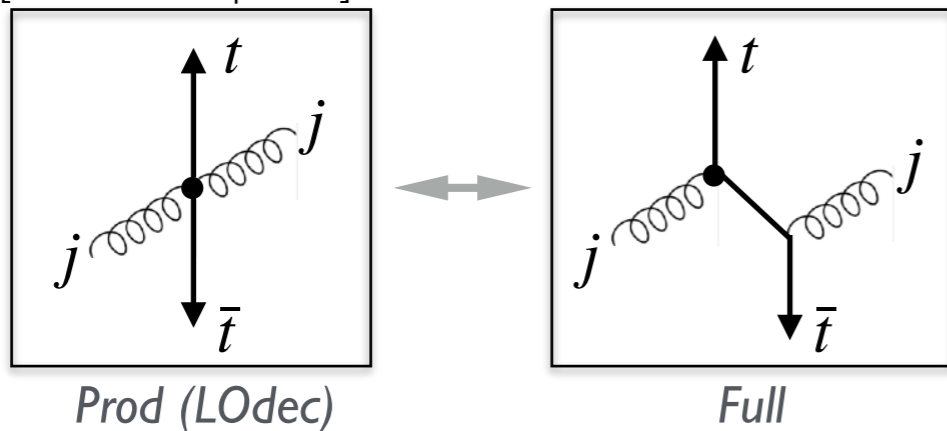
- Sensitivity to different $\Delta R(jb)$ cuts enhanced around the bulk

$t\bar{t}jj$: differential cross sections

[GB, Lupattelli, Stremmer and Worek, [Phys. Rev. D \(107\) 2023, 114027](#)]



[Transverse plane]



- Shape distortions up to 15% - 20 % in both dimensional and dimensionless observables

Lessons from $t\bar{t} + 1 \text{ jet}$

[Melnikov, Scharf and Schulze, [Phys.Rev.D 85 \(2012\) 054002](#)]

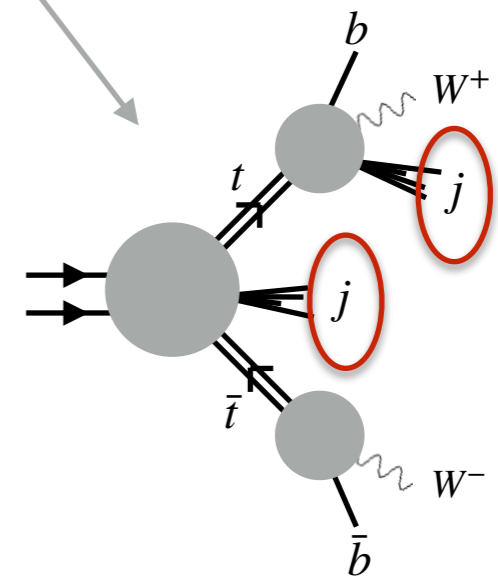
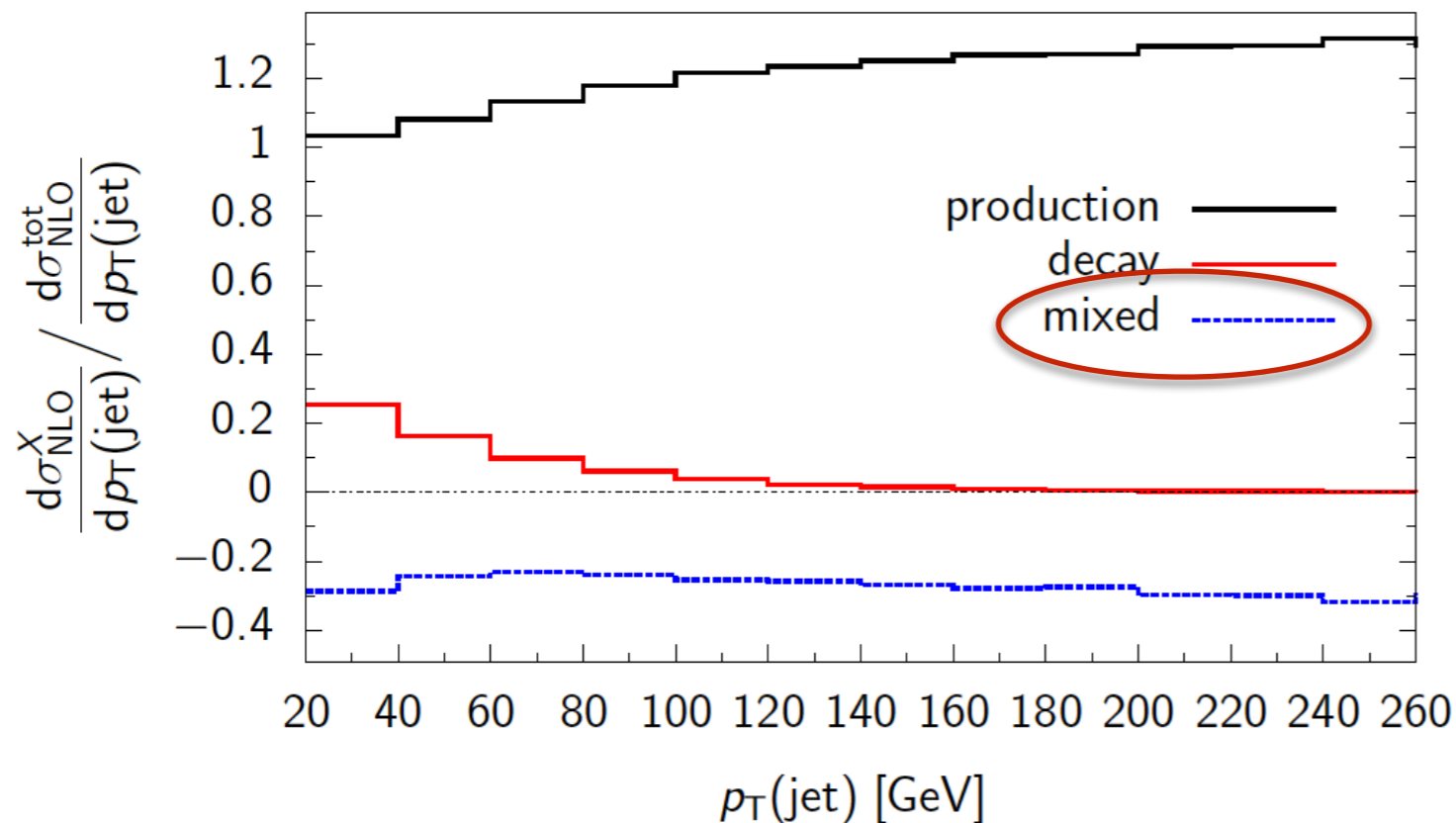
$$pp \rightarrow t\bar{t}j \rightarrow b\bar{b} \ell \nu_\ell jjj$$

$$\sqrt{s} = 7 \text{ TeV}$$

$$\sigma_{\text{LO}} = 316.9 \text{ (Pr)} + 33.4 \text{ (Dec)} = 350.3 \text{ fb,}$$

$$\sigma_{\text{NLO}} = 323 \text{ (Pr)} + 40.5 \text{ (Dec)} - 75.5 \text{ (Mix)} = 288 \text{ fb.}$$

-25% of σ_{NLO}

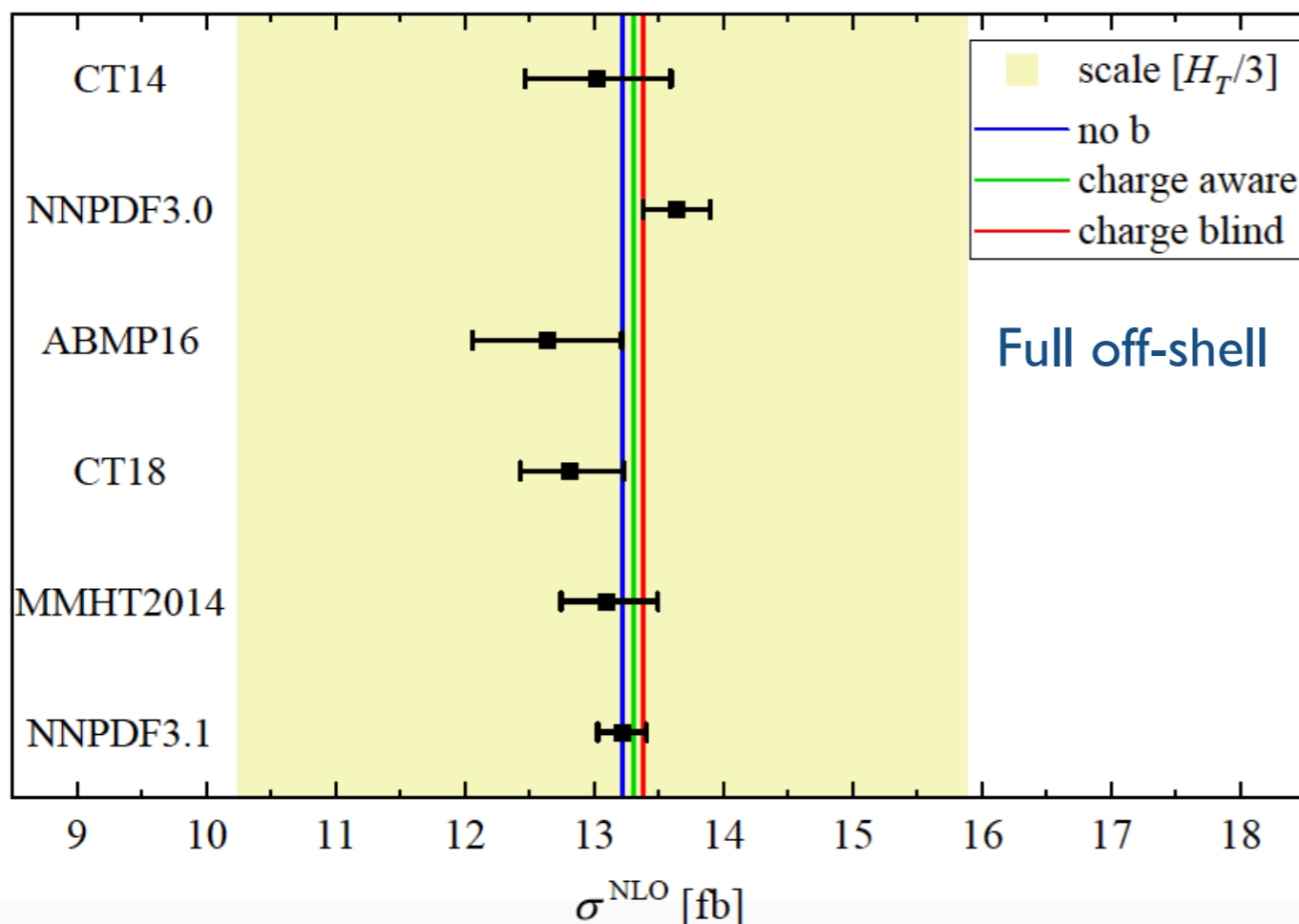


II. $t\bar{t}b\bar{b}$: off-shell effects and prompt b -jet categorisation

Off-shell $t\bar{t}b\bar{b}$: theoretical uncertainties at NLO

[GB, Bi, Hartanto, Kraus, Lupattelli and Worek, [JHEP 08 \(2021\) 008](#)]

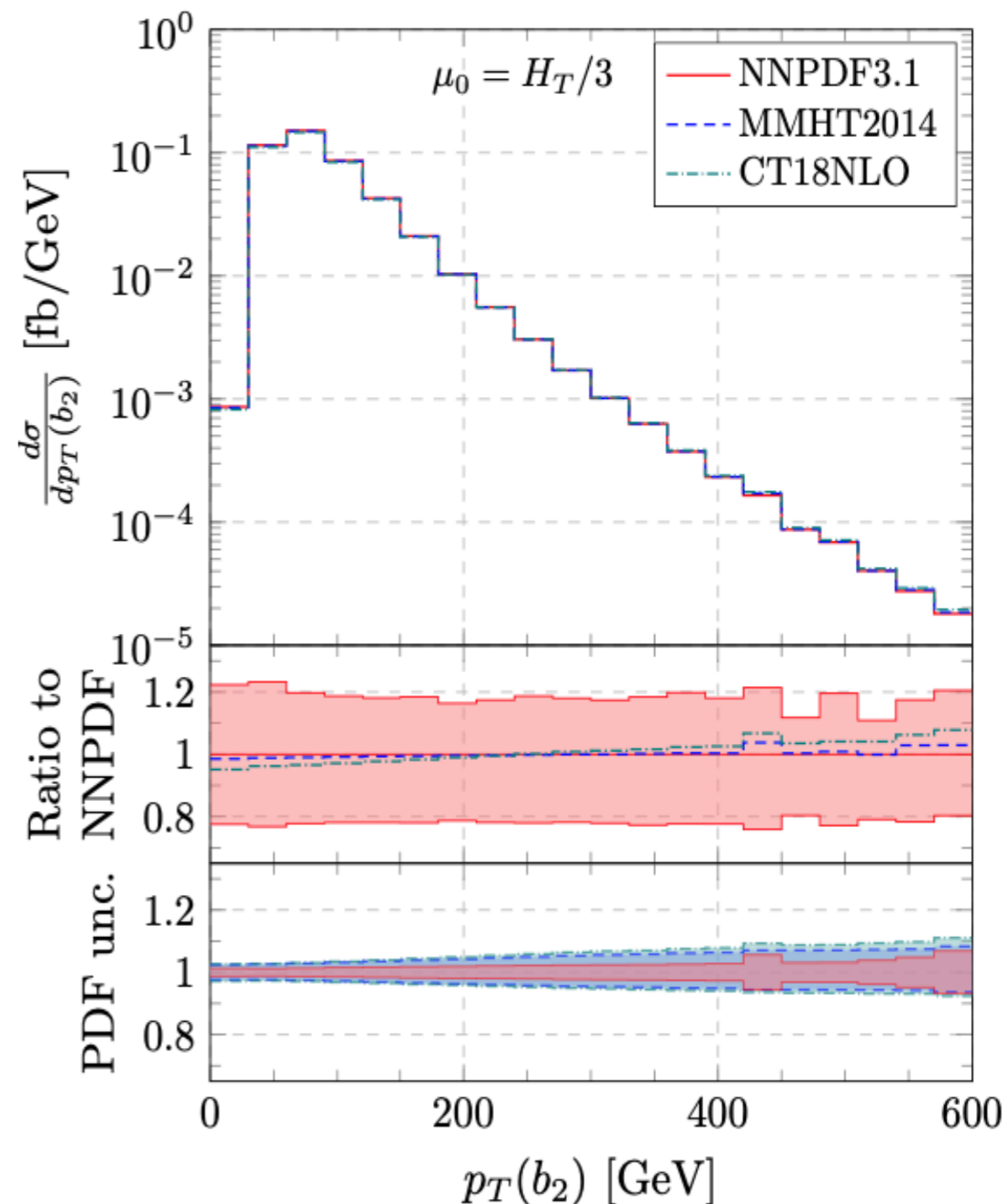
• $pp \rightarrow e^+ \nu_e \mu^- \bar{\nu}_\mu b\bar{b}b\bar{b}$ 13 TeV



NLO theory uncertainties

- Scale : $\mathcal{O}(20\%)$ - PDF : $\mathcal{O}(1\% - 3\%)$

↪ Relative sizes vary at high-energy tails



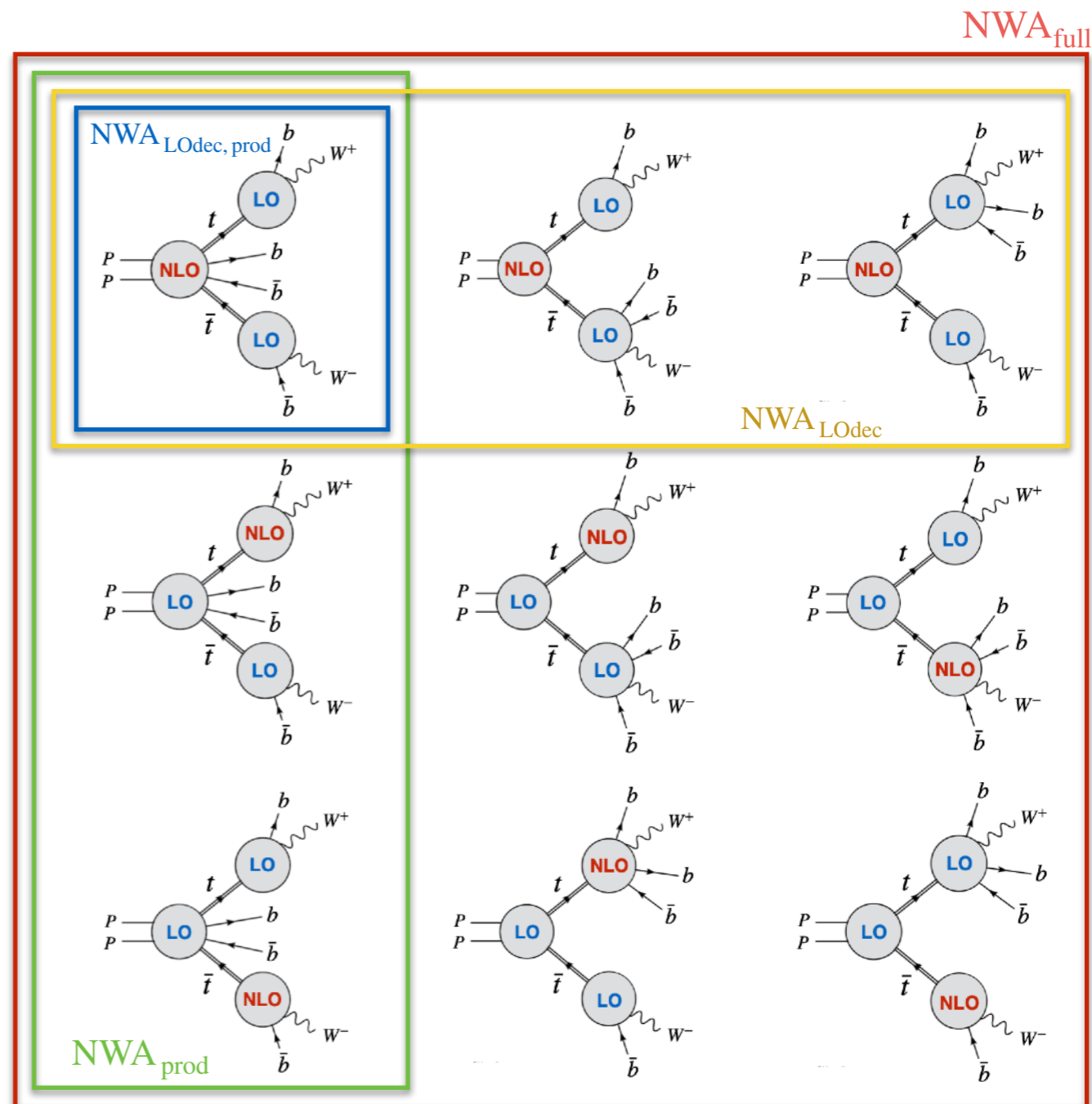
$t\bar{t}b\bar{b}$: comparing modelling approaches

[GB, Bi, Hartanto, Kraus, Lupattelli and Worek, [Phys. Rev. D \(107\) 2023](#)]

- Full off-shell vs NWA

Modelling	σ^{NLO} [fb]	δ_{scale} [fb]	$\frac{\sigma^{\text{NLO}}}{\sigma^{\text{NWA}_{\text{full}}}} - 1$
Off-shell	13.22(2)	+2.65 (20%) -2.96 (22%)	+0.5%
NWA_{full}	13.16(1)	+2.61 (20%) -2.93 (22%)	-
NWA _{LOdec}	13.22(1)	+3.77 (29%) -3.31 (25%)	+0.5%
NWA _{prod}	13.01(1)	+2.58 (20%) -2.89 (22%)	-1.1%
NWA _{prod, exp}	12.25(1)	+2.87 (23%) -2.86 (23%)	-6.9%
NWA_{prod, LOdec}	13.11(1)	+3.74 (29%) -3.28 (25%)	-0.4%

- NWA cross sections based on different levels of accuracy in top decay modelling



$t\bar{t}b\bar{b}$: comparing modelling approaches

[GB, Bi, Hartanto, Kraus, Lupattelli and Worek, [Phys. Rev. D \(107\) 2023](#)]

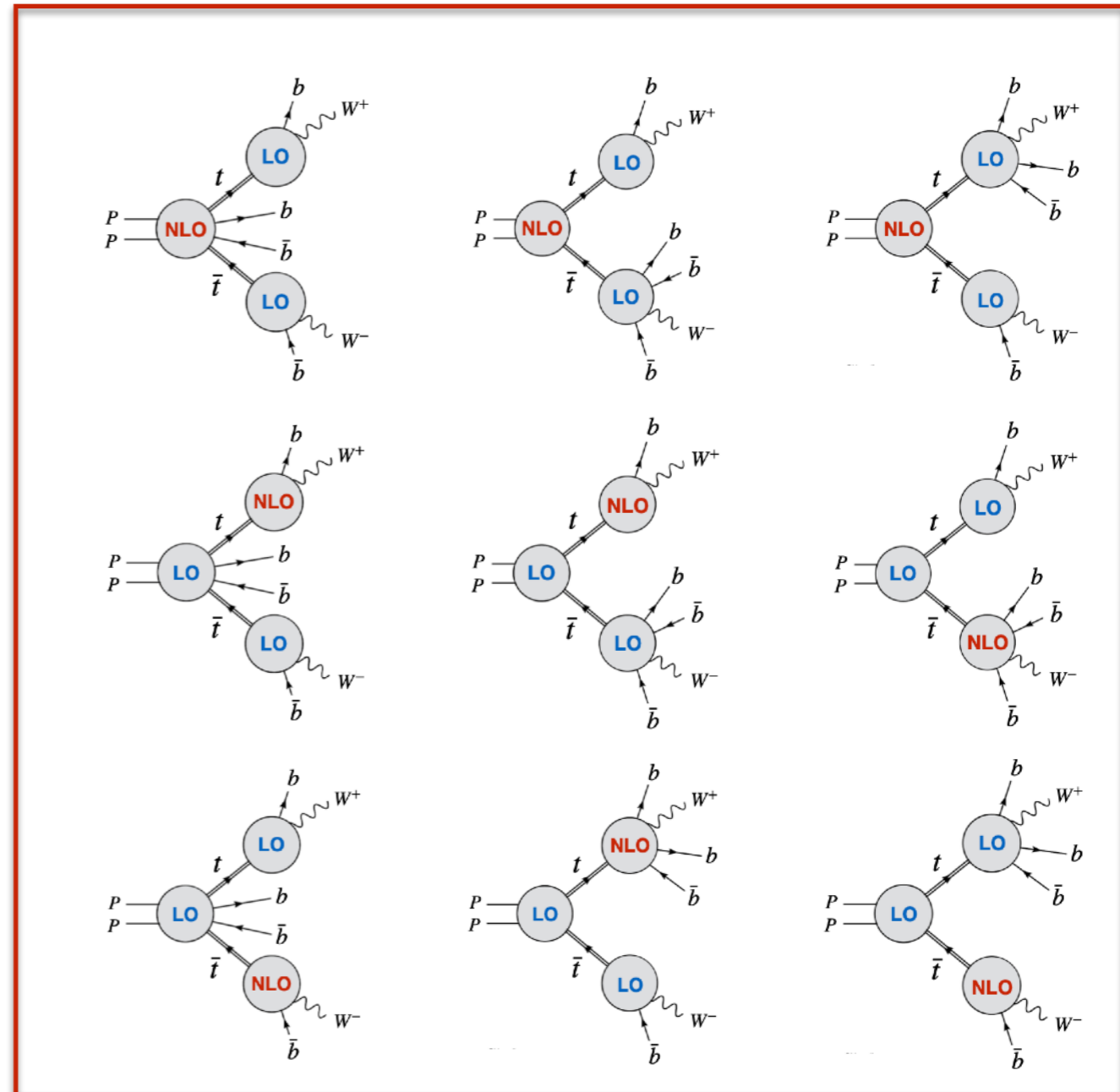
- **Full off-shell vs NWA**

Modelling	σ^{NLO} [fb]	δ_{scale} [fb]	$\frac{\sigma^{\text{NLO}}}{\sigma_{\text{NWA}_{\text{full}}^{\text{NLO}}}} - 1$
Off-shell	13.22(2)	+2.65 (20%) -2.96 (22%)	+0.5%
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NWA _{prod, LOdec}	13.11(1)	+3.74 (29%) -3.28 (25%)	-0.4%

- Off-shell vs NWA_{full} :

↪ Off-shell effects: **+0.5 %**

NWA_{full}



$t\bar{t}b\bar{b}$: comparing modelling approaches

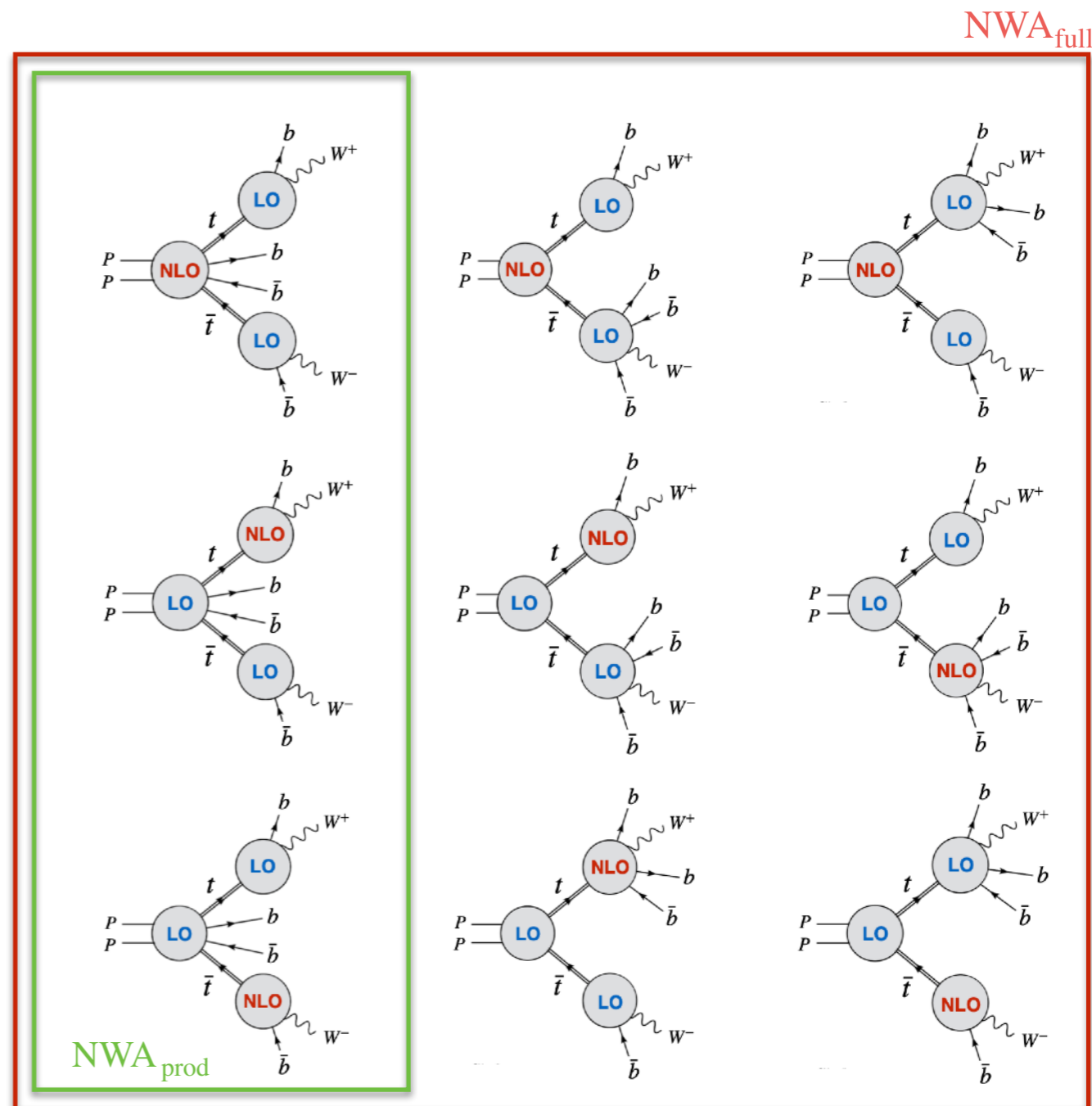
[GB, Bi, Hartanto, Kraus, Lupattelli and Worek, [Phys. Rev. D \(107\) 2023](#)]

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NWA _{prod, LOdec}	13.11(1)	+3.74 (29%) -3.28 (25%)	-0.4%

- NWA_{full} vs NWA_{prod} :

↪ Impact of $t \rightarrow Wb\bar{b}$ decays: +1%



$t\bar{t}b\bar{b}$: comparing modelling approaches

[GB, Bi, Hartanto, Kraus, Lupattelli and Worek, [Phys. Rev. D \(107\) 2023](#)]

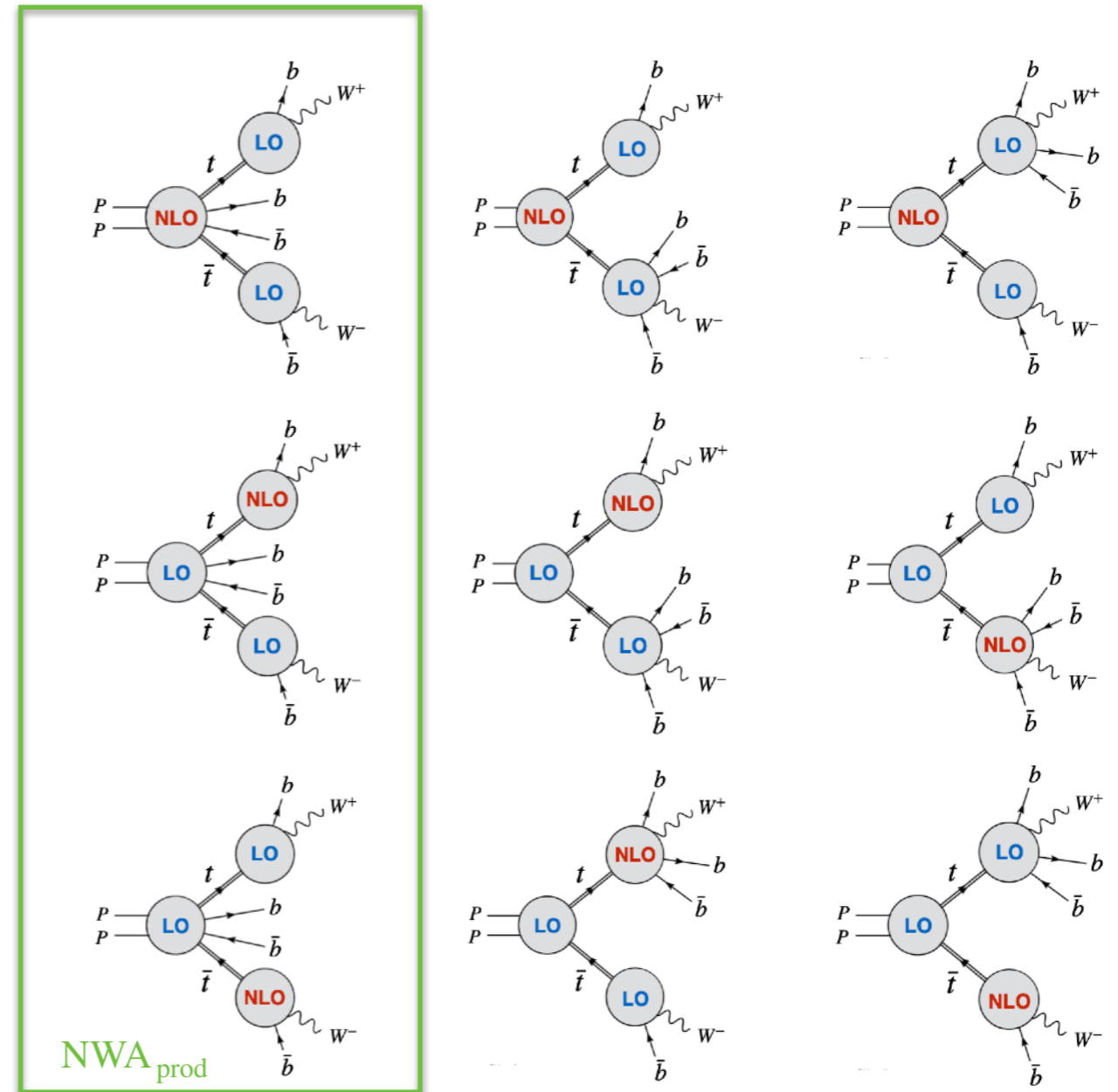
- Full off-shell vs NWA

Modelling	σ^{NLO} [fb]	δ_{scale} [fb]	$\frac{\sigma^{\text{NLO}}}{\sigma^{\text{NWA}_{\text{full}}}} - 1$
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NWA _{prod,LOdec}	13.11(1)	+3.74 (29%) -3.28 (25%)	-0.4%

- NWA_{prod} vs NWA_{prod,exp} :

↪ Impact of Γ_t^{NLO} expansion: +6%

$\mathcal{O}(\alpha_s^2)$ effects (within NLO scale uncertainty)



$t\bar{t}b\bar{b}$: comparing modelling approaches

[GB, Bi, Hartanto, Kraus, Lupattelli and Worek, [Phys. Rev. D \(107\) 2023](#)]

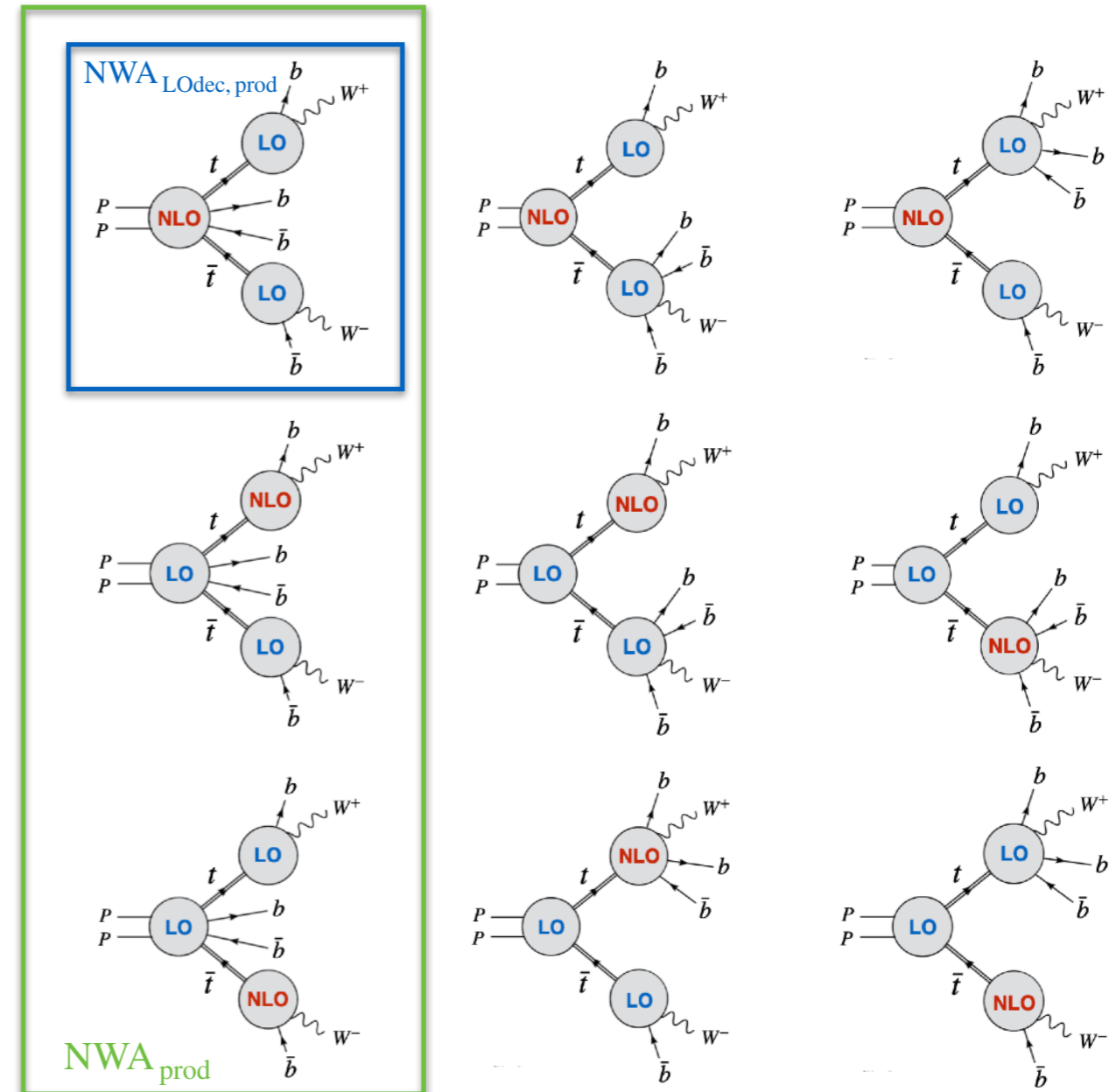
- Full off-shell vs NWA**

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NWA_{prod,LOdec}	13.11(1)	+3.74 (29%) -3.28 (25%)	-0.4%

- NWA_{prod,exp} vs NWA_{prod,LOdec} :

↪ **Genuine QCD corrections to top decays: -7%**

Accidental cancellations with $\mathcal{O}(\alpha_s^2)$ effects

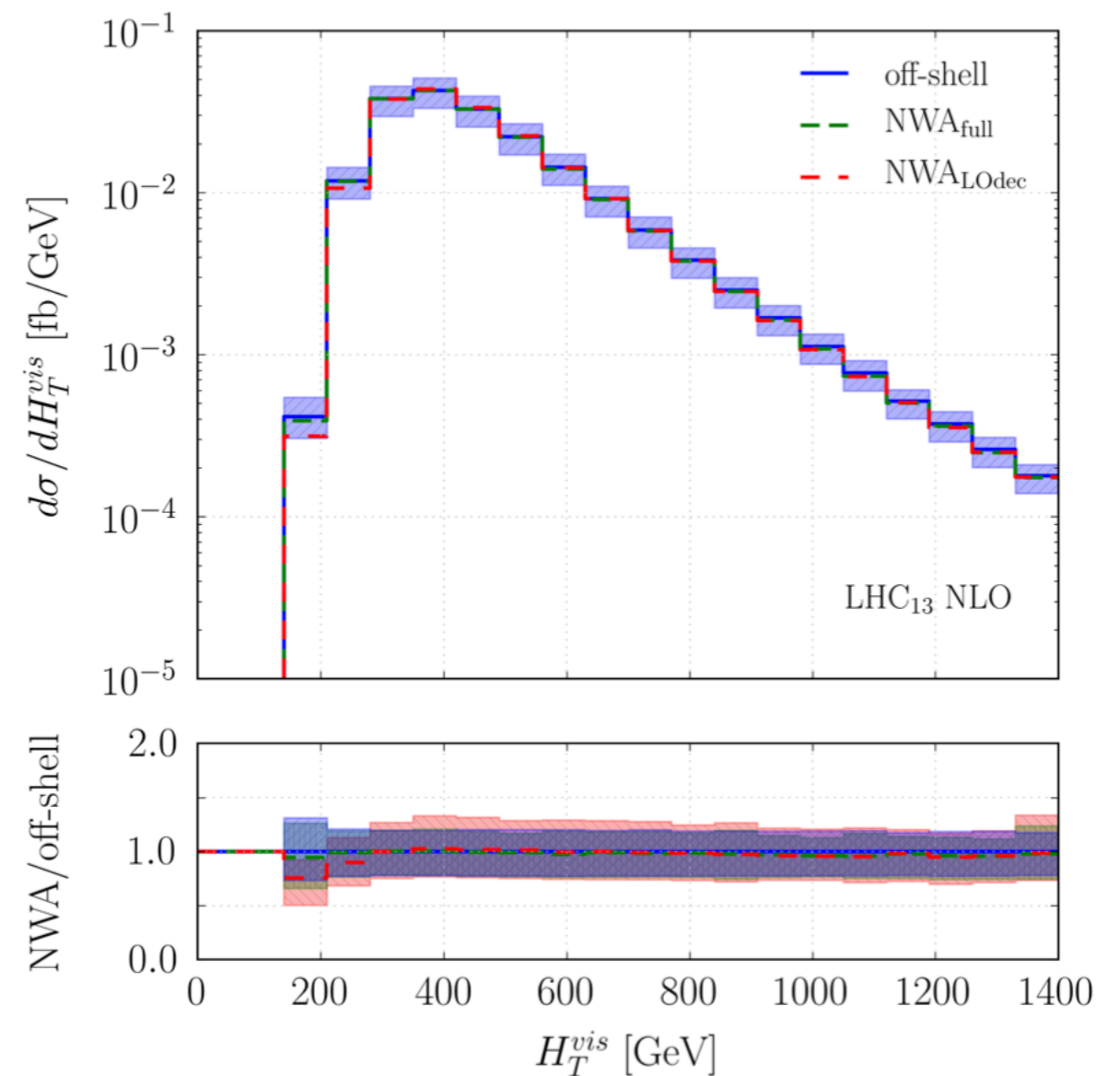
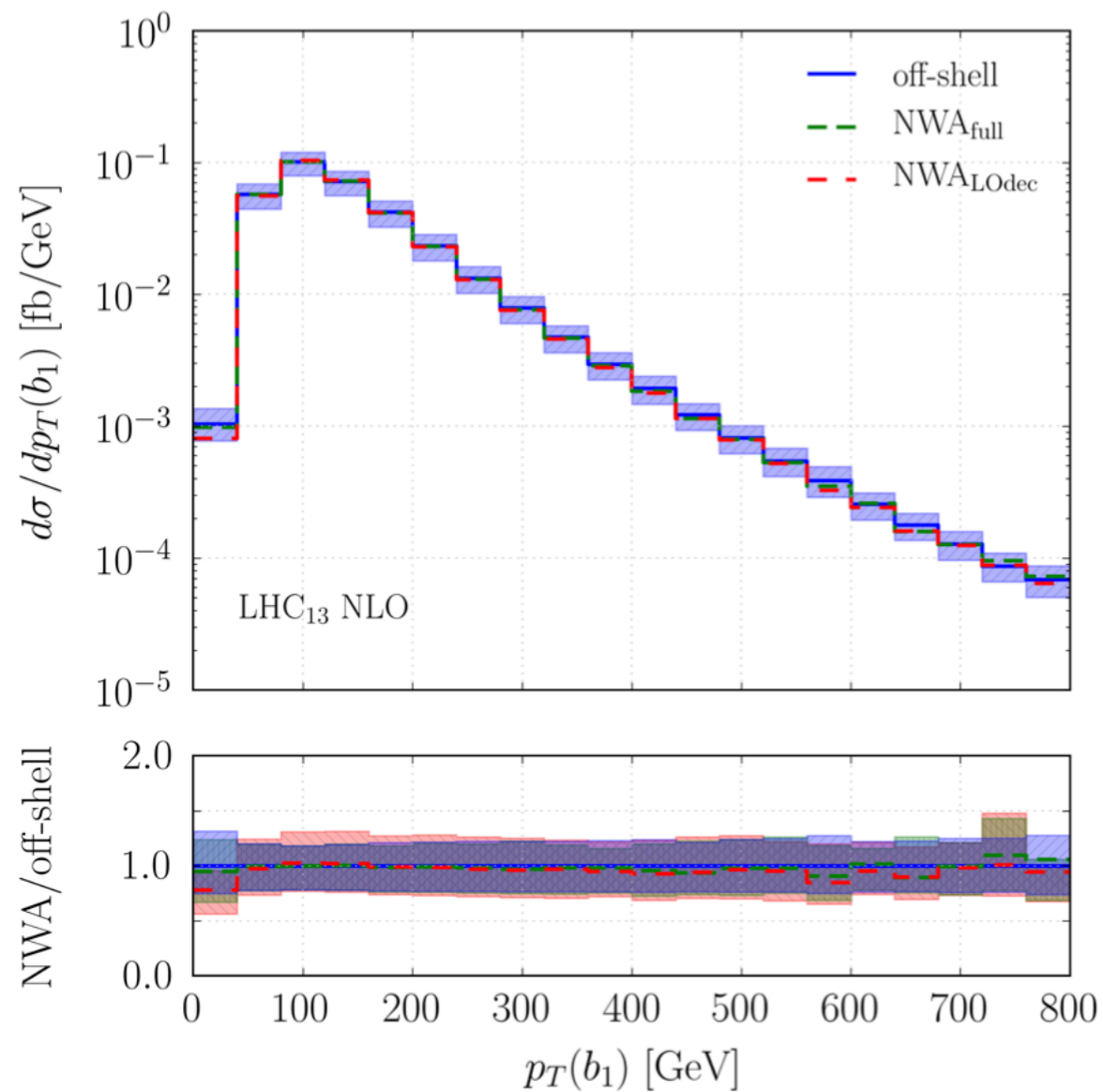


$t\bar{t}b\bar{b}$: impact of off-shell effects at differential level

[GB, Bi, Hartanto, Kraus, Lupattelli and Worek, [Phys. Rev. D \(107\) 2023](#)]

- Off-shell effects amount to few permille for most observables used in SM analyses

↪ Examples:

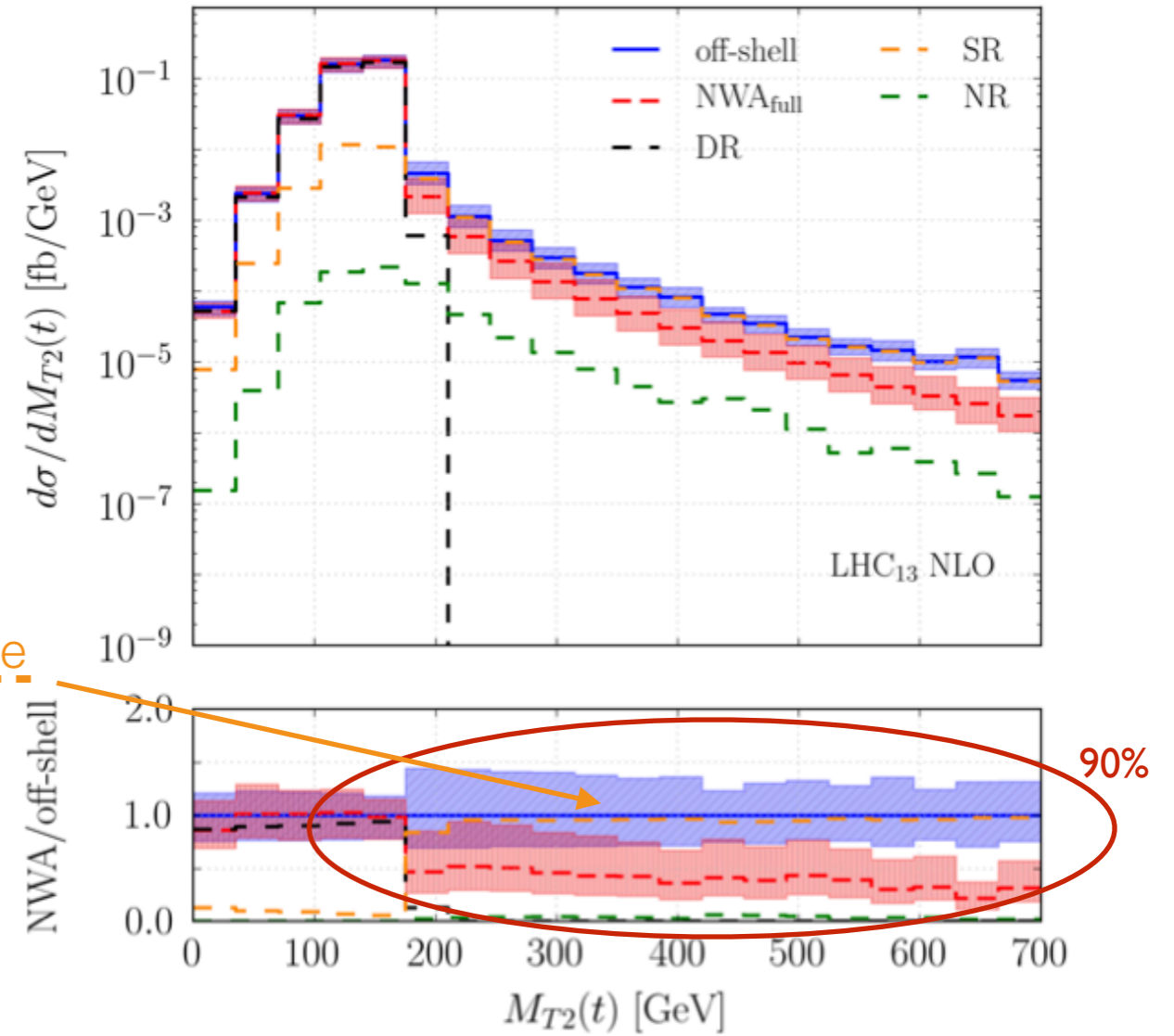
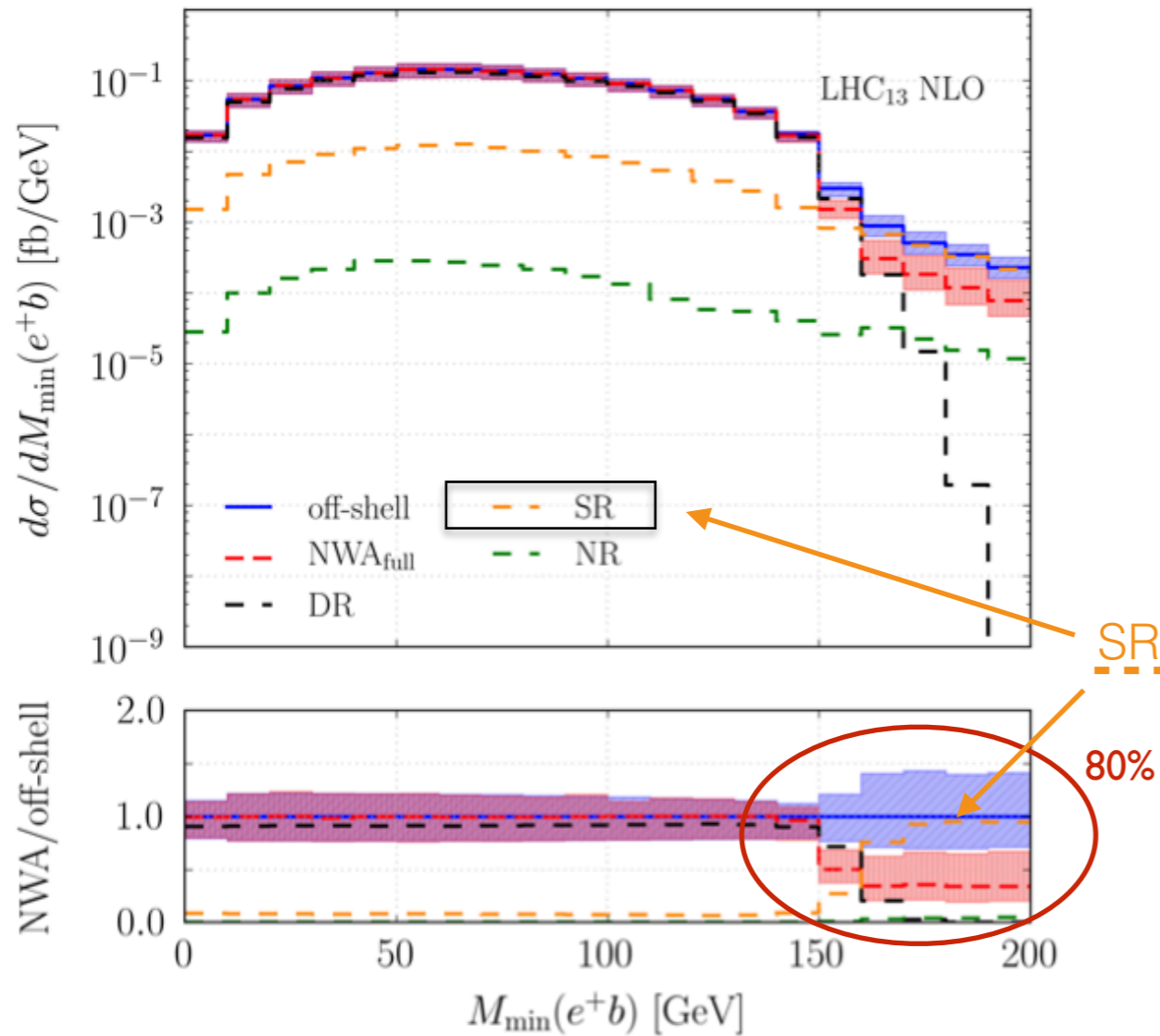


$$H_T^{vis} = \sum_{i=1}^4 p_T(b_i) + p_T(e^+) + p_T(\mu^-)$$

$t\bar{t}b\bar{b}$: anatomy of full off-shell effects at differential level

[GB, Bi, Hartanto, Kraus, Lupattelli and Worek, [Phys. Rev. D \(107\) 2023](#)]

- **Threshold observables** (mainly used in BSM studies) are more sensitive:



$$\text{LO}_{\text{NWA}} \rightarrow M_{\min}(e^+b) < \sqrt{m_t^2 - m_W^2} \approx 153 \text{ GeV}$$

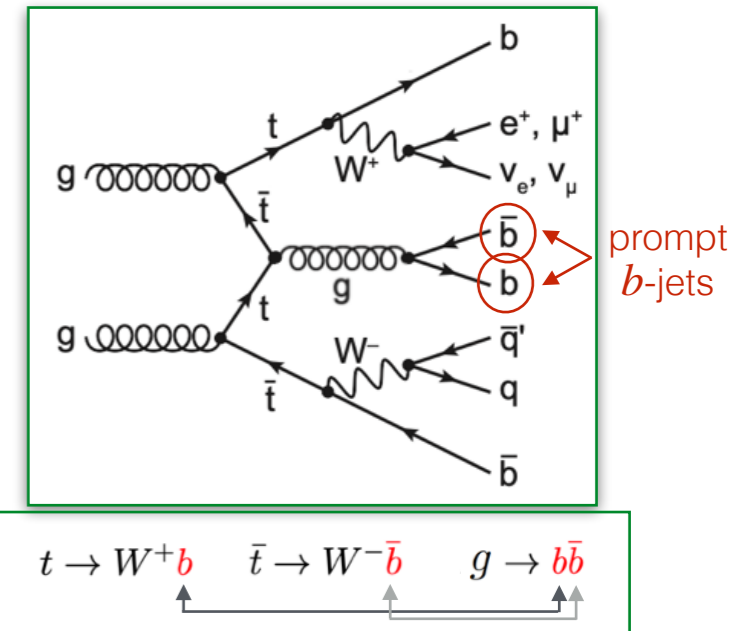
$$M_{T2}(t) = \min_{\sum p_T^{\nu_i} = p_T^{\text{miss}}} [\max\{M_T^2(p_T(e^+ X_t), p_T(\nu_1)), M_T^2(p_T(\mu^- X_{\bar{t}}), p_T(\nu_2))\}]$$

$t\bar{t}b\bar{b}$: prompt b -jet identification

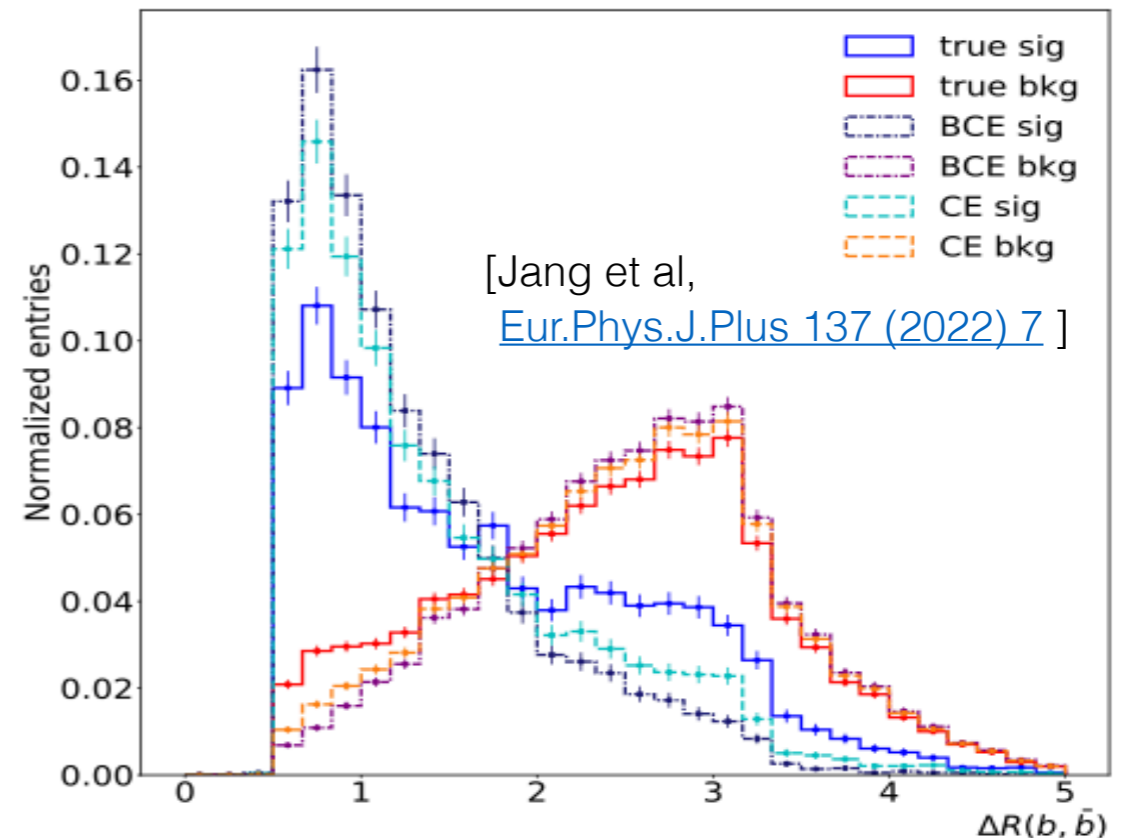
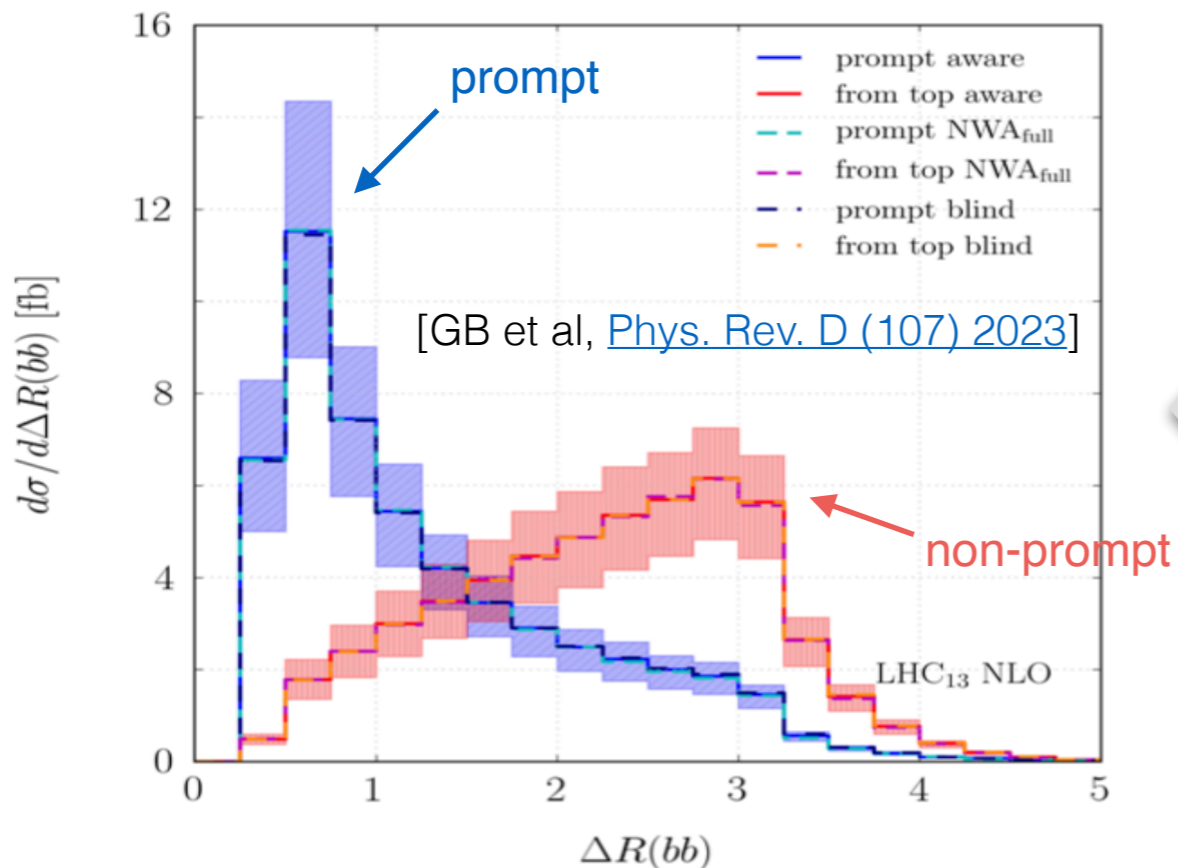
- Labelling prompt b -jets in $t\bar{t}b\bar{b}$ is not free of ambiguities in a full calculation (due to combinatorial background and quantum interference)

- **Kinematic-based prescription**: determine prompt b -jets according to a **minimum principle** for Q :

$$Q = |M(t) - m_t| \times |M(\bar{t}) - m_t| \times |M^{\text{prompt}}(bb)|$$



- Results consistent with DNN-based studies:



- We have examined some recent developments concerning matrix element description of $t\bar{t}b\bar{b}$ and $t\bar{t}jj$ @ NLO QCD

$t\bar{t}jj$ dilepton

- Jet radiation included in production and decays
- “Mix” contribution impacts NLO normalisation

$$\boxed{\Delta R(jb) > 0.8} \rightarrow \sigma_{\text{NLO}} = 1462 \text{ (Prod)} - 236 \text{ (Mix)} + 0.2 \text{ (Dec)} = 1225 \text{ fb}$$

-19% of σ_{NLO}

$$\boxed{\Delta R(jb) > 0.4} \rightarrow \sigma_{\text{NLO}} = 1662 \text{ (Prod)} - 205 \text{ (Mix)} + 2.4 \text{ (Dec)} = 1460 \text{ fb}$$

-14% of σ_{NLO}

$t\bar{t}b\bar{b}$ dilepton

- Off-shell effects small for most distributions used in SM analyses: $\mathcal{O}(0.5\%)$
- Impact of QCD corrections to top quark decays: -7%
- Impact of $\mathcal{O}(\alpha_s^2)$ effects in $\sigma_{\text{NWA}}^{\text{NLO}}$ and $\sigma_{\text{off-shell}}^{\text{NLO}}$: 6% (within the scale uncertainties)

Backup slides

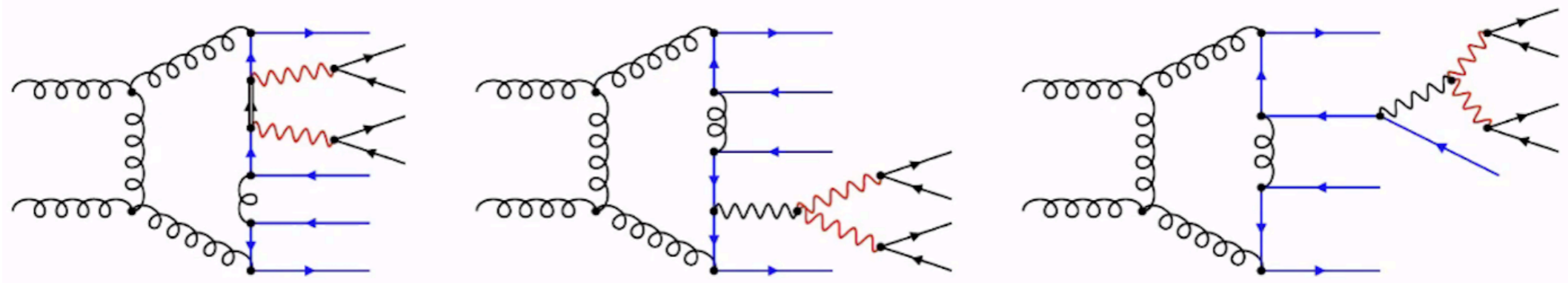
Technical aspects of off-shell $t\bar{t}b\bar{b}$ at NLO

Full off-shell calculations for multi-leg processes can be very challenging

One-loop correction type	Number of Feynman diagrams
Self-energy	93452
Vertex	88164
Box-type	49000
Pentagon-type	25876
Hexagon-type	11372
Heptagon-type	3328
Octagon-type	336
Total number [gg channel]	271528

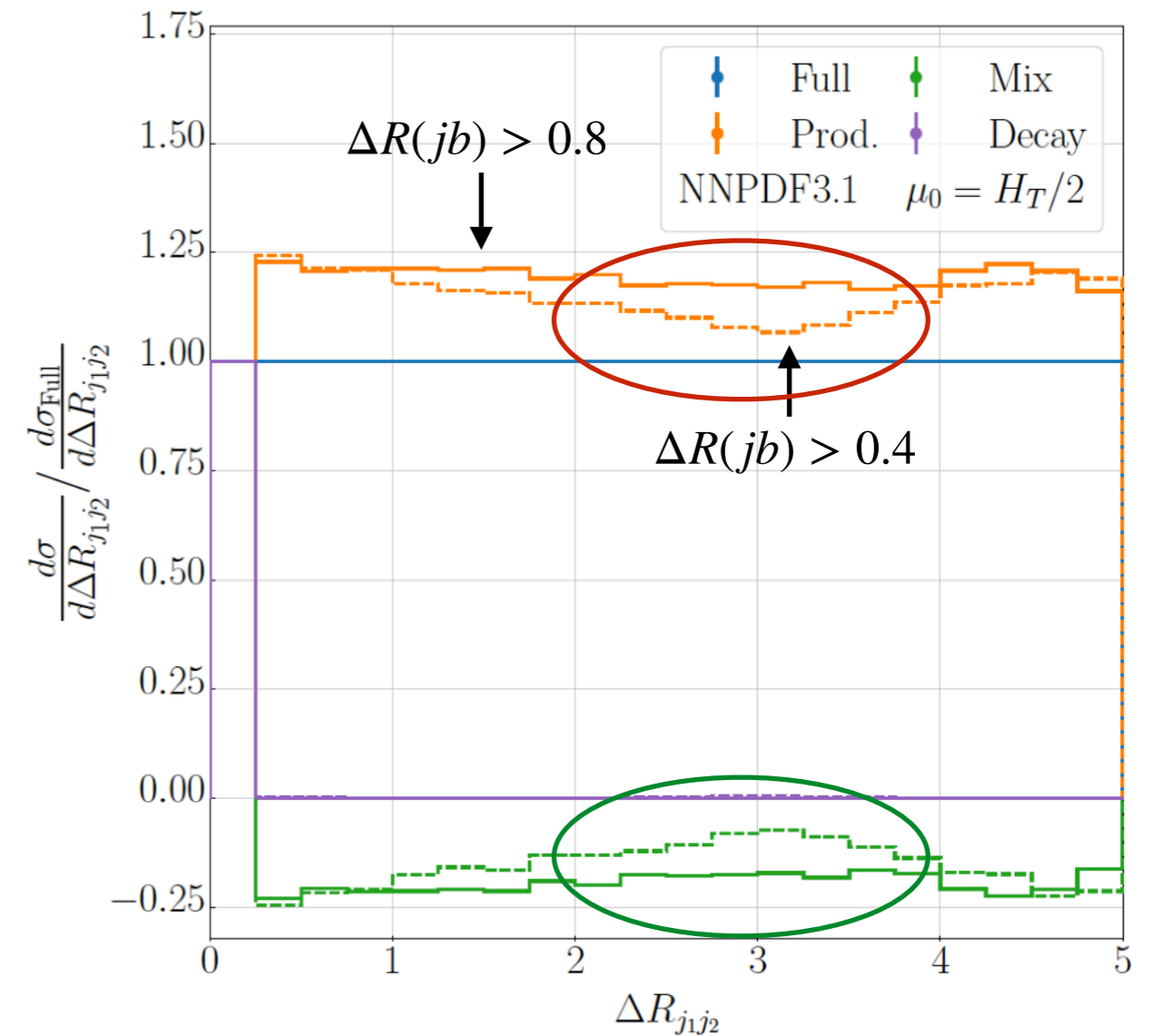
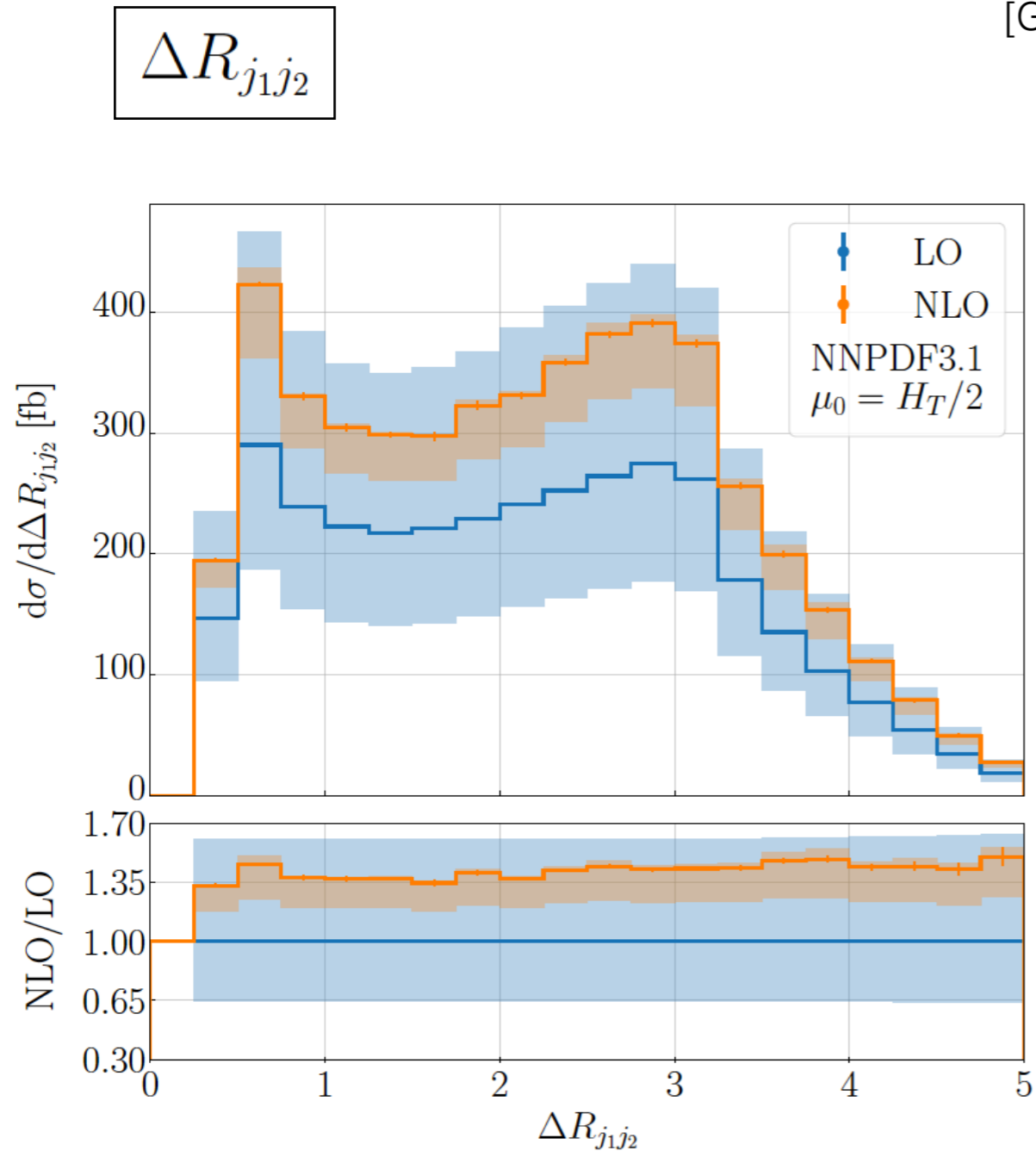
Example: $gg \rightarrow e^+ \nu_e \mu^- \bar{\nu}_\mu b\bar{b} b\bar{b}$

Partonic Subprocess	Number of Feynman diagrams	Number of CS Dipoles	Number of NS Subtractions
$gg \rightarrow e^+ \nu_e \mu^- \bar{\nu}_\mu b\bar{b} b\bar{b} g$	41364	90	18
$q\bar{q} \rightarrow e^+ \nu_e \mu^- \bar{\nu}_\mu b\bar{b} b\bar{b} g$	9576	50	10
$gq \rightarrow e^+ \nu_e \mu^- \bar{\nu}_\mu b\bar{b} b\bar{b} q$	9576	50	10
$g\bar{q} \rightarrow e^+ \nu_e \mu^- \bar{\nu}_\mu b\bar{b} b\bar{b} \bar{q}$	9576	50	10



$t\bar{t}jj$: differential cross sections

[GB, Lupattelli, Stremmer and Worek, [Phys. Rev. D \(107\) 2023, 114027](#)]



- Sensitivity to $\Delta R(jb)$ cuts enhanced around $\Delta R(j_1 j_2) = 3$

$t\bar{t}$ + jets: cross section ratios

[GB, Lupattelli, Stremmer and Worek, [Phys. Rev. D \(107\) 2023, 114027](#)]

Fiducial cross section ratios

$$\Delta R(jb) > 0.8$$

\mathcal{R}_n	\mathcal{R}^{LO}	\mathcal{R}^{NLO}	$\mathcal{R}_{\text{exp}}^{\text{NLO}}$
$\mathcal{R}_1 = \sigma_{t\bar{t}j} / \sigma_{t\bar{t}}$	$0.3686^{+12\%}_{-10\%}$	$0.3546^{+0\%}_{-5\%}$	$0.3522^{+0\%}_{-3\%}$
$\mathcal{R}_2 = \sigma_{t\bar{t}jj} / \sigma_{t\bar{t}j}$	$0.2539^{+11\%}_{-9\%}$	$0.2660^{+0\%}_{-5\%}$	$0.2675^{+0\%}_{-2\%}$

$$\mathcal{R}_{\text{exp}}^{\text{NLO}} = \frac{\sigma_{t\bar{t}j(j)}^0}{\sigma_{t\bar{t}(j)}^0} \left(1 + \frac{\sigma_{t\bar{t}j(j)}^1}{\sigma_{t\bar{t}j(j)}^0} - \frac{\sigma_{t\bar{t}(j)}^1}{\sigma_{t\bar{t}(j)}^0} \right)$$

- NLO QCD corrections: $\mathcal{R}_1 \rightarrow -4\%$ $\mathcal{R}_2 \rightarrow +4\%$
- NLO uncertainties: Scale $\rightarrow \mathcal{O}(5\%)$ PDF $\rightarrow \mathcal{O}(0.5\%)$

$t\bar{t}b\bar{b}$ with dilepton signatures: full off-shell predictions

- Integrated fiducial cross sections

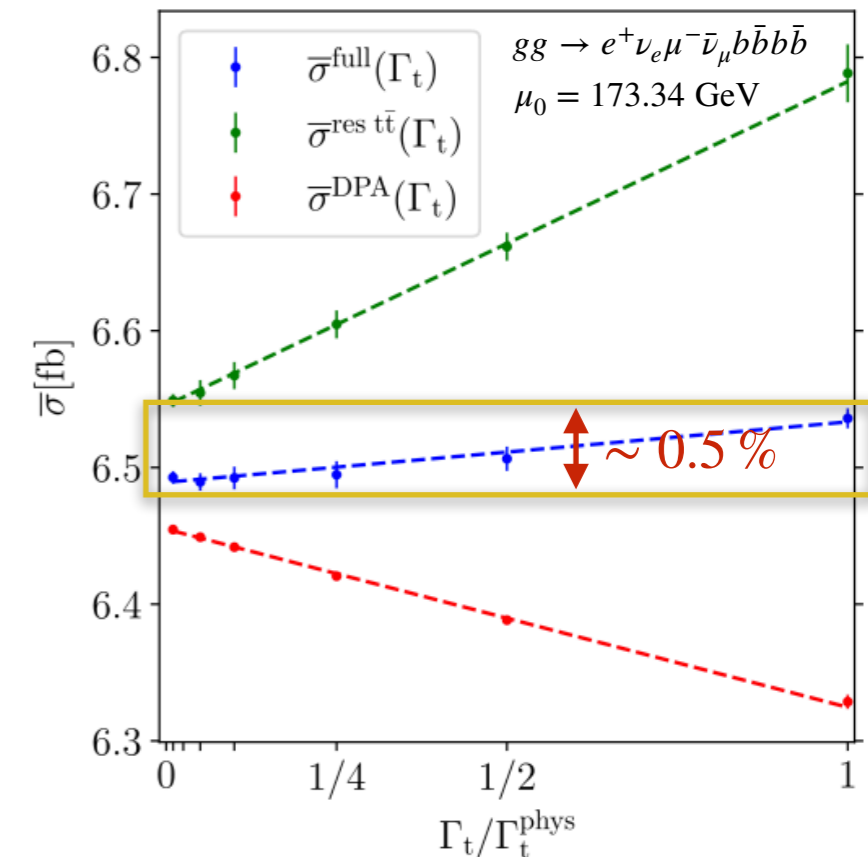
$$\sqrt{s} = 13 \text{ TeV}$$

Analysis cuts: $p_T(\ell) > 20 \text{ GeV}$, $p_T(b) > 25 \text{ GeV}$, $|y(\ell)| < 2.5$, $|y(b)| < 2.5$

[GB, Bi, Hartanto, Kraus, Lupattelli and Worek, [JHEP 08 \(2021\) 008](#)]

$p_T(b)$	σ^{LO} [fb]	δ_{scale}	σ^{NLO} [fb]	δ_{scale}	δ_{PDF}	$\mathcal{K} = \sigma^{\text{NLO}}/\sigma^{\text{LO}}$
$\mu_R = \mu_F = \mu_0 = H_T/3$ [NNPDF 3.1]						
25	6.813	+4.338 (64%) -2.481 (36%)	13.22	+2.66 (20%) -2.95 (22%)	+0.19 (1%) -0.19 (1%)	1.94
30	4.809	+3.062 (64%) -1.756 (37%)	9.09	+1.66 (18%) -1.98 (22%)	+0.16 (2%) -0.16 (2%)	1.89
35	3.431	+2.191 (64%) -1.256 (37%)	6.37	+1.07 (17%) -1.36 (21%)	+0.11 (2%) -0.11 (2%)	1.86
40	2.464	+1.582 (64%) -0.901 (37%)	4.51	+0.72 (16%) -0.95 (21%)	+0.09 (2%) -0.09 (2%)	1.83

[Denner, Lang, Pellen, [Phys. Rev. D 104 \(2021\), 056018](#)]

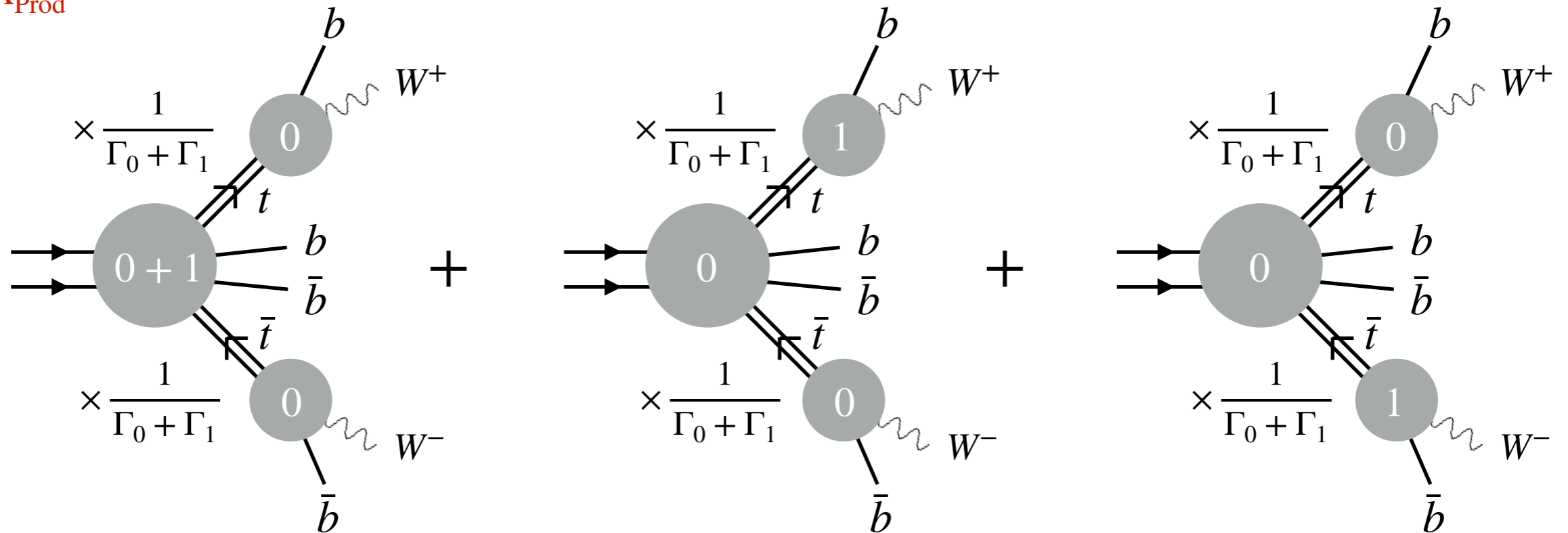


QCD corrections are large: $\mathcal{O}(80\%)$

Effects of Γ_t expansion in NWA

- "Unexpanded" cross section in NWA:

$$d\sigma_{\text{NWA}_{\text{Prod}}}^{\text{NLO}} \equiv$$



- $d\sigma_{\text{NWA}_{\text{Prod}}}^{\text{NLO}}$ yields the same Γ_t^{NLO} factors that are present in the full off-shell calculation
- Most suitable setup to assess the impact of genuine off-shell effects

Effects of Γ_t expansion in NWA

- “Unexpanded” cross section in NWA:

$$\Gamma_t^{NLO} = \Gamma_0 + \Gamma_1$$

\downarrow \downarrow
 $\mathcal{O}(\alpha_s^0)$ $\mathcal{O}(\alpha_s)$

$$\int d\sigma_{\text{NWA Prod}}^{NLO} =$$

The diagram illustrates the expansion of the NWA cross section into three terms. Each term features a central vertex with two incoming lines and four outgoing lines (two for b, \bar{b} and two for t, \bar{t}). The first term has a central vertex labeled $0+1$ and is multiplied by $\frac{1}{\Gamma_0 + \Gamma_1}$. It is connected to two vertices labeled 0 , one for W^+ production and one for W^- production. The second term has a central vertex labeled 0 and is multiplied by $\frac{1}{\Gamma_0 + \Gamma_1}$. The t and \bar{t} lines are connected to a vertex labeled 1 for W^+ production and a vertex labeled 0 for W^- production. The third term has a central vertex labeled 0 and is multiplied by $\frac{1}{\Gamma_0 + \Gamma_1}$. The t and \bar{t} lines are connected to a vertex labeled 0 for W^+ production and a vertex labeled 1 for W^- production.

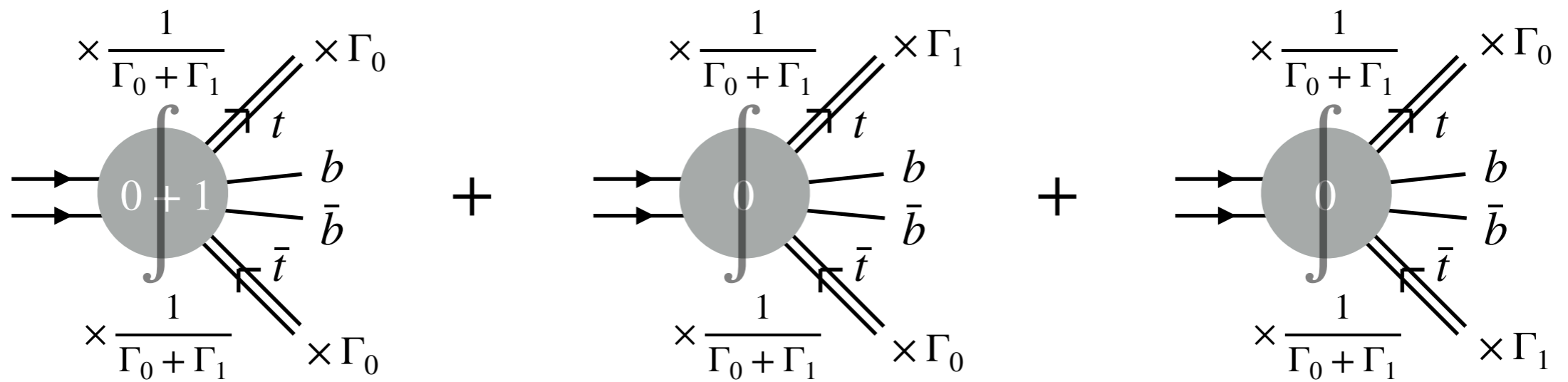
Effects of Γ_t expansion in NWA

- “Unexpanded” cross section in NWA:

$$\Gamma_t^{NLO} = \Gamma_0 + \Gamma_1$$

\downarrow \downarrow
 $\mathcal{O}(\alpha_s^0)$ $\mathcal{O}(\alpha_s)$

$$\int d\sigma_{\text{NWA Prod}}^{NLO} =$$



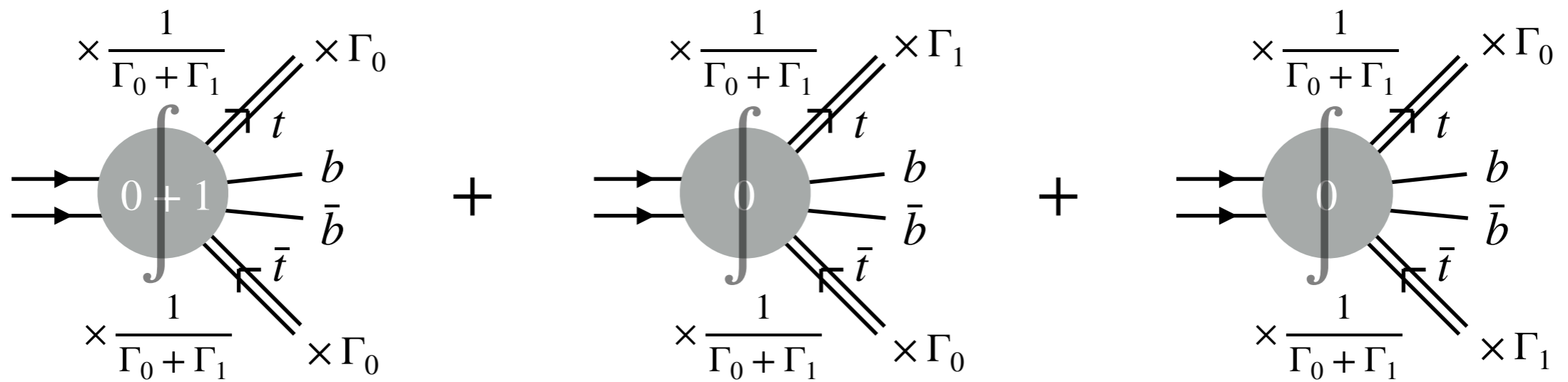
Effects of Γ_t expansion in NWA

- “Unexpanded” cross section in NWA:

$$\Gamma_t^{NLO} = \Gamma_0 + \Gamma_1$$

\downarrow \downarrow
 $\mathcal{O}(\alpha_s^0)$ $\mathcal{O}(\alpha_s)$

$$\int d\sigma_{\text{NWA Prod}}^{NLO} =$$



$$= \int d\sigma_{\text{NWA Prod}}^{NLO} + \mathcal{O}(\alpha_s^2) = \sigma_{t\bar{t}b\bar{b}, \text{incl.}}^{NLO} + \mathcal{O}(\alpha_s^2)$$

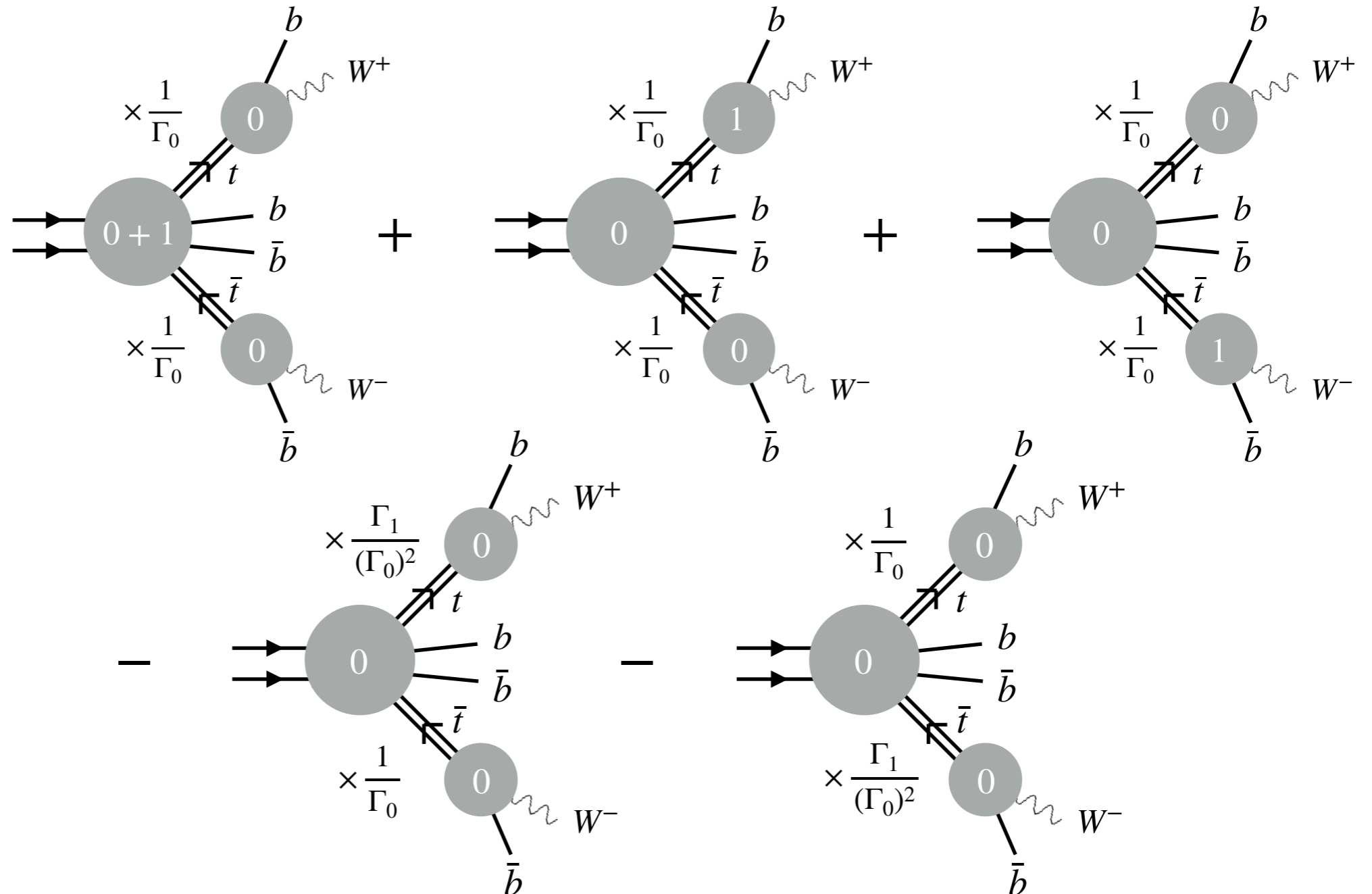
- Rigorous factorization at NLO spoiled by *formally suppressed $\mathcal{O}(\alpha_s^2)$ terms.*

Effects of Γ_t expansion in NWA

- "Expanded" cross section in NWA:

$$\frac{1}{\Gamma_0 + \Gamma_1} \sim \frac{1}{\Gamma_0} - \frac{\Gamma_1}{(\Gamma_0)^2} + \mathcal{O}(\alpha_s^2)$$

$d\sigma_{\text{NWA Prod, exp}}^{\text{NLO}} \equiv$



Effects of Γ_t expansion in NWA

- "Expanded" cross section in NWA:

$$\frac{1}{\Gamma_0 + \Gamma_1} \sim \frac{1}{\Gamma_0} - \frac{\Gamma_1}{(\Gamma_0)^2} + \mathcal{O}(\alpha_s^2)$$

$$\int d\sigma_{\text{NWA Prod, exp}}^{\text{NLO}} =$$

The diagrammatic expansion of the NWA cross section is shown as follows:

- Row 1 (Addition):
 - Diagram 1: Vertex $0+1$ (left), 0 (top), 0 (bottom). Factors: $\times \frac{1}{\Gamma_0}$ (top), $\times \frac{1}{\Gamma_0}$ (bottom).
 - Diagram 2: Vertex 0 (left), 1 (top), 0 (bottom). Factors: $\times \frac{1}{\Gamma_0}$ (top), $\times \frac{1}{\Gamma_0}$ (bottom).
 - Diagram 3: Vertex 0 (left), 0 (top), 1 (bottom). Factors: $\times \frac{1}{\Gamma_0}$ (top), $\times \frac{1}{\Gamma_0}$ (bottom).
- Row 2 (Subtraction):
 - Diagram 4: Vertex 0 (left), 0 (top), 0 (bottom). Factors: $\times \frac{\Gamma_1}{(\Gamma_0)^2}$ (top), $\times \frac{1}{\Gamma_0}$ (bottom).
 - Diagram 5: Vertex 0 (left), 0 (top), 0 (bottom). Factors: $\times \frac{1}{\Gamma_0}$ (top), $\times \frac{\Gamma_1}{(\Gamma_0)^2}$ (bottom).

Effects of Γ_t expansion in NWA

- "Expanded" cross section in NWA:

$$\frac{1}{\Gamma_0 + \Gamma_1} \sim \frac{1}{\Gamma_0} - \frac{\Gamma_1}{(\Gamma_0)^2} + \mathcal{O}(\alpha_s^2)$$

$$\int d\sigma_{\text{NWA Prod, exp}}^{NLO} =$$

Effects of Γ_t expansion in NWA

- "Expanded" cross section in NWA:

$$\frac{1}{\Gamma_0 + \Gamma_1} \sim \frac{1}{\Gamma_0} - \frac{\Gamma_1}{(\Gamma_0)^2} + \mathcal{O}(\alpha_s^2)$$

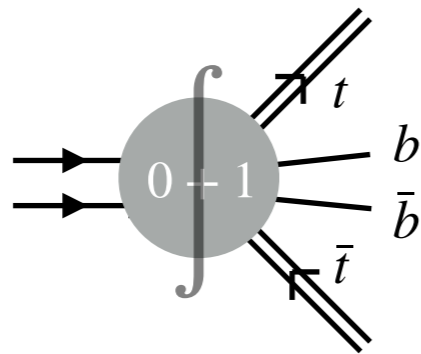
$$\int d\sigma_{\text{NWA Prod, exp}}^{NLO} =$$

The diagrammatic expansion shows the following terms:

- Term 1: Tree-level NWA cross section with a top quark propagator, multiplied by $\frac{1}{\Gamma_0}$ and Γ_0 .
- Term 2: NWA cross section with a top quark propagator, multiplied by $\frac{1}{\Gamma_0}$ and Γ_1 .
- Term 3: NWA cross section with a top quark propagator, multiplied by $\frac{1}{\Gamma_0}$ and Γ_0 .
- Term 4: NWA cross section with a top quark propagator, multiplied by $\frac{\Gamma_1}{(\Gamma_0)^2}$ and Γ_0 .
- Term 5: NWA cross section with a top quark propagator, multiplied by $\frac{1}{\Gamma_0}$ and Γ_0 .

Effects of Γ_t expansion in NWA

$$\int d\sigma_{\text{NWA Prod, exp}}^{\text{NLO}} =$$



$$\equiv \sigma_{t\bar{t}b\bar{b}, \text{incl.}}^{\text{NLO}}$$

- Rigorous expansion of Γ_t^{NLO} in NWA gets rid of “spurious” $\mathcal{O}(\alpha_s^2)$ contributions
- Note: the same procedure does not apply straightforwardly to the off-shell calculation

- The unusually large size of $\mathcal{O}(\alpha_s^2)$ effects, $\mathcal{O}(6\%)$, relates to the large size of QCD corrections to $t\bar{t}b\bar{b}$ production cross section
- These effects are well within the NLO scale uncertainties, $\mathcal{O}(20\%)$